





The structure of the proton: QCD at work from the LHC to IceCube

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Nuclear Physics Seminar

University of California in Los Angeles (UCLA)

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Particle Physics in the LHC precision era

The Higgs boson

Huge gap between weak and Plank scales?

Compositeness? Non-minimal Higgs sector?

Coupling to Dark Matter? Role in cosmological phase transitions?

Is the vacuum state of the Universe stable?









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Dark matter

- Weakly interacting massive particles? Neutrinos? Ultralight particles (axions)?
- Interactions with SM particles? Selfinteractions?
- Structure of the Dark Sector?



The Higgs boson

Huge gap between weak and Plank scales?

Compositeness? Non-minimal Higgs sector?

Coupling to Dark Matter? Role in cosmological phase transitions?

Is the vacuum state of the Universe stable?

Quarks and leptons

- Why **3 families?** Origin of **masses, mixings**?
- Origin of Matter-Antimatter asymmetry?
- Lepton Flavour Universality?
- Origin of neutrino masses? Are neutrinos Majorana or Dirac?

Dark matter

- Weakly interacting massive particles? Neutrinos? Ultralight particles (axions)?
- Interactions with SM particles? Selfinteractions?
- Structure of the Dark Sector?





Crucial information on these fundamental questions will be provided by the LHC: the **exploration of the high-energy frontier** has just started!

The inner life of protons















The proton in the spotlight

THE SCIENCES

Proton Spin Mystery Gains a New Clue



Non-zero gluon polarisation

Scientific American (2014)

Nucleon pressure

The inside of a proton endures more pressure than anything else we've seen

NEWS PARTICLE PHYSICS

For the first time, scientists used experimental data to estimate the pressure inside a proton



Science News (2018)

After 40 years of studying the strong nuclear force, a revelation

This was the year that analysis of data finally backed up a prediction, made in the mid 1970s, of a surprising emergent behaviour in the strong nuclear force



BFKL dynamics The Guardian (2017)

The proton keeps surprising us as an endless source of **fundamental discoveries**

From colliders to the cosmos



New elementary particles beyond the Standard Model?

Origins and properties of **cosmic neutrinos**?





Nature of Quark-Gluon Plasma in heavy-ion collisions?

From colliders to the cosmos



New elementary particles beyond the Standard Model?

Origins and properties of **cosmic neutrinos**?





Nature of Quark-Gluon Plasma in heavy-ion collisions?

Proton energy divided among constituents: quarks and gluons





Proton energy divided among constituents: **quarks** and **gluons**





Lattice QCD starting to also make an impact

Proton energy divided among constituents: quarks and gluons



Heavy quark content? Novel QCD dynamics?

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Proton energy divided among constituents: quarks and gluons



 $N_{\text{LHC}}(H) \sim g \otimes g \otimes \widetilde{\sigma}_{ggH}$

Parton Distributions



All-order structure: QCD factorisation theorems

g(x,Q)

Energy of hard-scattering reaction: inverse of resolution length

Probability of finding a gluon inside a

proton, carrying a fraction *x* of the proton momentum, when probed with energy *Q*

x: fraction of proton momentum carried by gluon

Dependence on *x* fixed by **non-perturbative QCD dynamics**: extract from experimental data

Energy conservation: momentum sum rule

$$\int_0^1 dx \, x \left(\sum_{i=1}^{n_f} \left[q_i((x, Q^2) + \bar{q}_i(x, Q^2)) \right] + g(x, Q^2) \right) = 1$$

Quark number conservation: valence sum rules

$$\int_0^1 dx \, \left(u(x, Q^2) + \bar{u}(x, Q^2) \right) = 2$$

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g(x,Q)

Energy of hard-scattering reaction: inverse of resolution length

Probability of **finding a gluon inside a proton**, carrying a fraction *x* of the proton

momentum, when probed with energy **Q**

x: fraction of proton momentum carried by gluon

Dependence on **Q** fixed by perturbative QCD dynamics: computed up to $\mathcal{O}(\alpha_s^4)$

$$\frac{\partial}{\partial \ln Q^2} q_i(x, Q^2) = \int_x^1 \frac{dz}{z} P_{ij}\left(\frac{x}{z}, \alpha_s(Q^2)\right) q_j(z, Q^2)$$

DGLAP parton evolution equations

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The Global QCD analysis paradigm

QCD factorisation theorems: PDF universality

$$\sigma_{lp \to \mu X} = \widetilde{\sigma}_{u\gamma \to u} \otimes u(x) \implies \sigma_{pp \to W} = \widetilde{\sigma}_{u\bar{d} \to W} \otimes u(x) \otimes \bar{d}(x)$$



Determine PDFs from deepinelastic scattering...

