



A global SMEFT analysis extended with HL-LHC and FCC-ee projections

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SMEFiT projections for future colliders

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Combined SMEFT interpretation of Higgs, diboson, and top quark data from the LHC

The SMEFiT Collaboration:

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Abstract

We present a global interpretation of Higgs, diboson, and top quark production and decay measurements from the LHC in the framework of the Standard Model Effective Field Theory (SMEFT) at dimension six. We constrain simultaneously 36 independent directions in its parameter space, and compare the outcome of the global analysis with that from individual and

Tommaso Giani, Giacomo Magni, and Juan Rojo Department of Physics and Astronomy, Vrije Universiteit Amsterdam, NL-1081 HV Amsterdam, The Netherlands Nikhef Theory Group, Science Park 105, 1098 XG Amsterdam, The Netherlands Abstract The Standard Model Effective Field Theory (SMEFT) provides a robust framework to interpret experimental measurements in the context of new physics scenarios while minimising assumptions on the nature of the underlying UV-complete theory. We present the PYTHON open source SMEF1T framework, designed to carry out parameter inference in the SMEFT within a global analysis of particle physics data. SMEF1T is suitable for inference problems involving a large number of EFT degrees of freedom, without restrictions on their functional dependence in the fitted observables, can include UV-inspired restrictions in the parameter space, and implements arbitrary rotations between operator bases. Posterior distributions are determined from two complementary approaches, Nested Sampling and Monte Carlo optimisation. SMEF1T is released together with documentation, tutorials, and post-analysis reporting tools, and can be used to carry out state-of-the-art EFT fits of Higgs, top quark, and electroweak production data. To illustrate its functionalities, we reproduce the results of the recent ATLAS EFT interpretation of Higgs and electroweak data from Run II and demonstrate how equivalent results are obtained in two different operator bases.

Public code paper

SMEFIT: a flexible toolbox for global interpretations of particle physics data with effective field theories

The automation of SMEFT-Assisted Constraints on **UV-Complete Models**

Nikhef-2022-023

Automation UV matching

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ep-ph]	ABSTRACT: The ongoing Effective Field Theory (EFT) program at the LHC and elsewhere is motivated by
6	streamlining the connection between experimental data and UV-complete scenarios of heavy new physics
5	beyond the Standard Model (BSM). This connection is provided by matching relations mapping the Wilson
<u>Ч</u>	coefficients of the EFT to the couplings and masses of UV-complete models. Building upon recent work on
	the automation of tree-level and one-loop matching in the SMEFT, we present a novel strategy automating
.04523v1	the constraint-setting procedure on the parameter space of general heavy UV-models matched to dimension- $% \mathcal{A}$
ń	six SMEFT operators. A new Mathematica package, MATCH2FIT, interfaces MATCHMAKEREFT, which
	derives the matching relations for a given UV model, and SMEFIT, which provides bounds on the Wilson
4	coefficients by comparing with data. By means of this pipeline and using both tree-level and one-loop
0.	matching, we derive bounds on a wide range of single- and multi-particle extensions of the SM from a
IV:2309	global dataset composed by LHC and LEP measurements. Whenever possible, we benchmark our results
ň	with existing studies. Our framework realises one of the main objectives of the EFT program in particle
2	physics: deploying the SMEFT to bypass the need of directly comparing the predictions of heavy UV
$1 \leq 1$	models with experimental data.
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SMEFiT projections for future colliders

Extend SMEFiT2.0 with recent LHC Run-2 datasets on top, diboson, and Higgs production

✓ New, dedicated implementation of the constraints from LEP's EWPOs

Include tailored projections for HL-LHC pseudo-data based on extrapolating Run-2 data

☑ Include ``Snowmass" projections for FCC-ee & CEPC pseudo-data, updated with FCC midterm Feasibility Report

☑ Results both Wilson coefficient and matched to UVcomplete models that have SMEFT as low-energy limit

Public code, data, and theory: user-friendly and straightforward to apply to any other future collider

extensive documentation and **example runcards**, fully reproducible results

SMEFiT3.0 & Projections for HL-LHC & FCC-ee

Mapping the SMEFT at High-Energy Colliders: from LEP and the (HL-)LHC to the FCC-ee

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ABSTRACT: We present SMEFIT3.0, an updated global SMEFT analysis of Higgs, top quark, and diboson data from the LHC complemented by electroweak precision observables (EWPOs) from LEP and SLD. We consider the most recent inclusive and differential measurements from the LHC Run II, alongside with a novel implementation of the EWPOs based on independent calculations of the relevant EFT contributions. We estimate the impact of HL-LHC measurements on the SMEFT parameter space, when added on top of SMEF1T3.0, by means of dedicated projections extrapolating from Run II data. We quantify the significant constraints that measurements from future high-energy circular electron-positron colliders would impose on both the SMEFT parameter space and on representative UV-complete models. Our analysis considers projections for the FCC-ee and the CEPC based on the latest running scenarios and including Z-pole EWPOs, fermion-pair, Higgs, diboson, and top quark production, using optimal observables for both the W^+W^- and the $t\bar{t}$ channels. The framework presented in his work may be extended to other proposed colliders and running scenarios, providing a timely input to ongoing studies towards future high-energy particle colliders.

