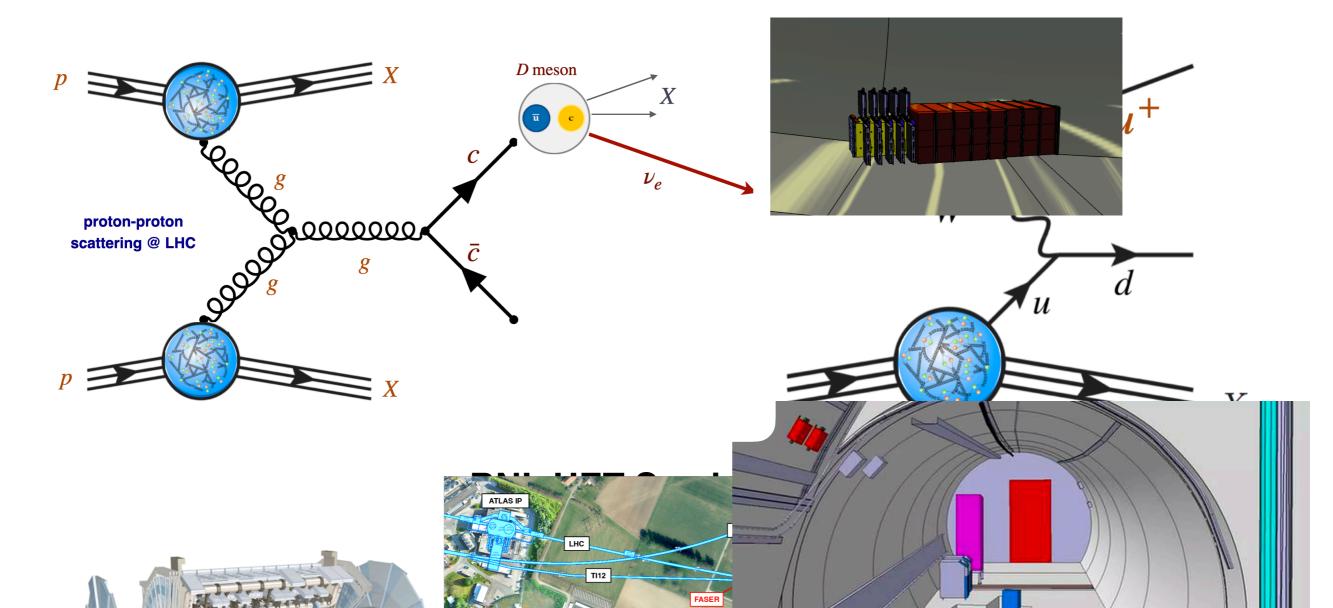


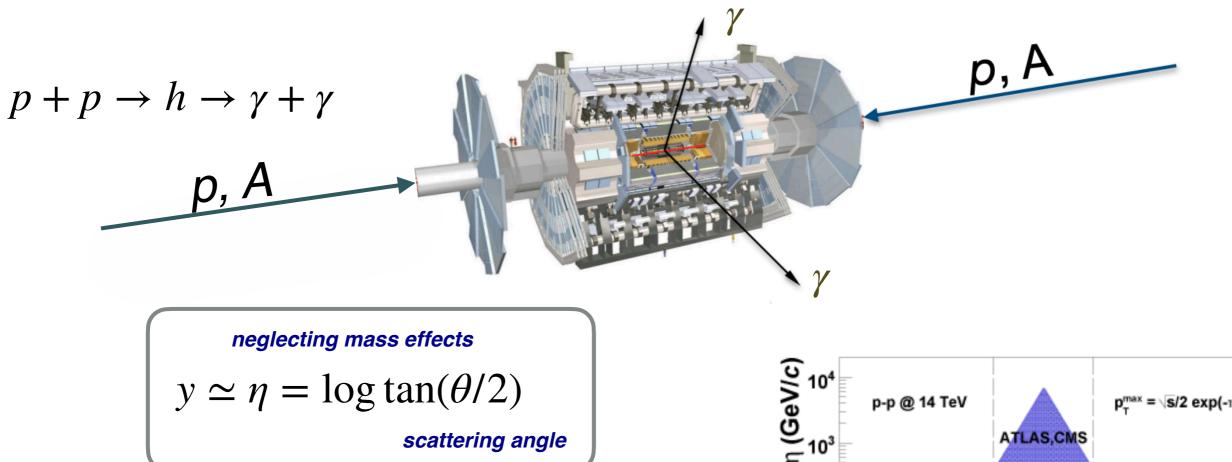


# **Physics with TeV Neutrinos at the LHC**

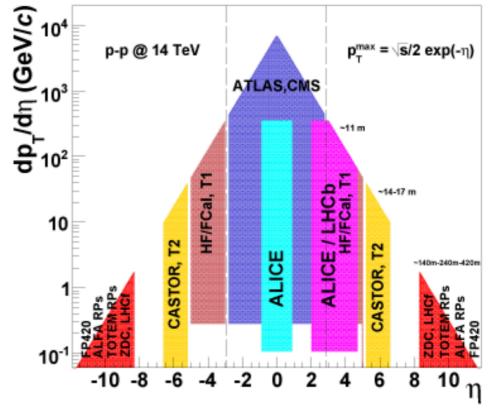
Juan Rojo, VU Amsterdam & Nikhef



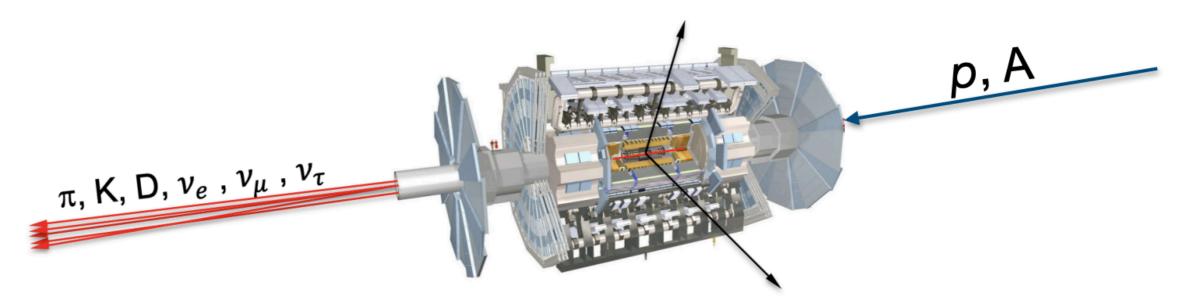
The ATLAS and CMS detectors were designed with a focus on identifying weak-scale and heavier particles, whose decay products lie in the central rapidity acceptance region



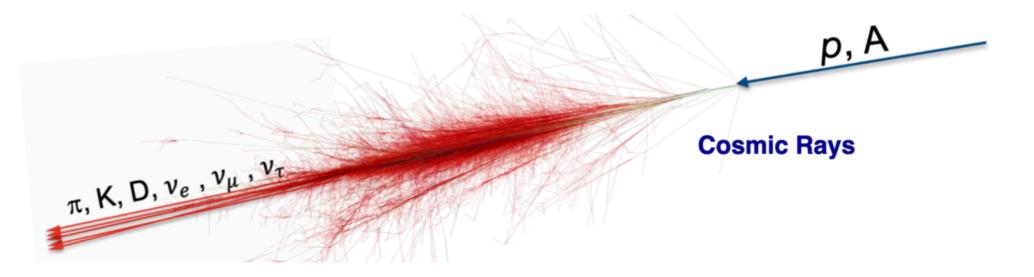
- Solution Forward region for hard-scattering physics restricted to **LHCb** ( $\eta < 4.5$ ) and in the future **ALICE-FoCal** ( $\eta < 5.0$ )
- Far-forward region essentially beyond access for current LHC detectors, except for *e.g.* total cross-section analyses