... and use them to compute predictions for **proton-proton collisions**

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25

A proton structure snapshop



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PDF uncertainties in the production of New Physics heavy resonances up to 100%

Due to limited coverage of the large Bjorken-x region



PDF uncertainties one of dominant theory errors in Higgs production cross-sections

any small deviations of Higgs couplings from SM predictions: smoking gun for BSM

Inclusive Higgs production rates



Snowmass 13

BSM model	Deviations in Higgs coupling to		
	W, Z weak bosons	bottom quarks	photons
New heavy Higgs boson	6%	6%	6%
Two-Higgs Doublet model	1%	10%	1%
Composite Higgs	-3%	-9%	-9%
New heavy top-like quark	-2%	-2%	+2%



Heavy bSM physics beyond the direct reach of the LHC can be parametrised in a model-independent in terms of a complete basis of higher-dimensional operators: this is the Standard Model Effective Field Theory

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{i}^{N_{d6}} \frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \sum_{j}^{N_{d8}} \frac{b_j}{\Lambda^4} \mathcal{O}_j^{(8)} + \dots ,$$

Some operators induce **growth with the partonic centre-of-mass energy**: increased sensitivity in LHC cross-sections in the TeV region

$$\sigma(\boldsymbol{E}) = \sigma_{\rm SM}(\boldsymbol{E}) \left(1 + \sum_{i}^{N_{d6}} \omega_{i} \frac{c_{i} m_{\rm SM}^{2}}{\Lambda^{2}} + \sum_{i}^{N_{d6}} \widetilde{\omega}_{i} \frac{c_{i} \boldsymbol{E}^{2}}{\Lambda^{2}} + \mathcal{O}\left(\Lambda^{-4}\right) \right)$$

enhanced sensitivity from **TeV-scale processes:** unique feature of LHC





SMEFT interpretation: from a massive particle at high energies ...



... or reflecting our limited understating of proton structure?

Hot and Cold Nuclear Matter

- Cold nuclear matter effects modify the PDFs of bound nucleons as compared to the free-proton case
- Rich QCD phenomenology: EMC effect, shadowing, non-linear evolution,
- Onset of new gluon-dominated state of matter: the Color Glass Condensate





astrophysics e.g. neutrino telescopes

Neutrino telescopes as QCD microscopes

Ultra-high energy (cosmic) neutrino - nucleus scattering: unique probe of small-x PDFs and QCD



Sensitive to **small-***x* **quarks** (and gluons via evolution) down to $\mathbf{x} \approx \mathbf{10}^{-8}$ at $\mathbf{Q} \approx \mathbf{M}_{\mathbf{W}}$

36

Bertone, Gauld, JR 18
Machine Learning in particle physics

Machine Learning in Particle Physics

- Machine Learning: families of computer algorithms that learn how to excel on a task based on a large sample of examples rather than on some a priori fixed rules
- ML algorithms are nowadays ubiquitous, from driverless cars to Amazon's purchase suggestions, to automated medical imaging recognition to beating the words best players at Go and chess
- Nowadays a standard element of the HEP toolbox as well







Artificial Neural Networks

Inspired by **biological brain models**, **Artificial Neural Networks** (ANNs) are mathematical algorithms designed to excel where domains as their evolution-driven counterparts outperforms traditional algorithms in tasks such as **pattern recognition**, **forecasting**, **classification**, ...



Artificial Neural Networks and PDFs



Generative Adversarial Networks

- New architecture for an **unsupervised neural network training** (unlabelled samples)
- Based on two independent nets that work separately and act as adversaries:
 - the Discriminator (D) undergoes training and plays the role of classifier

Ithe Generator (G) and is tasked to generate random samples that resemble real samples with a twist rendering them as fake samples.