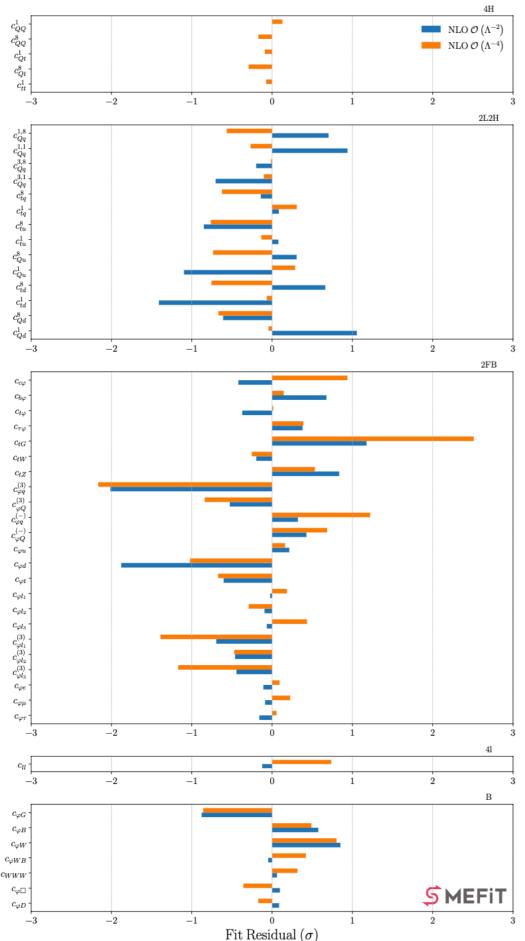
to appear Friday on the arXiv

Baseline SMEFT analysis

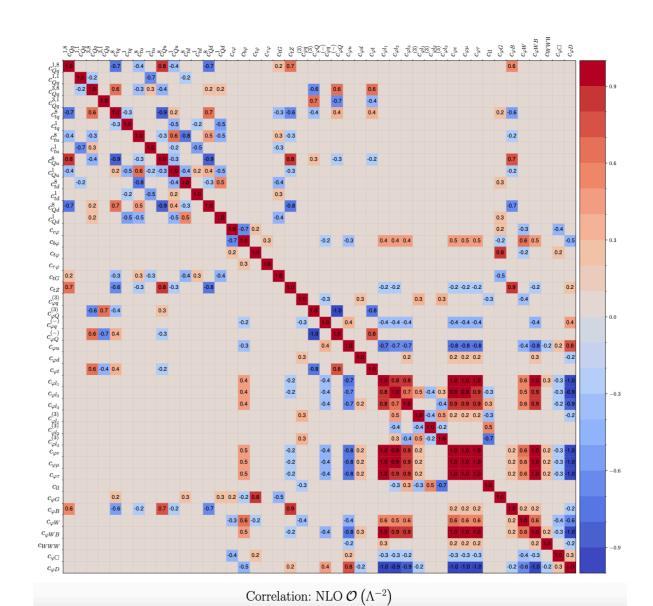
- 445 fitted cross-sections from top (189), Higgs (129), and diboson production (81), complemented with LEP's EWPOs (44 data points)
- § 50 (45) Wilson coefficients in the quadratic (linear) EFT fits
- Posterior distributions provided by Nested Sampling (Bayesian inference). Linear fits cross-checked with analytic solution.
- Flavour assumptions: $U(2)_q \otimes U(3)_q \otimes U(2)_u \otimes (U(1)_{\ell} \otimes U(1)_e)^3$

Catagoria	Dresser	$n_{ m dat}$						
Category	Processes	SMEFIT2.0	SMEFIT3.0					
	$t\bar{t} + X$	94	115					
	$tar{t}Z,tar{t}W$	14	21					
	$t\bar{t}\gamma$	-	2					
Top quark production	single top (inclusive)	27	28					
	tZ, tW	9	13					
	$t\bar{t}t\bar{t}$, $t\bar{t}b\bar{b}$	6	12					
	Total	150	189					
	Run I signal strengths	22	22					
Higgs production	Run II signal strengths	40	36 (*)					
and decay	Run II, differential distributions & STXS	35	71					
	Total	97	129					
	LEP-2	40	40					
Diboson production	LHC	30	41					
	Total	70	81					
Z-pole EWPOs	LEP-2	-	44					
Baseline dataset	Total	317	445					