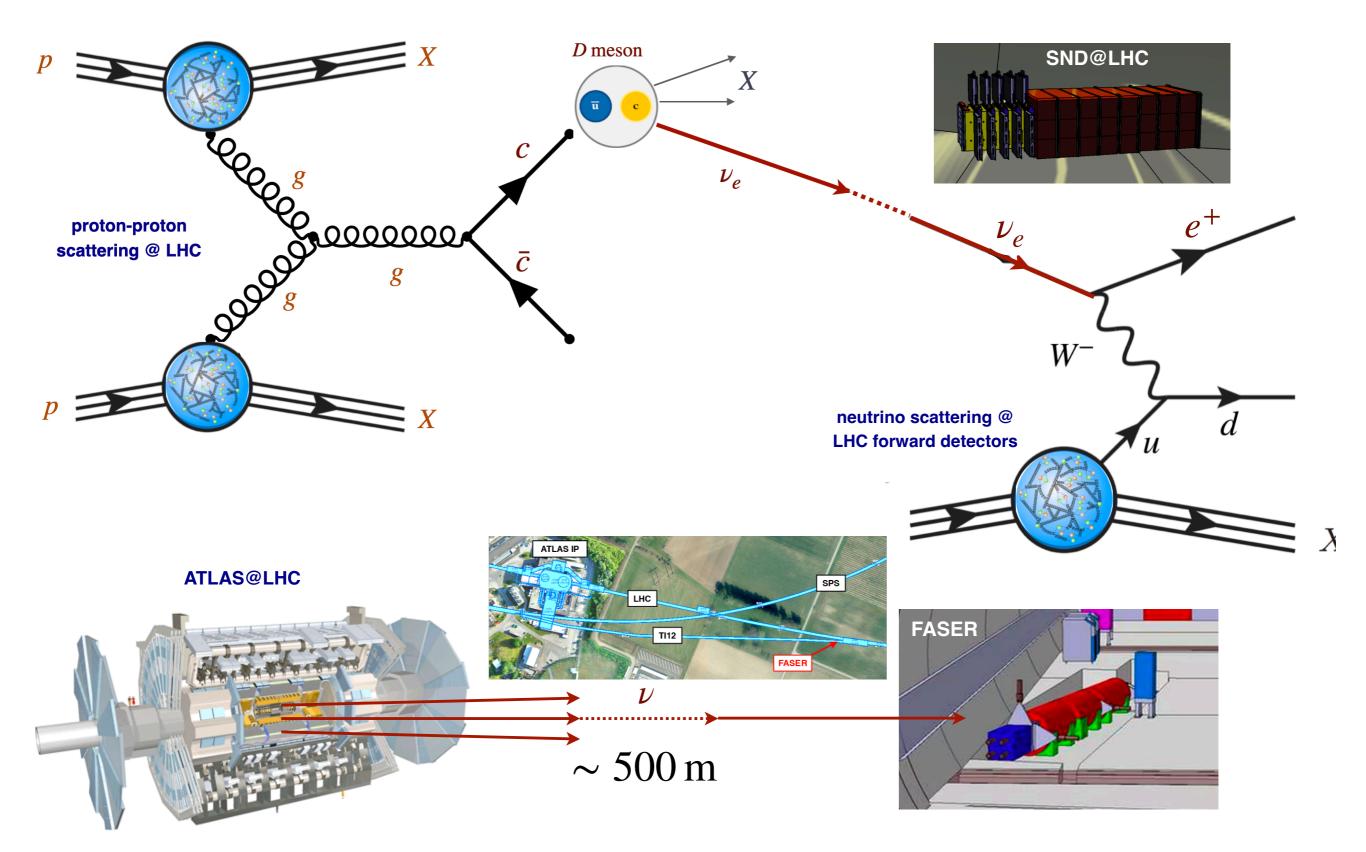
LHC collisions result into a large flux of energetic neutrinos which escape the detectors unobserved: major blind spot of the LHC



Being able to detect and utilise the most energetic human-made neutrinos ever produced would open many exciting avenues in QCD, neutrino, and astroparticle physics

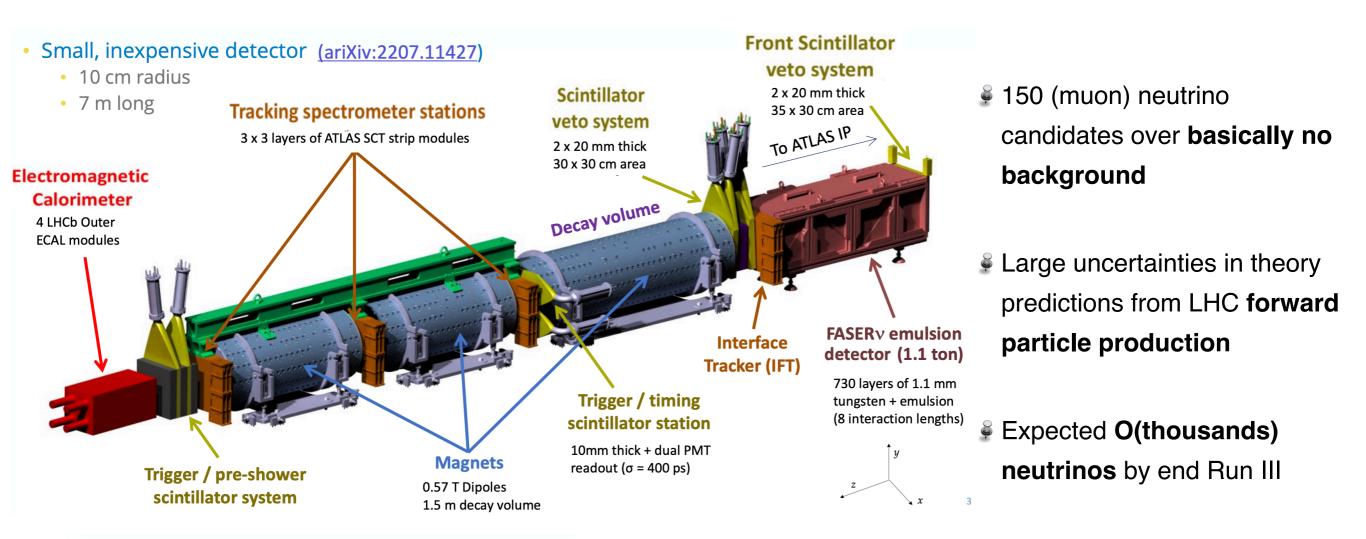


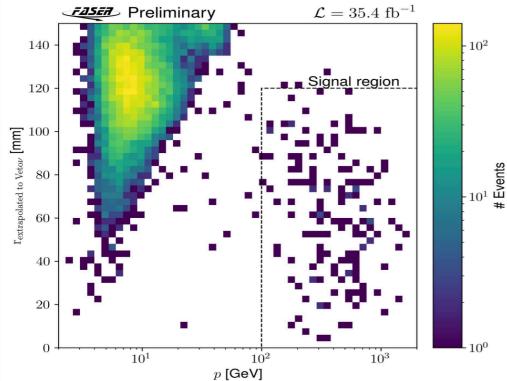
LHC neutrinos (CoM frame)  $\approx$  Cosmic Rays and UHE neutrinos (lab frame)



FASER & SND@LHC demonstrated how far detectors can be deployed to identify collider neutrinos