Image generation with GANs



https://thispersondoesnotexist.com/

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42

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Recurrent Neural Networks

RNNs use as inputs not just the current "training examples" but also what they have perceived previously: **built-in notion of time ordering** useful for time-dependent functions



The output of a RNN at time time, *y(t)*, depends both on the current input example *x(t)* as well as of its previous output *y(t-1)* (or activation states of hidden neutrons at *t-1*)

Al-based artificial text generation

SYSTEM PROMPT (HUMAN-WRITTEN) In a shocking finding, scientist discovered a herd of unicorns living in a remote, previously unexplored valley, in the Andes Mountains. Even more surprising to the researchers was the fact that the unicorns spoke perfect English.

MODEL COMPLETION (MACHINE-WRITTEN, 10 TRIES) The scientist named the population, after their distinctive horn, Ovid's Unicorn. These four-horned, silver-white unicorns were previously unknown to science.

Now, after almost two centuries, the mystery of what sparked this odd phenomenon is finally solved.

Dr. Jorge Pérez, an evolutionary biologist from the University of La Paz, and several companions, were exploring the Andes Mountains when they found a small valley, with no other animals or humans. Pérez noticed that the valley had what appeared to be a natural fountain, surrounded by two peaks of rock and silver snow.

https://openai.com/

Pérez and the others then ventured further into the valley. "By the time we reached the top of one peak, the water looked blue, with some crystals on top," said Pérez.

Pérez and his friends were astonished to see the unicorn herd. These creatures could be seen from the air without having to move too much to see them — they were so close they could touch their horns.

While examining these bizarre creatures the scientists discovered that the creatures also spoke some fairly regular English. Pérez stated, "We can see, for example, that they have a common 'language,' something like a dialect or dialectic."

The optimisation conundrum

(Stochastic) Gradient Descent



- Deterministic algorithms: follow the gradient of the cost function
- Evolutionary algorithms: act on population of solutions with random mutations and selection

A challenge for any QCD global analysis is the efficient exploration of huge parameter space (hundreds/thousands free params)

Several pitfalls to be avoided: under-fitting, over-fitting, **local minima**, saddle points,

Genetic Algorithms







Generation 2

Generation 5



Generation 3



Generation 6



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Parton Distributions with Theoretical Uncertainties

Based on NNPDF Collaboration: R. Abdul Khalek, R. D. Ball, S. Carrazza, S. Forte, T. Giani, Z. Kassabov, R. L. Pearson, E. R. Nocera, J. Rojo, L. Rottoli, M. Ubiali, C. Voisey, M. Wilson

arXiv:1905.04311, submitted to PRL arXiv:1906.10698, submitted to EPJC

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PDF uncertainties

PDF uncertainties receive contributions from different sources:



QCD uncertainties in PDF fits

Standard global PDF fits are based on fixed-order QCD calculations

$$\sigma = \alpha_s^p \sigma_0 + \alpha_s^{p+1} \sigma_1 + \alpha_s^{p+2} \sigma_2 + \mathcal{O}(\alpha_s^{p+3})$$

The truncation of the perturbative series has associated a theoretical uncertainty: **Missing Higher Order (MHO)** uncertainty

How severe is **ignoring MHOUs** in modern global PDFs fits?



A theoretical covariance matrix

Construct a **theory covariance matrix** from **scale-varied cross-sections** and combine it with the experimental covariance matrix

$$\chi^{2} = \frac{1}{N_{\text{dat}}} \sum_{i,j=1}^{N_{\text{dat}}} \left(D_{i} - T_{i} \right) \left(C + S \right)_{ij}^{-1} \left(D_{j} - T_{j} \right)$$
experimental theoretical

assumption: theory errors are Gaussianly distributed around true value

Formally the theory covariance matrix is defined as

$$S_{ij} = \left\langle (\mathcal{T}_i - T_i)(\mathcal{T}_j - T_j) \right\rangle \equiv \left\langle \Delta_i \Delta_j \right\rangle$$

true result
actual calculation
How to estimate these **theory systematic shifts**?