Baseline SMEFT analysis

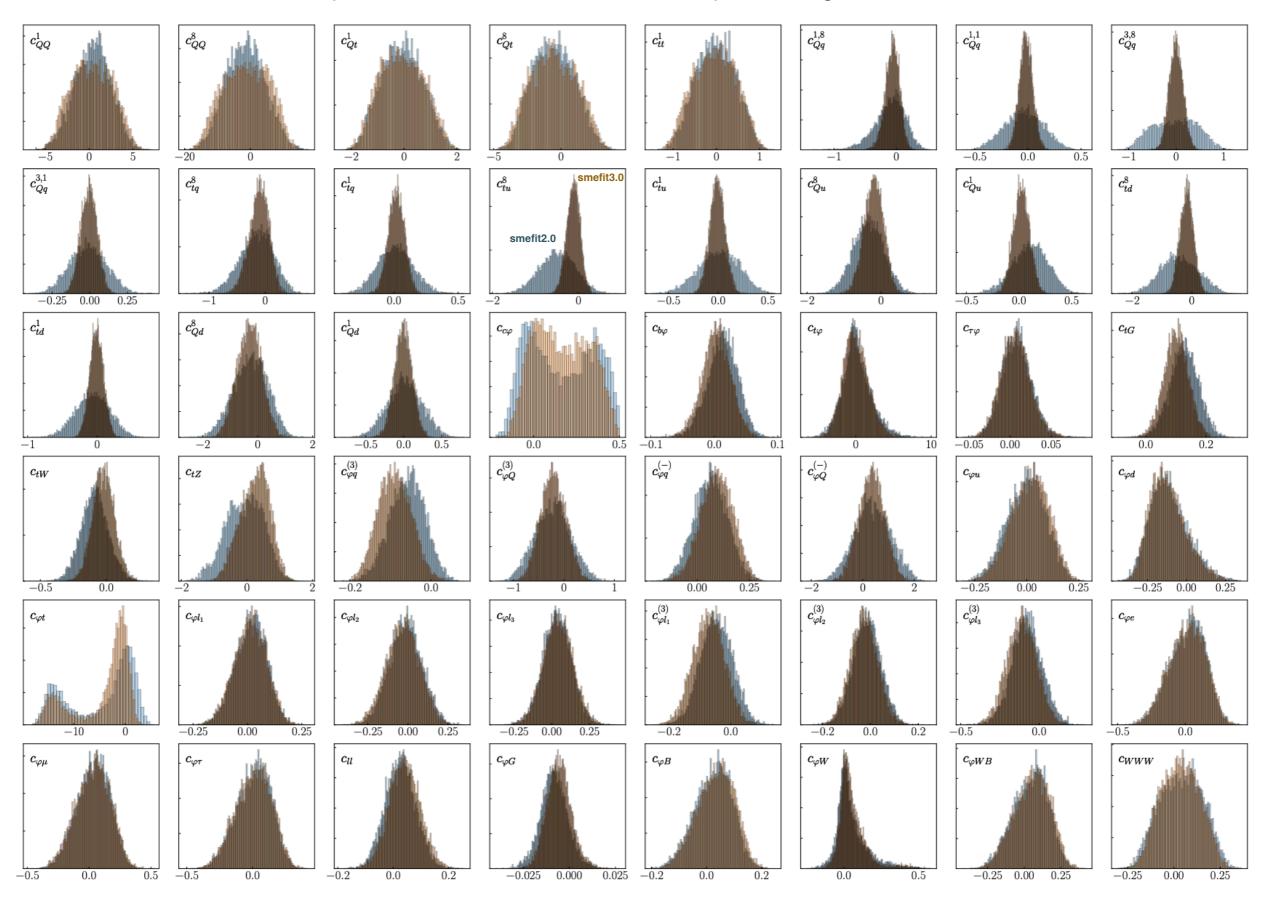


- Good agreement with the SM expectations, with a few (well understood) exceptions
- Large correlations between fitted coefficients, partially washed out in the quadratic EFT fits
- Full posterior distribution available, can be marginalised to the 1-coeff, 2-coeff, 3-coeff level

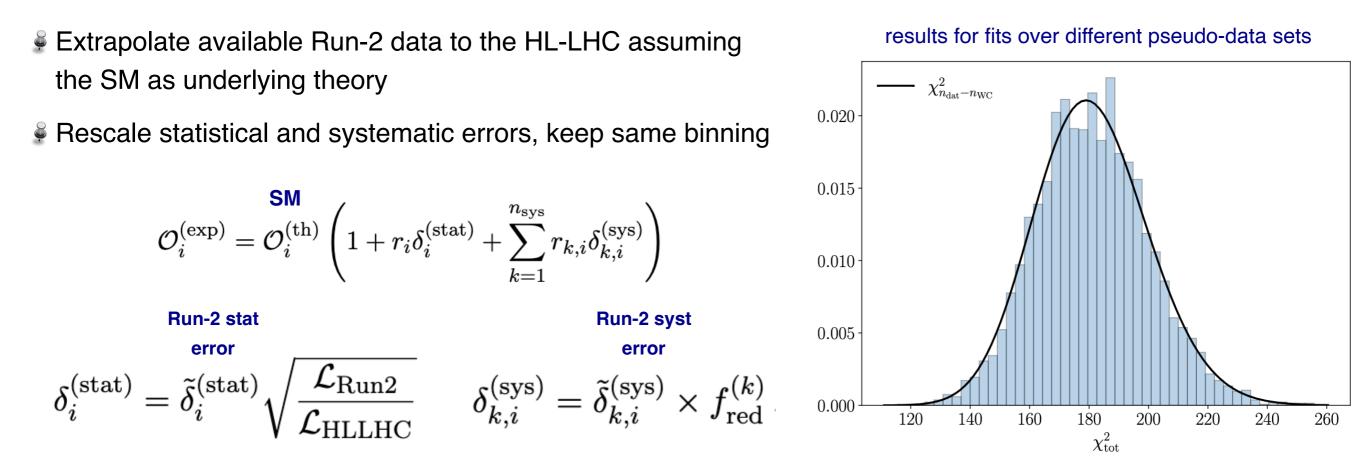


Baseline SMEFT analysis

Impact of recent LHC data in the quadratic global fit



HL-LHC projections



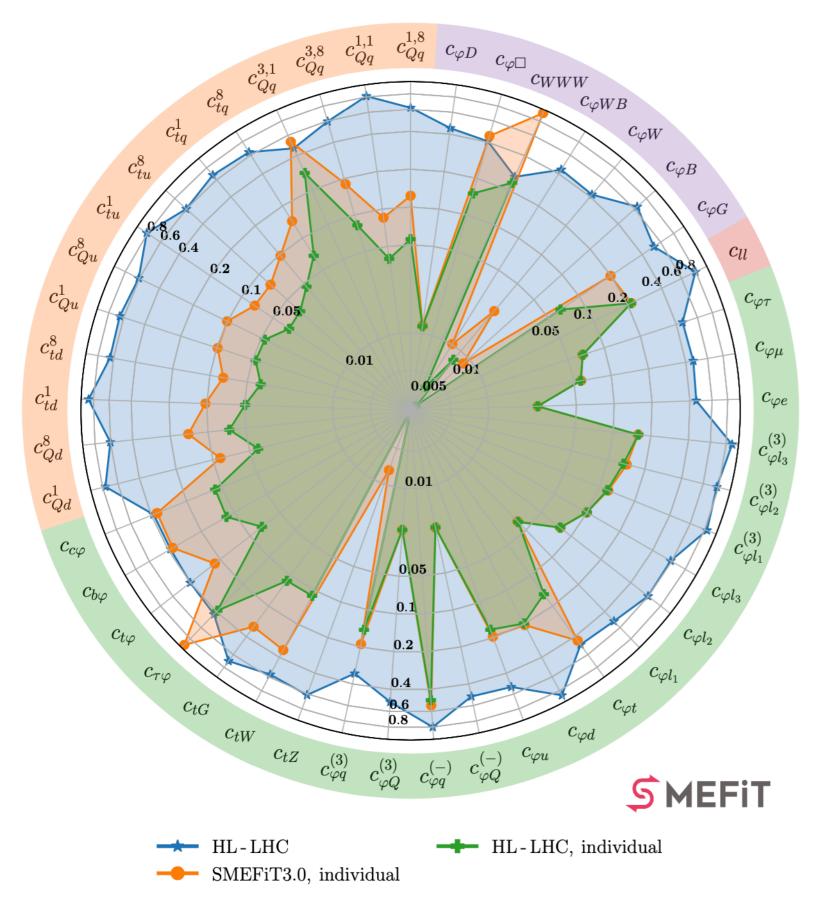
pro: keep in sync global fit to actual data and HL-LHC projections