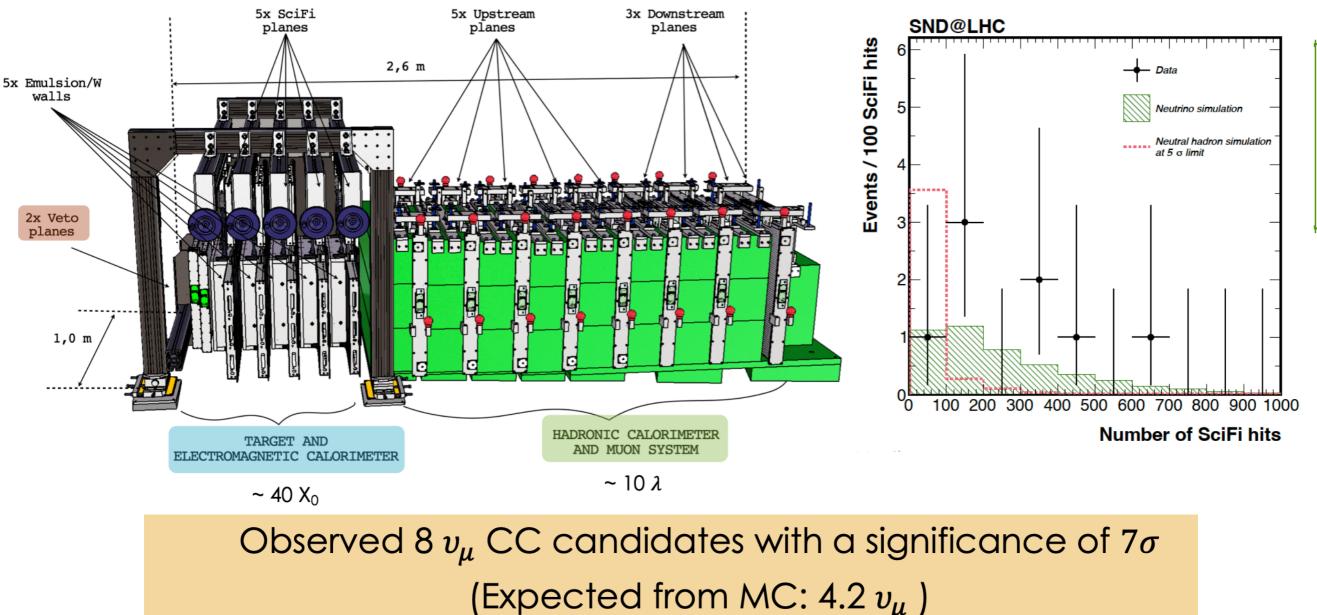
### Neutrinos at the LHC: FASER





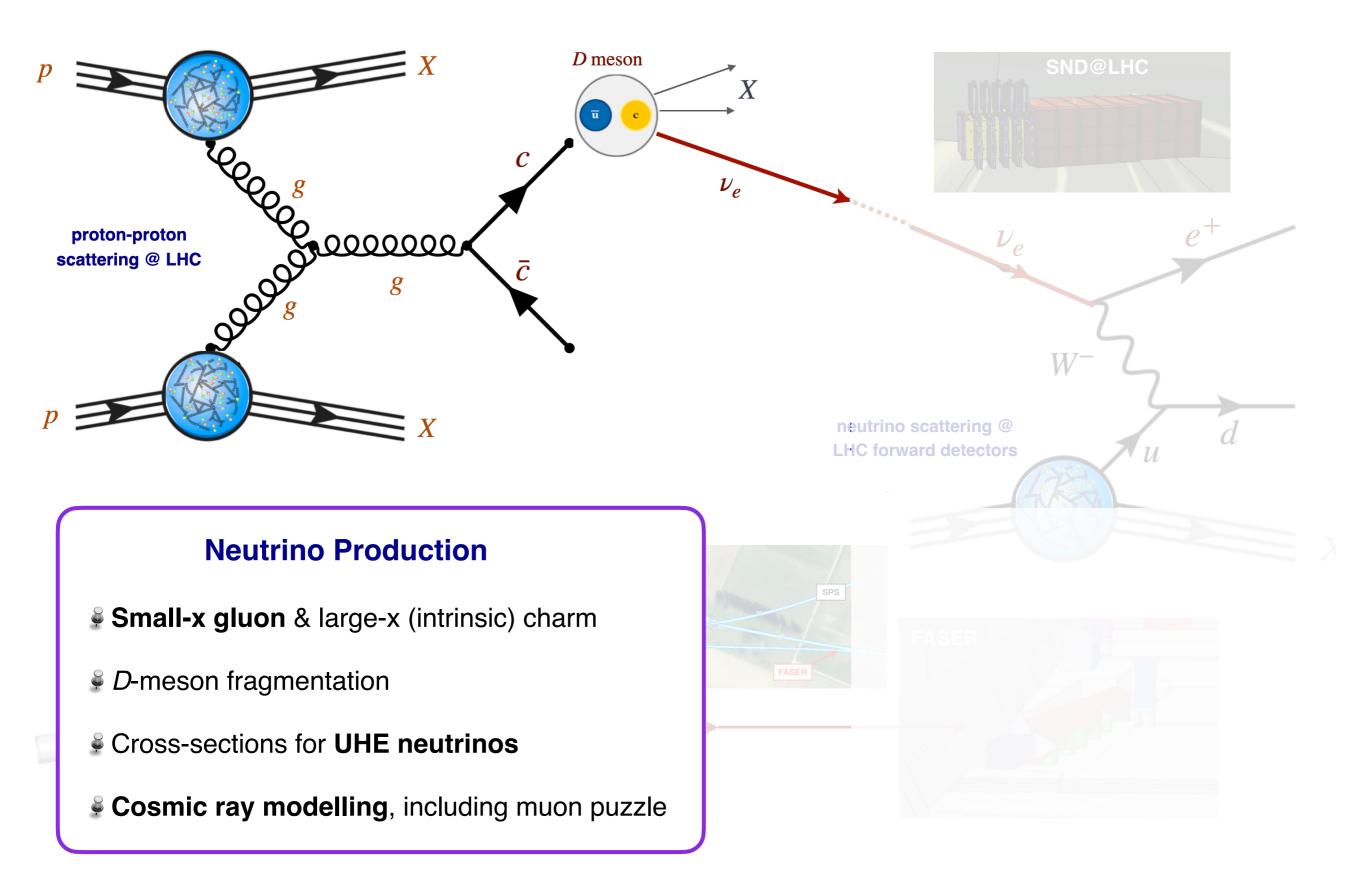
Candidate	Events
n <sub>o</sub>	<b>153</b> (151 ± 41)
<b>n</b> <sub>10</sub>	4
<b>n</b> <sub>01</sub>	6
n <sub>2</sub>	64014695