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A theoretical covariance matrix

Here we use **scale variations** to estimate the MHOUs



note: renormalisation scale variations are only correlated within the same process



Different prescriptions for scale variations possible: Need to validate which ones exhibit the best performance

Point prescriptions



The theory covariance matrix

covariance matrices



Theory Covariance matrix (9 pt)

Rich pattern of **theory-induced correlations**: Absent if only experimental errors considered

Experimental Covariance Matrix

The theory covariance matrix

correlation matrices



Experimental + Theory Correlation Matrix (9 pt)

Rich pattern of **theory-induced correlations**: Absent if only experimental errors considered

Experimental Correlation Matrix

Validation

Systematic validation of NLO theory covariance matrix on the `exact' result, the NLO=>NNLO shift, for O(3000) data points of the global PDF fit



Scale variations: good estimate of MHOU for processes of relevance in PDF fits

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Impact on PDFs



10⁻²

10⁻²

10⁻¹

10⁻¹

55

Impact for LHC phenomenology



Depending on process, main consequence of **MHOUs in PDF fit for LHC pheno** is shift in central values, increase in overall PDF uncertainties, or both

Can New Physics Hide Inside the Proton?

Based on S. Carrazza, C. Degrande, S. Iranipour, J. Rojo, M. Ubiali

arXiv:1905.05215, submitted to PRL

Heavy bSM physics beyond the direct reach of the LHC can be parametrised in a model-independent in terms of a complete basis of higher-dimensional operators: this is the Standard Model Effective Field Theory

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{i}^{N_{d6}} \frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \sum_{j}^{N_{d8}} \frac{b_j}{\Lambda^4} \mathcal{O}_j^{(8)} + \dots,$$

Some operators induce **growth with the partonic centre-of-mass energy**: increased sensitivity in LHC cross-sections in the TeV region

$$\sigma(\boldsymbol{E}) = \sigma_{\rm SM}(\boldsymbol{E}) \left(1 + \sum_{i}^{N_{d6}} \omega_i \frac{c_i m_{\rm SM}^2}{\Lambda^2} + \sum_{i}^{N_{d6}} \widetilde{\omega}_i \frac{c_i \boldsymbol{E}^2}{\Lambda^2} + \mathcal{O}\left(\Lambda^{-4}\right) \right)$$

enhanced sensitivity from **TeV-scale processes:** unique feature of LHC

NNPDF3.1



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59





60

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62

Naive approach

Separate LHC data into input for PDF fits and input for SMEFT studies?



Can we do better?

63

Simultaneous PDF+SMEFT fits

Our goal: constrain **simultaneously** both the PDFs and SMEFT degrees of freedom

Proof of concept: DIS-only fits where SM **augmented** by four *d=6* SMEFT operators

$$\mathscr{L}_{\text{SMEFT}} = \mathscr{L}_{\text{SM}} + \sum_{q=u,d,s,c} \frac{a_q}{\Lambda^2} \left(\bar{l}_R \gamma^{\mu} l_R \right) \left(\bar{q}_R \gamma_{\mu} q_R \right)$$

which can arise *e.g.* from a **Z' boson** with non-universal couplings to quarks

These SMEFT operators modify the DIS structure functions and thus affect the PDF fit

$$\Delta F_2^{\text{SMEFT}} \supset \frac{x}{12e^4} \begin{pmatrix} 4a_u e^2 \frac{Q^2}{\Lambda^2} (1 + 4K_Z \sin^4 \theta_W) + 3a_u^2 \frac{Q^4}{\Lambda^4} \end{pmatrix} (u + \bar{u})$$

$$\stackrel{\text{SMEFT effects enhanced by } Q^2:$$

$$\stackrel{\text{constrain from HERA data!}}{\text{from interference with SM}} \qquad \text{from squared amplitude}$$

Impact on the PDFs

For a large region of the allowed parameter space,

SMEFT effects can be partially (but not completely) reabsorbed into the PDFs

NNPDF3.1 DIS-only, Q = 10 GeV



Fingerprinting BSM effects

Tell-tale sign of SMEFT effects: rapid variation with Q (DGLAP evolution slower)



Fingerprinting BSM effects

We can compare bounds on SMEFT degrees of freedom in the joint fit as compared to the usual approach where PDFs are kept fixed



90%CL allowed region

Ultimate goal (HL-LHC timescale): **simultaneous PDF+SMEFT global analysis!**

Neutrino Telescopes as QCD Microscopes

Based on V. Bertone, R. Gauld, J. Rojo arXiv:1808.02034 (JHEP)

The small-x gluon from HERA data

- Small-x gluon unconstrained: information from HERA ends for x<10⁻⁴
- Large uncertainties in global fits