Con: does not exploit HL-LHC optimisation e.g. more bins in high-pT region, multi-differential measurements todo list: include the available tailored projections for HL-LHC observables

Goal: provide a **realistic baseline**, in preparation of the subsequent inclusion of the FCC-ee and CEPC projections, in terms of constraints in the SMEFT parameter space

HL-LHC projections

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-2})$, Marginalised

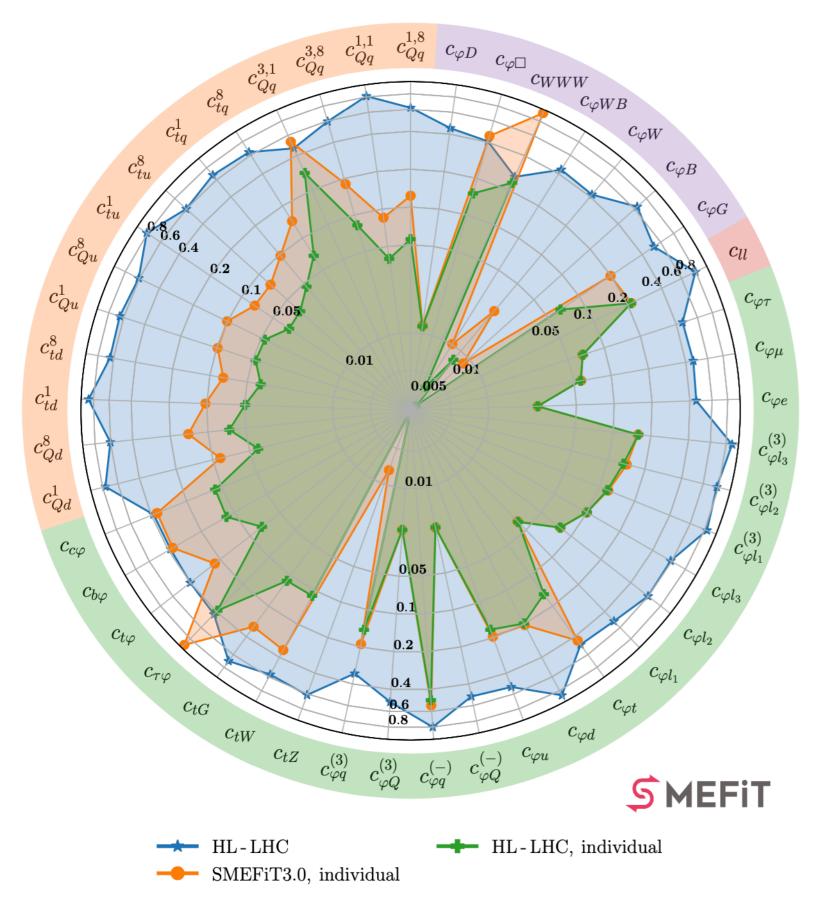


HL-LHC: bounds on the EFT coefficients improved by up to a factor 3 as compared to current dataset

- Large differences between marginalised bounds from global fit and individual (one-parameter fits) bounds: the latter are overly optimistic
- HL-LHC reach can be improved with inclusion of optimised observables
- Qualitatively similar picture in the quadratic fits

HL-LHC projections

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-2})$, Marginalised



HL-LHC: bounds on the EFT coefficients improved by up to a factor 3 as compared to current dataset

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FCC-ee projections

FCC-ee and CEPC projections follow the Snowmass study (except for 4 IPs for FCC-ee, following feasibility study)

Energy (\sqrt{s})	$\mathcal{L}_{\mathrm{int}}$ (Ru			
Energy (\sqrt{s})	FCC-ee (4 IPs)	$\mathcal{L}_{ ext{FCC-ee}}/\mathcal{L}_{ ext{CEPC}}$		
91 GeV (Z-pole)	$300 \text{ ab}^{-1} (4 \text{ years})$	$100 \text{ ab}^{-1} (2 \text{ years})$	3	
161 GeV $(2 m_W)$	$20 \text{ ab}^{-1} (2 \text{ years})$	$6 \text{ ab}^{-1} (1 \text{ year})$	3.3	
$240~{ m GeV}$	10 ab^{-1} (3 years)	$20 \text{ ab}^{-1} (10 \text{ years})$	0.5	
$350~{ m GeV}$	$0.4 \text{ ab}^{-1} (1 \text{ year})$	$0.2~{ m ab}^{-1}$	2	
$\begin{array}{ c c c c c }\hline 365 {\rm GeV} (2 m_t) \end{array}$	3 ab^{-1} (4 years)	1 ab^{-1} (5 years)	3	