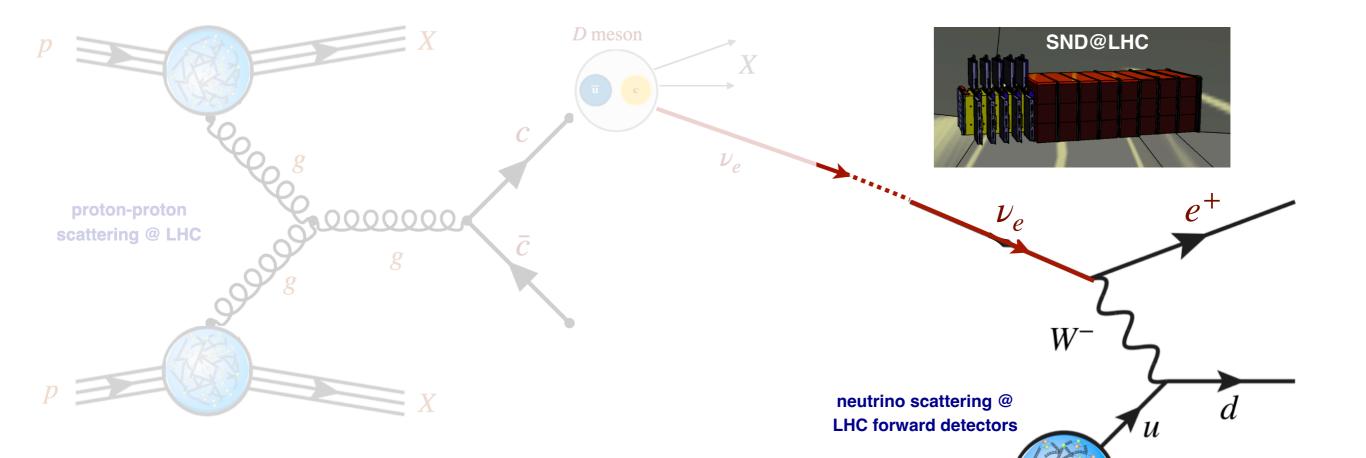
#### Neutrinos at the LHC: SND@LHC



#### 8 neutrino candidates (likely muon neutrinos)

Different rapidity range from FASER (slightly off-axis): SND@LHC covers 7.2 < y<sub>v</sub> < 8.4</p>



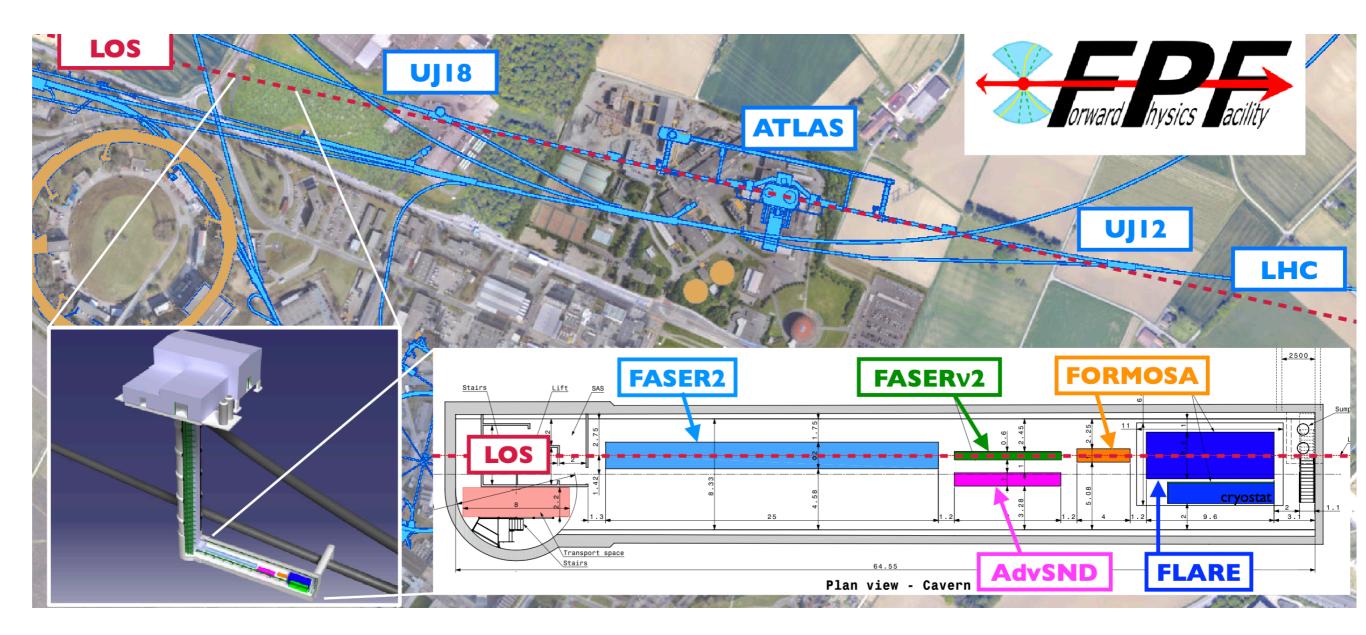


FASER

#### **Neutrino Scattering**

- DIS with TeV neutrinos (``Neutrino-Ion Collider')
- Neutrino (EFT) interactions at the TeV
- Cross-sections for atmospheric neutrinos
- Nuclear PDFs, strangeness from charm
- Neutrino flavor (non-)university (with tau neutrinos)

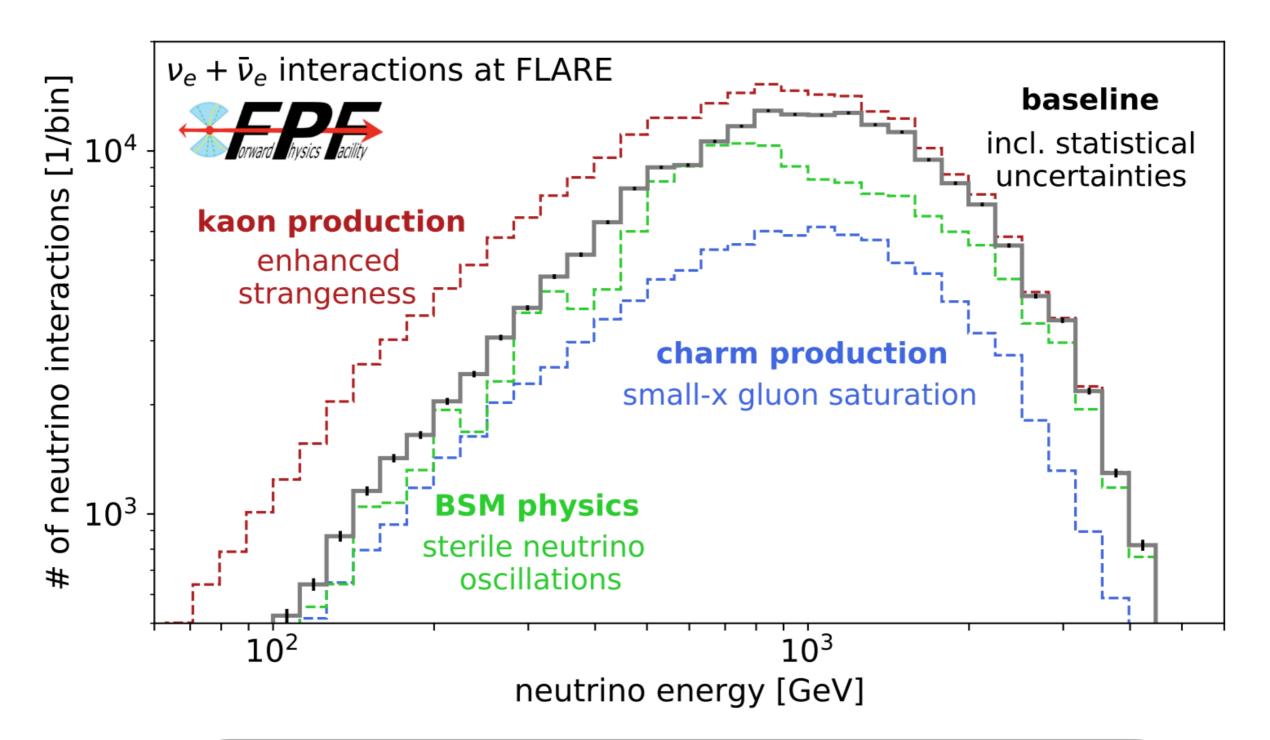
# **Forward Physics Facility**



Dedicated new cavern equipped with a suite of far-forward experiments

- Operating concurrently with the HL-LHC, exploit intense high-energy beam of forward particles
- Start civil engineering during LS3 or shortly thereafter: positive outcome of ongoing site investigation studies (including drill down to the cavern depth)