Need processes covering x<10⁻⁴ region





Forward charm production



Forward charm production

- Include LHCb D meson production at 5, 7, 13 TeV
- Fit normalised distributions & ratios between CoM energies to reduce MHOUs

$$N_X^{ij} = \frac{d^2\sigma(\text{X TeV})}{dy_i^D d(p_T^D)_j} \left/ \frac{d^2\sigma(\text{X TeV})}{dy_{\text{ref}}^D d(p_T^D)_j} \right|_{T_{\text{ref}}}$$
$$R_{13/X}^{ij} = \frac{d^2\sigma(13 \text{ TeV})}{dy_i^D d(p_T^D)_j} \left/ \frac{d^2\sigma(\text{X TeV})}{dy_i^D d(p_T^D)_j} \right|_{T_{\text{ref}}}$$

gluon PDF uncertainties reduced by factor 10 at $x \approx 10^{-6}$



Excellent description of all LHCb datasets

and ratios (after errata corrected)

$N_5(84)$	$N_{7}(79)$	$N_{13}(126)$	$R_{13/5}(107)$	$R_{13/7}(102)$
1.97	1.21	2.36	1.36	0.80
0.86	0.72	1.14	1.35	0.81
1.31	0.91	1.58	1.36	0.82
0.74	0.66	1.01	1.38	0.80
1.08	0.81	1.27	1.29	0.80
1.53	0.99	1.73	1.30	0.81
1.07	0.81	1.34	1.35	0.81
0.82	0.70	1.07	1.35	0.81
0.84	0.71	1.10	1.36	0.81

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BFKL dynamics at small-x

- QCD calculations in the DGLAP factorisation framework successful in describing data from proton-proton and electron-proton collisions
- Need to go beyond DGLAP: at small-x, logarithmically enhanced terms in 1/x become dominant and need to be resummed to all orders
- BFKL (high-energy, small-x) resummation can be matched to DGLAP collinear framework and included into PDF fits

$$\begin{array}{ll} \begin{array}{ll} \textbf{DGLAP} \\ \textbf{Evolution in } Q^2 \end{array} & \begin{array}{l} \frac{\partial}{\partial \ln Q^2} f_i(x,Q^2) = \int_x^1 \frac{dz}{z} P_{ij}\left(\frac{x}{z},\alpha_s(Q^2)\right) f_j(z,Q^2) \\ \end{array} \\ \begin{array}{l} \textbf{BFKL} \\ \textbf{Evolution in } x \end{array} & \begin{array}{l} \frac{\partial}{\partial \ln 1/x} f_+(x,Q^2) = \int_0^\infty \frac{d\nu^2}{\nu^2} K\left(\frac{Q^2}{\nu^2},\alpha_s(Q^2)\right) f_+(x,\nu^2) \\ \end{array} \\ \begin{array}{l} \textbf{ABF, CCSS, TW} \\ \textbf{+ others, 94-08} \end{array} & P_{ij}^{N^k LO + N^h LLx}(x) = P_{ij}^{N^k LO}(x) + \Delta_k P_{ij}^{N^h LLx}(x) \end{array}$$
BFKL dynamics at small-x



Monitor the **fit quality** as one includes more data from the **small-***x* **region**

Best description of **small-***x***HERA data** only possible with **BFKL effects!**

Forward charm production revisited

LHCb D meson production included in NNPDF3.1sx (N)NLO+NLLx fits

Similar reduction of gluon PDF errors at **small-***x* + **increase in central value**



Neutrino telescopes

Ultra-high energy (UHE) neutrinos: novel window to the extreme Universe



Neutrino telescopes as QCD microscopes

signal: cosmic neutrino - nucleus scattering

background: prompt charm production





Sensitive to **small-***x* **quarks** (and thus gluons via evolution) down to $\mathbf{x} \approx \mathbf{10^{-8}}$ and $\mathbf{Q} \approx \mathbf{M_W}$

Sensitive to small-x gluons down to $x \approx 10^{-6}$ and $Q \approx M_{charm}$ in the centre-of-mass frame

Neutrino telescopes as QCD microscopes

signal: cosmic neutrino - nucleus scattering

background: prompt charm production



UHE neutrino-nucleus cross-section



State-of-the-art predictions for ultra-high energy neutrino interactions

- BFKL small-x effects in PDFs and deep-inelastic structure functions
- Constraints on small-x PDFs from LHCb charm production
- Accounting for nuclear corrections and heavy-quark-initiated contributions

UHE neutrino-nucleus cross-section



- Differences both at intermediate (updated PDFs, improved treatment of heavy quarks) and high energies (LHCb constraints, BFKL effects)
- Nuclear effects important: constrain them with LHCb charm production in p+Pb
- IceCube and other neutrino telescopes are the ultimate QCD microscopes!