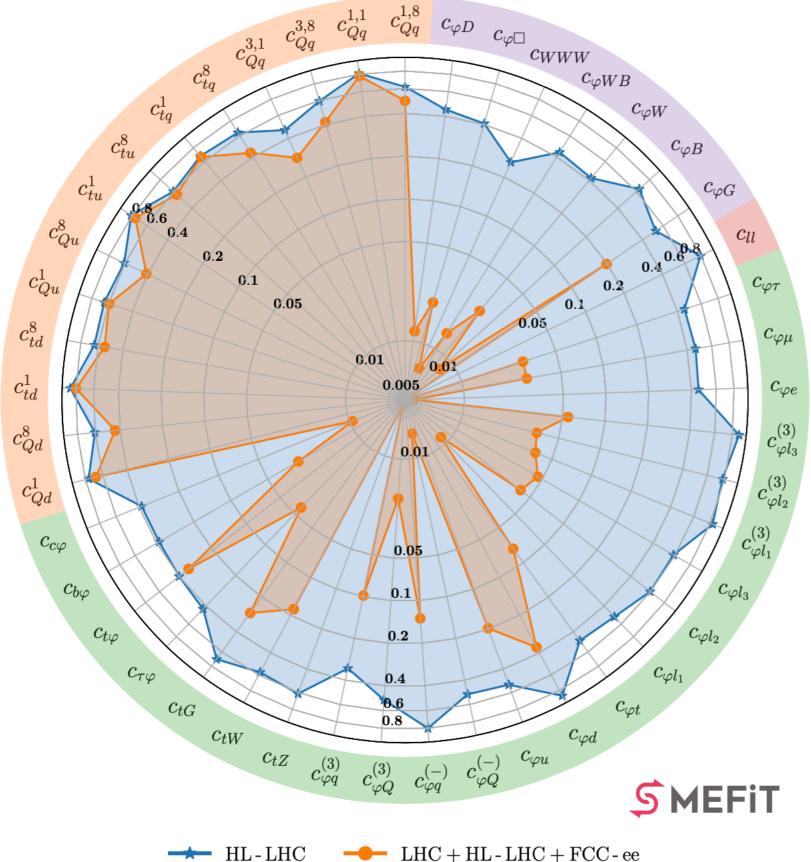
Variations upon this configuration would be easy to implement

In our analysis, we consider five different classes of observables that are accessible at high-energy circular electron-positron colliders such as the FCC-ee and the CEPC. These are the EWPOs at the Z-pole; light fermion (up to b quarks and τ leptons) pair production; Higgs boson production in both the hZ and $h\nu\nu$ channels; gauge boson pair production; and top quark pair production. Diboson (W^+W^-) production becomes available at $\sqrt{s} = 161$ GeV (WW threshold), Higgs production opens up at $\sqrt{s} = 240$ GeV, and top quark pair production is accessible starting from $\sqrt{s} = 350$ GeV, above the $t\bar{t}$ threshold.

Among these processes, the Z-pole EWPOs, light fermion-pair, W^+W^- , and Higgs production data are included at the level of inclusive cross-sections, accounting also for the corresponding branching fractions. The complete list of observables considered, together with the projected experimental uncertainties entering the fit, are collected in App. E. For diboson and top quark pair production, we consider also unbinned normalised measurements within the optimal observables approach, described in App. F. We briefly review below these groups of processes.

FCC-ee projections

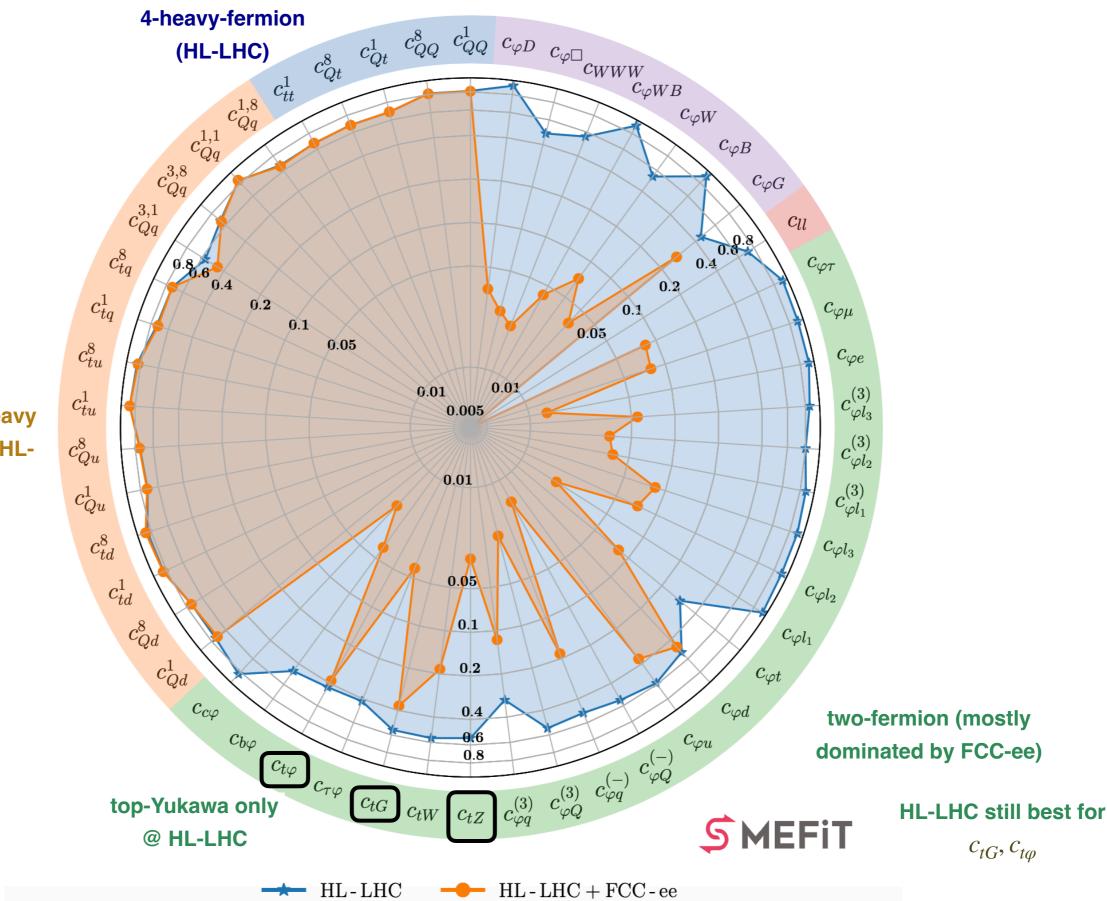
Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-2})$, Marginalised



- FCC-ee: huge improvements (up to factor 100) for most coefficients
- Clear impact on two-fermion, purely bosonic, and four-lepton operators
- Four-fermion operators involving top quarks are unaffected by FCC-ee
- Note: these are global marginalised
 bounds. If one performs individual (one-parameter) fits, impact of FCC-ee is even stronger.