# **Forward Physics Facility**

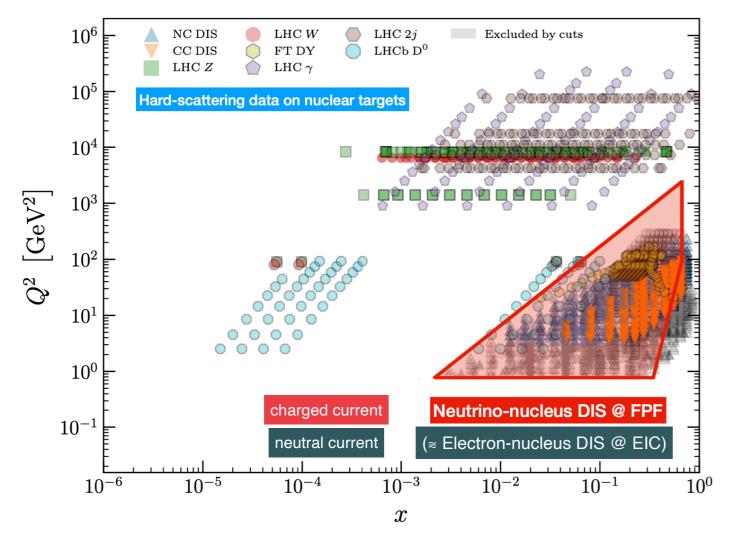


Broad, far-reaching program on **QCD** (small-x gluon, saturation), **cosmic rays** (muon puzzle), **neutrino BSM** (sterile neutrinos), hadronic structure, **UHE neutrinos**, **FCC-pp cross-sections** ...

# Proton and Nuclear Structure with TeV Neutrinos

M. Fieg, T. Giani, P. Krack, G. Magni, T. Makela, T. Rabemananjara, J. Rojo, *paper in preparation* 

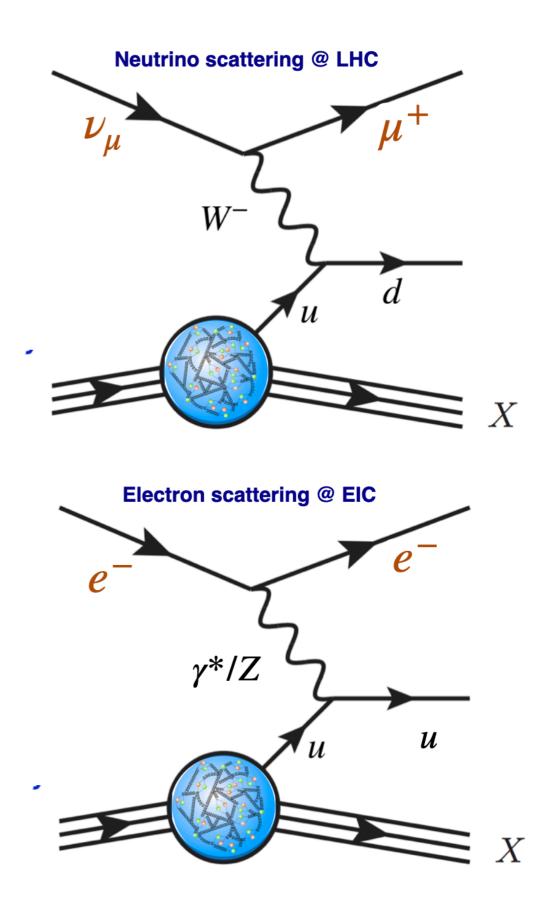
### Neutrino DIS at the LHC





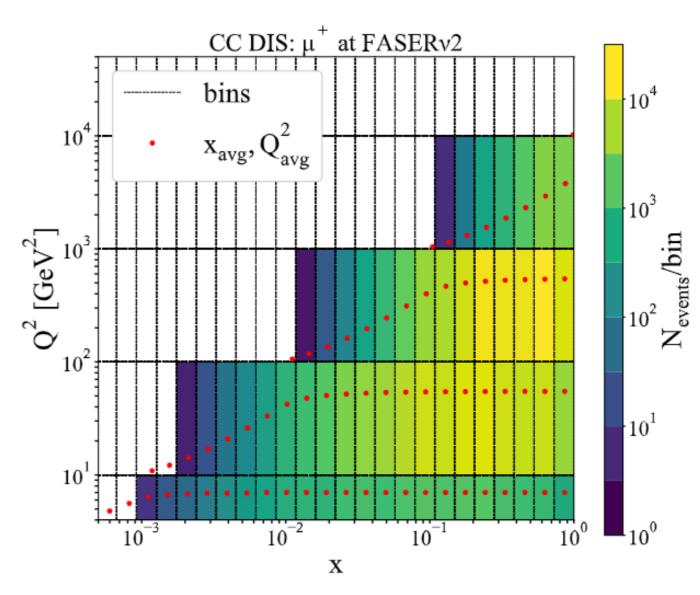
- Continue succesful program of neutrino DIS experiments
   CERN & expand kinematic coverage
- Charged-current analog of the Electron-Ion Collider: the LHC as a Neutrino-Ion Collider

Constrain proton & nuclear light (anti-)quark PDFs



# Impact projections

- Neutrino fluxes from Kling-Nevay calculation
- Focus on muon neutrinos: higher rates, dominated by light hadron production
- Generate pseudo-data for DIS structure functions for FASER, SND@LHC, and the proposed FPF experiments, both inclusive and charm production
- Assume outgoing lepton charge separation
- Model systematic errors based on the feedback provided by the experiments



$$N_{\rm ev}/{\rm bin} = n_T L_T \int_{Q^2_{\rm min}}^{Q^2_{\rm max}} \int_{x_{\rm min}}^{x_{\rm max}} \int_{E_v^{\rm min}}^{E_v^{\rm max}} \frac{dN_v(E_v)}{dE_v} \frac{d^2 \sigma^{vA}(x, y, E_v)}{dx dy} dQ^2 dx dE_v$$
Geometry/Target
Binning
neutrino fluxes
(include rapidity  
acceptance)
DIS differential  
cross-section

#### Both for inclusive production and for charm-tagged final states

# Impact projections

- Neutrino fluxes from Kling-Nevay calculation
- Focus on muon neutrinos: higher rates, dominated by light hadron production
- Generate pseudo-data for DIS structure functions for FASER, SND@LHC, and the proposed FPF experiments, both inclusive and charm production

$$\begin{array}{lll} F_2^{\nu p}(x,Q^2) &=& 2x \left(f_{\bar{u}}+f_d+f_s+f_{\bar{c}}\right) \left(x,Q^2\right), \\ F_2^{\bar{\nu} p}(x,Q^2) &=& 2x \left(f_u+f_{\bar{d}}+f_{\bar{s}}+f_c\right) \left(x,Q^2\right), \\ xF_3^{\nu p}(x,Q^2) &=& 2x \left(-f_{\bar{u}}+f_d+f_s-f_{\bar{c}}\right) \left(x,Q^2\right), \\ xF_3^{\bar{\nu} p}(x,Q^2) &=& 2x \left(f_u-f_{\bar{d}}-f_{\bar{s}}+f_c\right) \left(x,Q^2\right), \end{array}$$