Towards Ultimate PDFs from Collider Data

Based on R. Abdul Khalek, S. Bailey, J. Gao, L. Harland-Land, J. Rojo

arXiv:1810.03639 (EPJC) *arXiv:1906.10127* (Submitted to SciPost)

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Fully complementary in terms of PDF constraints, possible synchronous operation

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81

Exploit novel facilities for precision studies of the proton structure







HL-LHC: generate our own pseudo-data based on extrapolations from Run II

Hessian profiling

Hessian Profiling: analytic minimisation of the figure of merit:

$$\chi^{2}(\beta_{\exp},\beta_{\mathrm{th}}) = \frac{1}{\left(\delta_{\mathrm{tot},i}^{\exp}\right)^{2}} \sum_{i=1}^{N_{\mathrm{dat}}} \left(\sigma_{i}^{\exp} + \sum_{j} \Gamma_{ij}^{\exp} \beta_{j,\exp} - \sigma_{i}^{\mathrm{th}} + \sum_{k} \Gamma_{ik}^{\mathrm{th}} \beta_{k,\mathrm{th}}\right)^{2} + \sum_{j} \beta_{j,\exp}^{2} + T^{2} \sum_{k} \beta_{k,\mathrm{th}}^{2} + \sum_{j} \beta_{j,\exp}^{2} + T^{2} \sum_{k} \beta_{j,\exp}^{2} + T^{2} \sum_{k} \beta_{j,\exp}^{2} + T^{2} \sum_{k} \beta_{j,\exp}^{2} + T^{2} \sum_{k} \beta_{j,\exp}^{2} + T^{2} \sum_{j} \beta_{j,\exp}^{2} + T^{2} \sum$$

To mimic the effects of the correlation model for the HL-LHC pseudo-data, we rescale the total experimental systematic error by a suitable factor *f_{corr}*

$$\delta_{\text{tot},i}^{\text{exp}} \equiv \left(\left(\delta_{\text{stat},i}^{\text{exp}} \right)^2 + \left(f_{\text{corr}} \times f_{\text{red}} \times \delta_{\text{sys},i}^{\text{exp}} \right)^2 \right)^{1/2}$$
$$\delta_{\text{stat},i}^{\text{exp}} = \left(f_{\text{acc}} \times N_{\text{ev},i} \right)^{-1/2}$$

For the LHeC pseudo-data the correlation model already available for inclusive structure functions

HL-LHC: focus on process where systematics are not already limiting factor

Quantify the ultimate **PDF constraining power** of HL-LHC



HL-LHC measurements will be specially useful to constrain gluon and

quark flavour separation, including strangeness

Pinning down **antiquarks at large-***x* seems challenging, suggestions welcome!

Quantify the ultimate **PDF constraining power** of HL-LHC and LHeC



A reduction of PDF uncertainties by up to a factor 10 could be within reach

Quantify the ultimate **PDF constraining power** of HL-LHC and LHeC



A reduction of PDF uncertainties by up to a factor 10 could be within reach

Impact on phenomenology



Summary and outlook

The accurate determination of the **quark and gluon structure of the proton** is an essential ingredient for **LHC phenomenology** and **beyond**

- Recent progress in **longstanding issues**: QCD uncertainties on PDFs, lattice QCD constraints, strange and charm content of the proton, connection with astrophysics, ...
- ... but also wrapping up: QED effects on PDFs, BFKL dynamics in HERA data,
- Long-term goals of QCD analyses: exploiting future facilities (HL-LHC, LHeC, EIC) and integrating its multiple dimensions: joint fits of (p)PDFs + FFs + nPDFs + TMDs
- Simultaneous constraints of both proton structure and BSM effects will play a central role for the interpretation of the HL-LHC measurements

The fascinating study of the proton structure never stops surprising us, stay tuned!

Summary and outlook

The accurate determination of the **quark and gluon structure of the proton** is an essential ingredient for **LHC phenomenology** and **beyond**



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