Top couplings

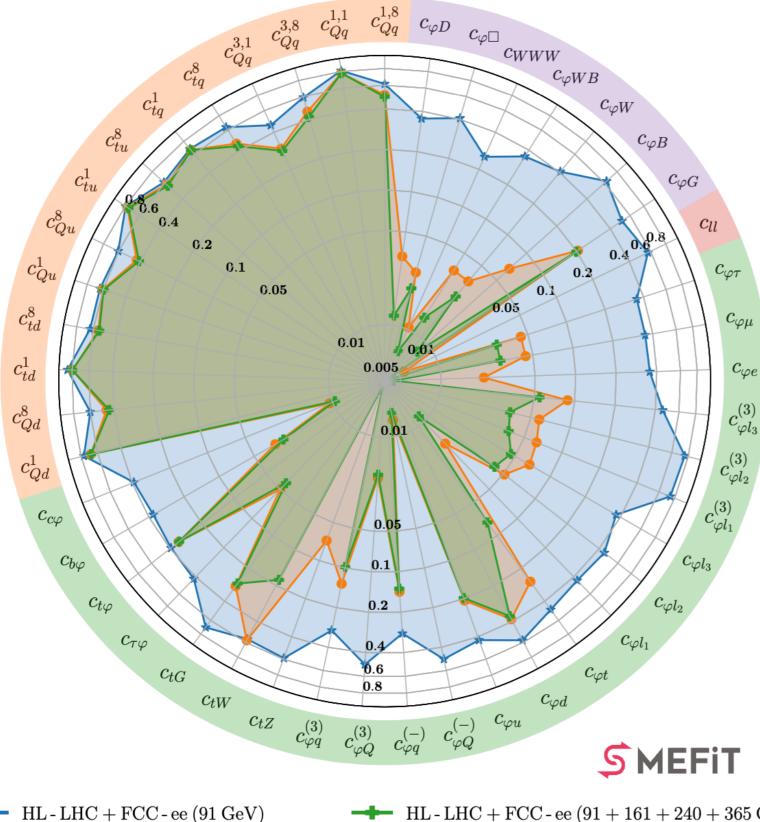
Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-4})$, Marginalised



2-light-2-heavy 4-fermion (HL-LHC)

FCC-ee projections

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-2})$, Marginalised



HL-LHC + FCC-ee (91 + 240 GeV)

Study impact of sequentially adding FCC-ee datasets for different energies

Sombining the **Z-pole run** with the Higgs factory run at 240 GeV dominates the final reach

Sean try **any other combination**, also to help determining which run order is most effective at least in terms of the SMEFT

HL - LHC + FCC - ee (91 + 161 + 240 + 365 GeV)