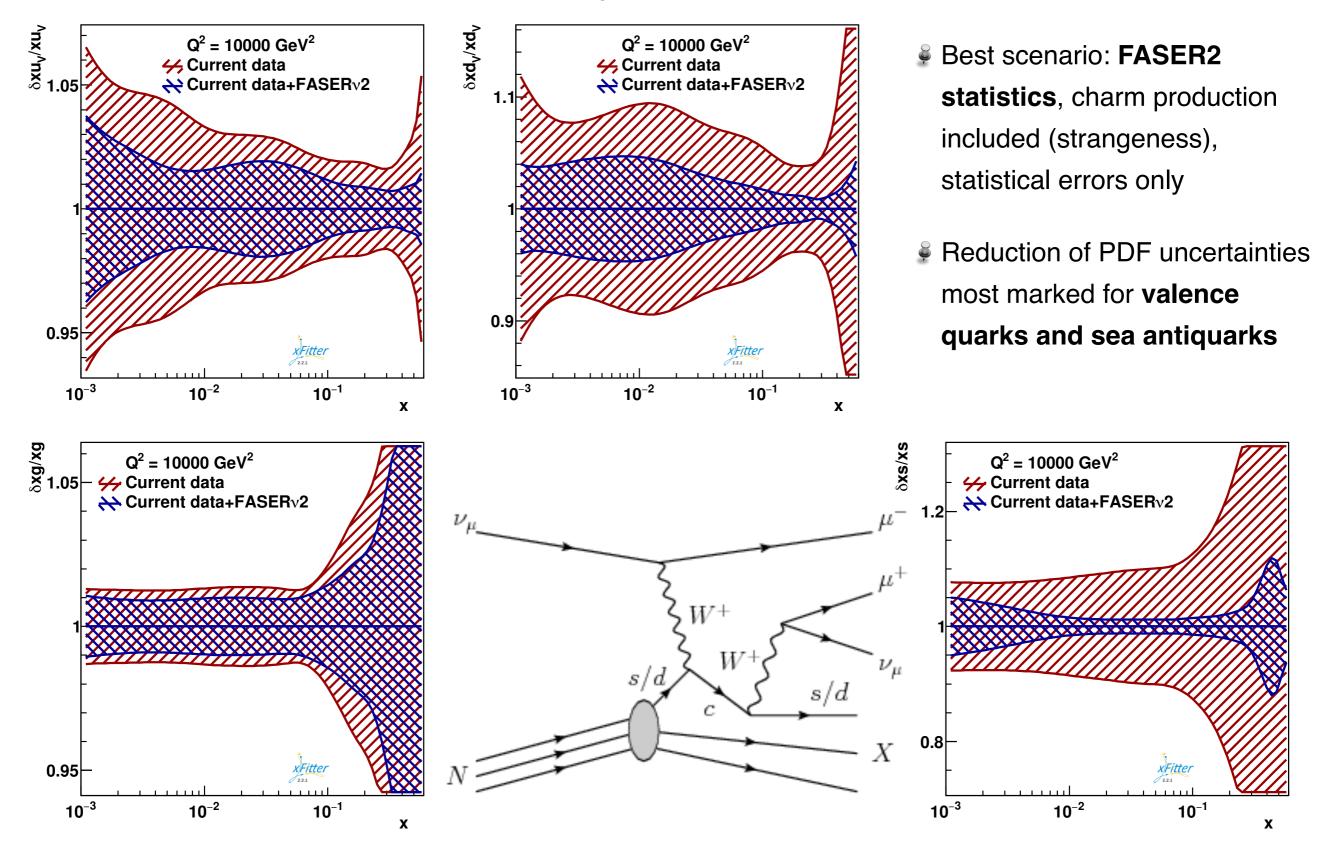
- Assume outgoing lepton charge separation
- Model systematic errors based on the feedback provided by the experiments

$$\frac{d^2 \sigma^{\nu A}(x,Q^2,y)}{dxdy} = \frac{G_F^2 s/4\pi}{\left(1+Q^2/m_W^2\right)^2} \left[Y_+ F_2^{\nu A}(x,Q^2) - y^2 F_L^{\nu A}(x,Q^2) + Y_- x F_3^{\nu A}(x,Q^2)\right]$$
$$\frac{d^2 \sigma^{\bar{\nu}A}(x,Q^2,y)}{dxdy} = \frac{G_F^2 s/4\pi}{\left(1+Q^2/m_W^2\right)^2} \left[Y_+ F_2^{\bar{\nu}A}(x,Q^2) - y^2 F_L^{\bar{\nu}A}(x,Q^2) - Y_- x F_3^{\bar{\nu}A}(x,Q^2)\right]$$

Differential measurements with charge-separation key to achieve sensitivity to proton and nuclear structure