Fisher Information analysis

	LEP	$t\bar{t}$ 8 TeV	$t\bar{t}$ 13 TeV	$t\bar{t}\gamma$	$t\bar{t}W$	$t\bar{t}Z$	t 8 TeV	t 13 TeV	tW	tZ	$t\bar{t}A_c$	W helicities	$t\bar{t}t\bar{t} + t\bar{t}b\bar{b}$	Higgs-run I	Higgs-run II	VV	$t\bar{t}$ 13 TeV HL-LHC	HT-THC	$t\bar{t}Z$ HL-LHC	t 13 TeV HL-LHC	tw HL-LHC	tz HL-LHC	$t\bar{t}A_c$ HL-LHC	W helicities HL-LHC	$t\bar{t}t\bar{t} + t\bar{t}b\bar{b}$ HL-LHC	Higgs HL-LHC	VV HL-LHC	FCC-ee 91 GeV	FCC-ee 161 GeV	FCC-ee 240 GeV	FCC-ee 365 GeV	 100
c_{QQ}^1													14.0												86.0							100
$\begin{smallmatrix} 1 \\ c_{QQ}^{1} \\ c_{QQ}^{8} \\ c_{QQ}^{1} \\ c_{Qq}^{1} \\ c_{Qq}^{1} \\ c_{Qq}^{1} \\ c_{Qq}^{1} \\ c_{Qq}^{1} \\ c_{Qq}^{3,8} \\$													15.1		0.0										84.9	0.0						
c_{Qt}^1													18.1												81.9						_	
c_{Qt}^8													14.1		0.0										85.9	0.0						
c_{tt}^1													14.0												86.0						_	Ļ
$c_{Qq}^{1,8}$		0.4	8.4	0.2	1.6	1.3					9.1		0.0	0.0	0.1		22.7	7.9	6.3				41.7		0.1	0.1						
$c_{Qq}^{1,1}$			10.4								11.6		0.0				31.2						46.4		0.2						_	
$c_{Qq}^{3,8}$			2.2	0.3	1.9	1.0					13.6			0.0	0.1		4.3	9.2	4.6	1.3			59.6		0.1							
$c_{Qq}^{3,1}$			0.0				15.2	7.7		4.8	0.1		0.0		0.0		0.1			40.0		31.6			0.0	0.0					_	
c_{tq}^{s}			6.9	1.0	4.1	2.3					8.1		0.1	0.0	0.3			20.1	10.4				38.6			0.1						- 80
c_{tq}^1			10.1								12.3		0.0				29.1						48.2		0.1							
c_{tu}^8			8.9	0.3		0.1					13.5		0.0	0.0	0.1		14.9		0.8				60.7		0.2	0.1					-	
c_{tu}^1			8.9								12.7		0.0				26.9						51.1		0.2							
c_{Qu}^8			3.7	2.5		1.0					13.7		0.1	0.0	0.4		6.9		5.2				64.8 48.5		0.7	0.2					-	
c^1_{Qu} c^8_{td}			11.0								12.4		0.0				27.7								0.1						_	ł
c_{td}^{o}			14.4 13.8	0.3		0.4					9.7 9.6		0.0	0.0	0.2		29.1 38.8		2.0				42.8 37.1		0.2	0.1						
c_{td}^{1}			8.7	0.2		2.4					9.4			0.0	0.5		21.2		12.1				42.9		0.8	0.2						
c_{Qd}^8			13.8	0.2		2.4					10.2		0.0	0.0	0.5		35.6		12.1				40.0		0.0	0.2						
c_{Qd}^1		0.4	13.0								10.2		0.0	0.0	0.0		33.0						40.0		0.1	0.1				70.0	21.1	
$c_{c\varphi}$														0.0	0.0											0.1				-	29.1	- 60
$c_{b\varphi}$														0.5	3.9											16.9				53.6		
$c_{t\varphi}$														0.0	0.1											0.0					21.2	
$c_{\tau\varphi}$		1.8	1.3	0.1	0.0	01			0.0		0.0	0.0	0.1		9.1		75	0.1	0.9		0.0		0.0	0.0	0.4	39.9				25.4		
c_{tG}			1.0	0.0	0.0	0.0	0.0	0.0		0.0	0.0	1.9	0.11	2.3	12.5		1.0	0.1	0.0	0.1		0.0	0.0	4.1		41.8				26.1		
c_{tW}				0.0		0.0	0.0	0.0	0.0	0.0				2.5	13.3				0.0	•		0.0				44.6				27.9		ł
C_{tZ} (3)	3.2				0.0	0.0	0.0	0.0		0.0				0.0	0.1	0.0		0.0		0.0		0.0				1.8	0.5	84.8	3.4	3.5		
$c_{\varphi q}$ (3)	1.8						0.0		0.0	0.0				0.0	0.0	0.0			0.0	0.0	0.0	0.0					0.0			0.0	-	
$c_{\varphi Q}$	1.5					0.0				0.0				0.0	0.0	0.0			0.0			0.0					0.0	82.2		14.5		
$c^{(3)}_{arphi q q} \ c^{(3)}_{arphi Q} \ c^{(3)}_{arphi Q} \ c^{(-)}_{arphi q} \ c^{(-)}_{arphi Q}$	1.5					0.0				0.0						0.0			0.0			0.0				0.0		80.7		16.1		
	3.8					0.0								0.0	0.1	0.0			0.0							1.1		95.1		0.0		- 40
$c_{arphi u} \ c_{arphi d}$	4.5					0.0								0.0		0.0			0.0								0.0	95.2		0.0		
$c_{\varphi t}$						11.2				0.1				0.3	1.8				74.8			0.5				6.2				3.6	1.5	
$c_{\varphi l_1}$	1.6													0.0	0.0	0.0											0.0	42.5	0.0	28.7		
$c_{\varphi l_2}$	4.6													0.0	0.0													78.1		15.6		
$c_{\varphi l_3}$	3.1													0.0	0.0	0.0										0.0	0.0	81.4		13.9	1.5	ł
c ⁽³⁾	0.1			0.0	0.0	0.0	0.0	0.0	0.0	0.0				0.0	0.0	0.0		0.0	0.0	0.0	0.0	0.0				0.0	0.0	3.1	4.2	79.6	12.9	
$c^{(3)}_{arphi l_1} \ c^{(3)}_{arphi l_2} \ c^{(3)}_{arphi l_2} \ c^{(3)}_{arphi l_3}$	0.1			0.0	0.0	0.0	0.0	0.0	0.0	0.0				0.0	0.0	0.0		0.0	0.0	0.0	0.0	0.0				0.0	0.0	1.1	5.1	82.5	11.2	
$c^{(3)}$	2.4													0.0	0.0	0.0										0.0	0.0	68.5	6.7	16.2	6.3	
$c_{\varphi l_3}$ $c_{\varphi e}$	1.5													0.0	0.0	0.0										0.0	0.0	31.0	0.0	41.5	25.9	
$c_{\varphi\mu}$	4.3													0.0	0.0	0.0										0.0	0.0	78.6		15.4	1.7	- 20
$c_{\varphi\tau}$	3.5													0.0	0.0	0.0										0.0	0.0	81.7		13.3	1.5	
c_{ll}	0.0			0.0	0.0	0.0	0.0	0.0	0.0	0.0				0.0	0.0	0.0		0.0	0.0	0.0	0.0	0.0				0.0	0.0	0.1	2.5	52.9	44.5	
$c_{\varphi G}$														0.3	2.5											10.9				58.7	27.6	
$c_{\varphi B}$														2.5	13.2											44.1				28.6	11.7	
$c_{\varphi W}$	1													1.1	5.8											19.4				46.4	27.3	Ì
$c_{\varphi WB}$	0.0			0.0		0.0				0.0				0.0	0.0	0.0			0.0			0.0				0.1	0.0	0.0	0.0	88.6	11.1	
CWWW	0.2									0.0						0.1						0.0					4.8		0.0	63.4	31.4	
$c_{\varphi \Box}$														0.0	0.1											0.2				75.2	24.5	
$c_{\varphi D}$	0.1			0.0		0.0				0.0				0.0	0.0	0.0			0.0			0.0				0.0	0.0	0.1	0.0	88.8	11.0	
72							1		1	-		-		_										_	_				_			± 0

Various statistical measures available to determine which datasets dominate the reach for specific EFT directions