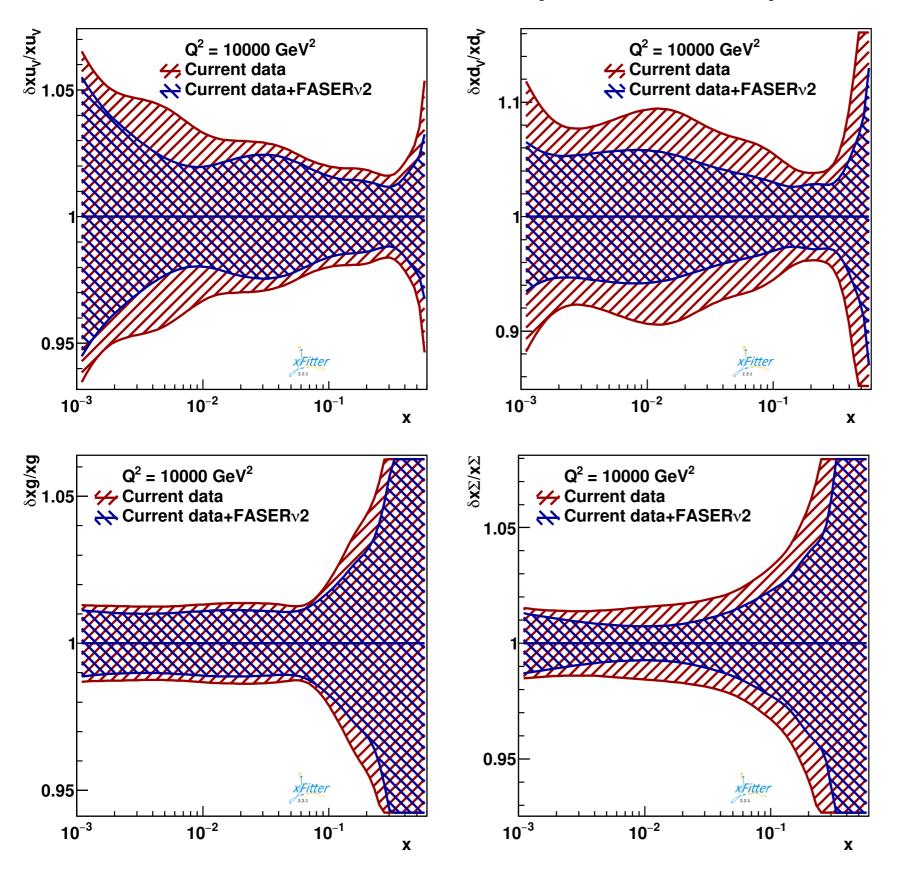
### **Results: proton PDFs**

Statistical error only, inclusive + charm data

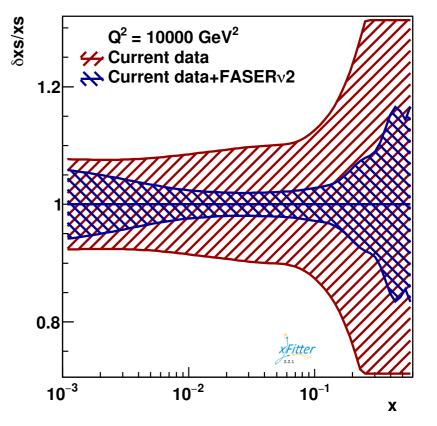


### **Results: proton PDFs**

Statistical + Systematic errors only, inclusive + charm data

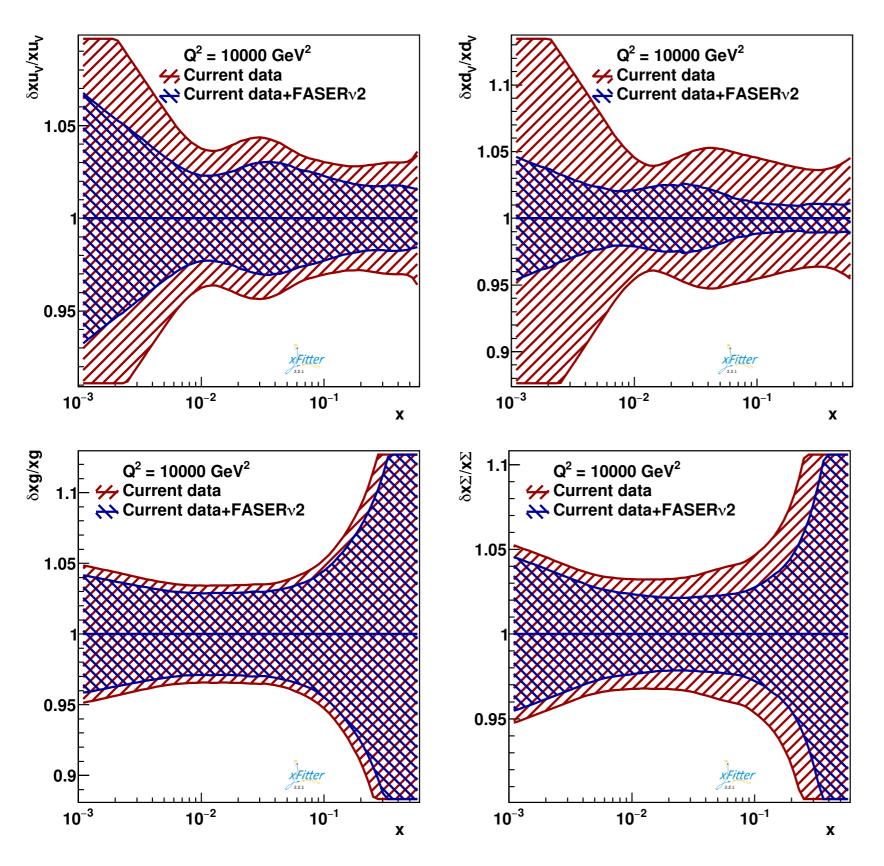


- Results are robust upon inclusion of systematic errors
- Depends on assumptions on
   correlation model, in
   particular bin-by-bin
   correlations
- Study of different scenarios in progress