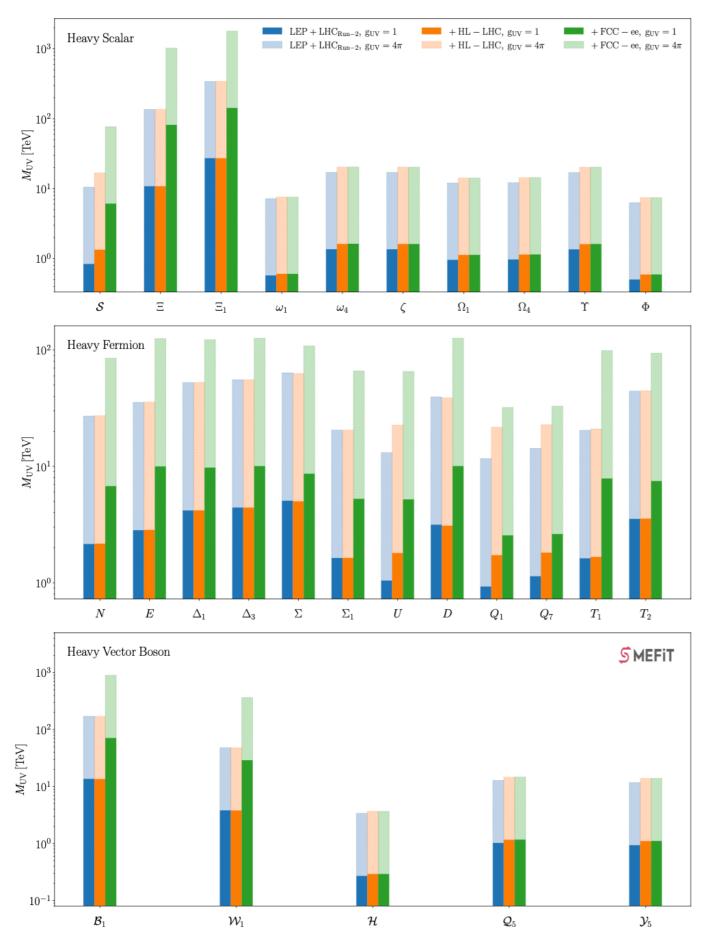
FCC-ee fully dominates the constraints on all operators except the two-light-two-heavy ones

In terms of the SMEFT, the run at 161 GeV is the one providing less information

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Normalized Value

Matching to UV models

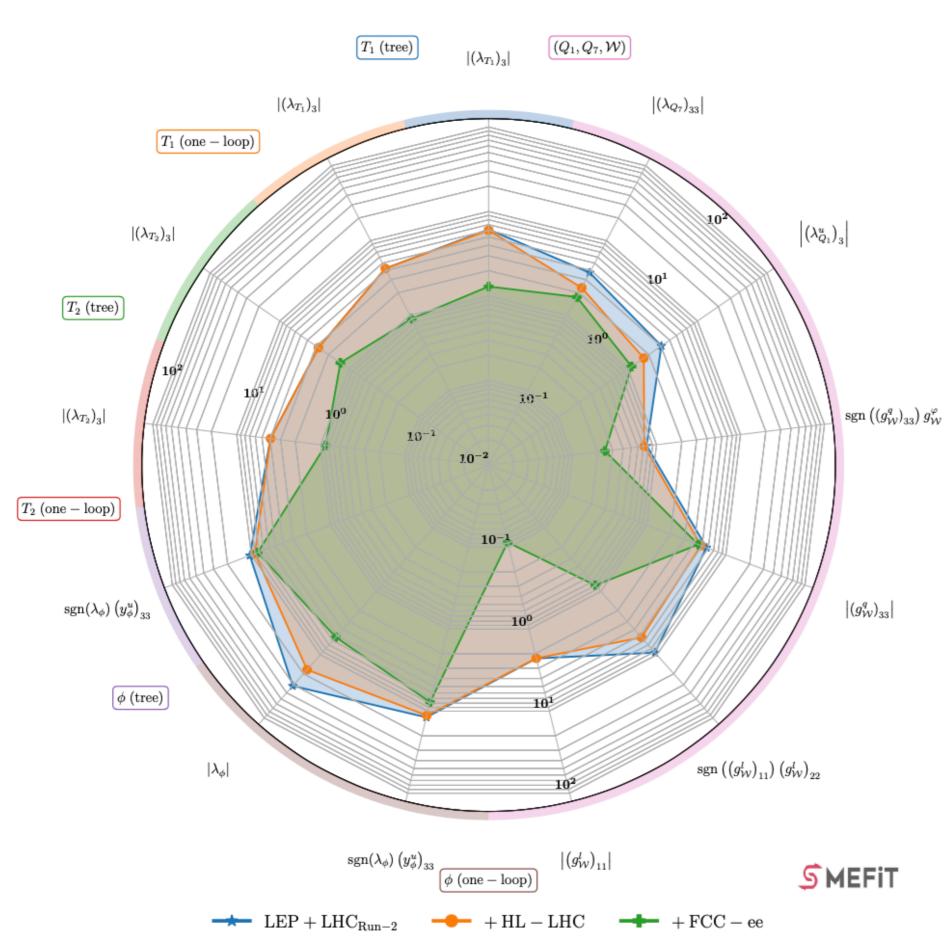


Lower bounds (95% CL) on different oneparticle extensions of the SM matched to the SMEFT at tree level

FCC-ee has an (indirect) reach on heavy particles with masses between a few TeV and up to around 100 TeV, for O(1) UV couplings

Strongest impact for UV models that induce the purely bosonic and two-fermion operators, which are tightly constrained by FCC-ee measurements

Matching to UV models



Analysis can be extended to
 multi-particle models as
 well as to one-loop
 matching to the SMEFT

Fully automated pipeline,
 determining constraints from
 FCC-ee data on general UV
 models matched to the
 SMEFT is now streamlined

Summary and outlook

- The SMEFiT framework enables detailed projection studies quantifying the impact of future colliders in the parameter space of the SMEFT and of UV-complete models matched upon it
- Starting point is a state-of-the-art global EFT fit, extended first with HL-LHC and then FCC-ee projections
- All results shown are fully reproducible by means of our **open-source**, **user-friendly code**
- Results for **CEPC** are in general similar to those of the FCC, with some noteworthy differences
- Projections for HL-LHC can be refined with the used of optimised observables (better coverage of highmass region, multi-differential observables, benefitting from increased statistics)
- In the pipeline is incorporating projections for other proposed future colliders (ILC, CLIC, C3, muon collider), more extensive investigations of the impact on UV complete models, and making more general the operator basis (flavour assumptions) defining our baseline analysis

Interested in projections for future colliders? **Try SMEFiT**, and let us know if any required functionality is missing!