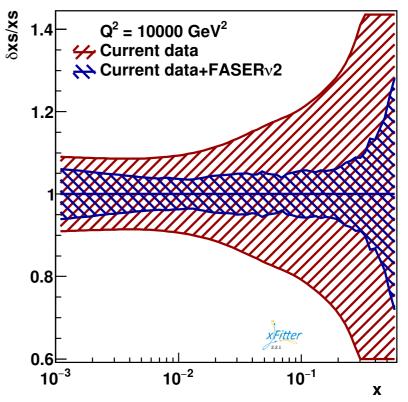


# Results: nuclear PDFs

#### Statistical error only, inclusive + charm data

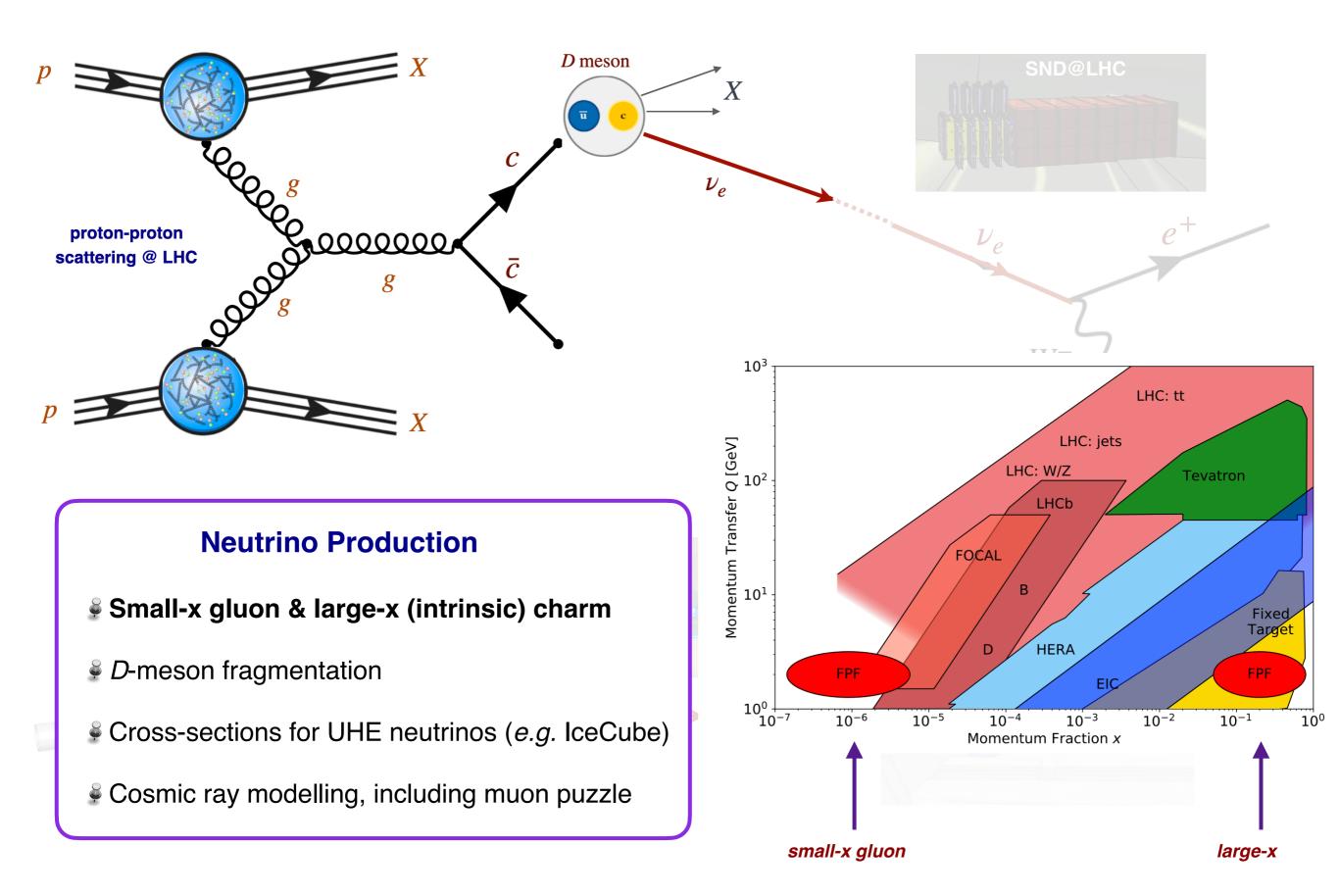


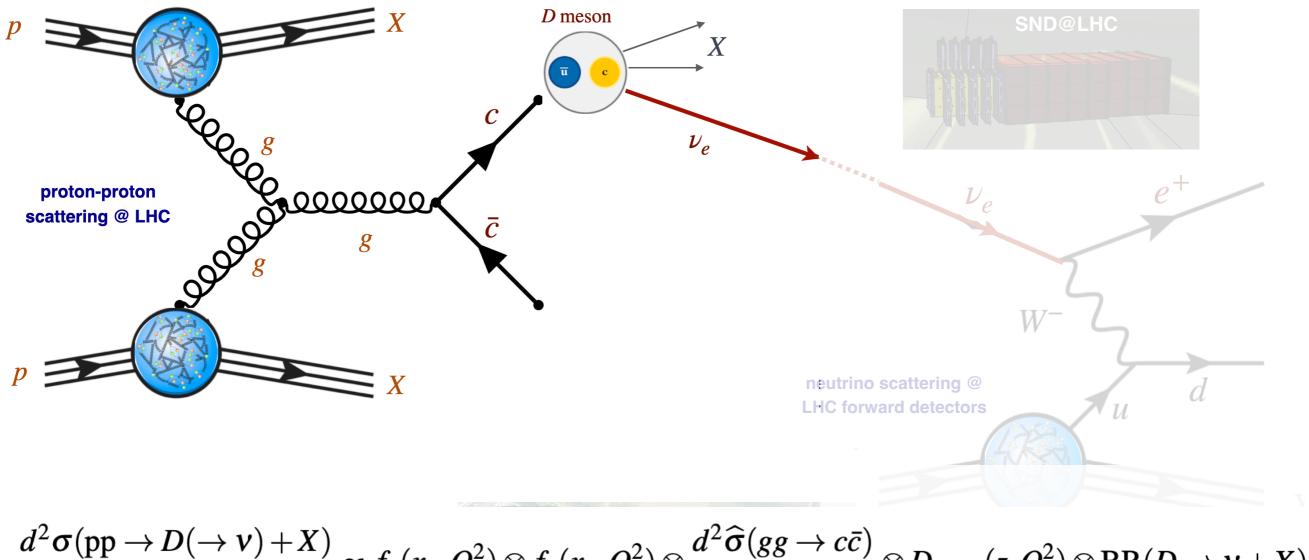
- Profiling of the EPPS21 global nPDF fit (Tungsten) reveals a consistent picture
- Excellent sensitivity to quarkflavour separation &strangeness
- Ideally, use different nuclear targets in the detector



# Pinning down ultra-small-x gluon with LHC neutrinos

P. Krack, S. Niedenzu, J. Rojo, J. Sola-Cava, *work in progress* 



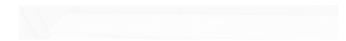


 $\frac{d^2\sigma(\mathrm{pp}\to D(\to\nu)+X)}{p_T^{\nu}y_{\nu}} \propto f_g(x_1,Q^2) \otimes f_g(x_2,Q^2) \otimes \frac{d^2\widehat{\sigma}(gg\to c\bar{c})}{p_T^{c}y_{c}} \otimes D_{c\to D}(z,Q^2) \otimes \mathrm{BR}(D\to\nu+X)$ 

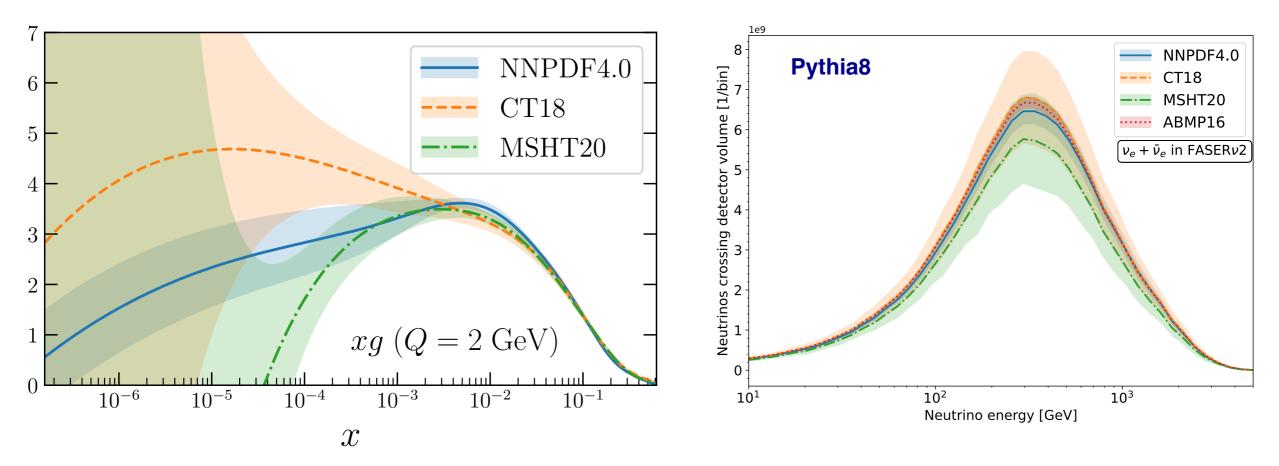
Extract from measured neutrino fluxes

Constrain from FASER/FPF data

QCD prediction: NLO + PS large theory uncertainties



# Impact projections



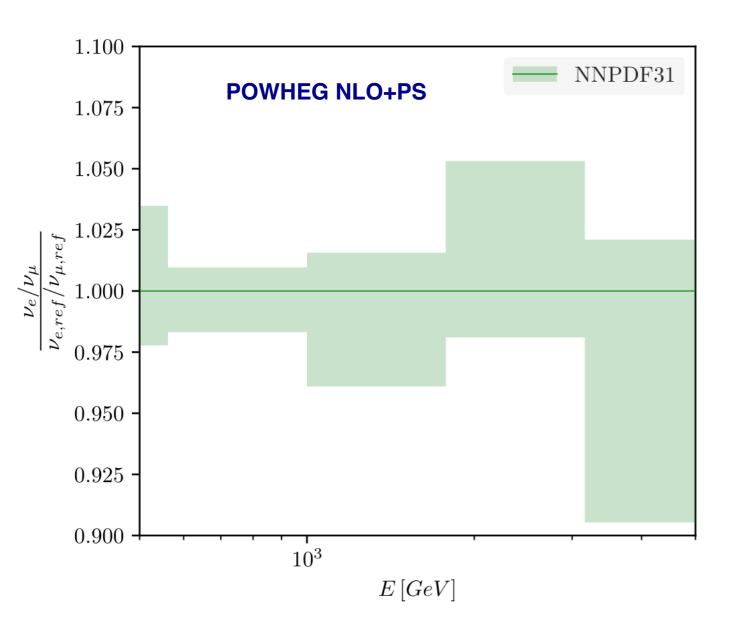
Spread of PDF predictions (e.g. small-x gluon) modifies predicted fluxes up to factor 2

- Focus on electron and tau neutrinos, with the largest contribution from charm production where QCD factorisation can be applied
- Seconstruct tailored observables where QCD uncertainties (partially) cancel out

$$R_{\tau/e}(E_{\nu}) \equiv \frac{N(\nu_{\tau} + \bar{\nu}_{\tau}; E_{\nu})}{N(\nu_{e} + \bar{\nu}_{e}; E_{\nu})}, \qquad R_{\exp}^{\nu_{e}}(E_{\nu}) = \frac{N_{\text{FASER}\nu}(\nu_{e} + \bar{\nu}_{e}E_{\nu})}{N_{\text{SND}@LHC}(\nu_{e} + \bar{\nu}_{e}; E_{\nu})}$$

#### Retain PDF sensitivity while reducing the large QCD uncertainties in the theory prediction

# Impact projections

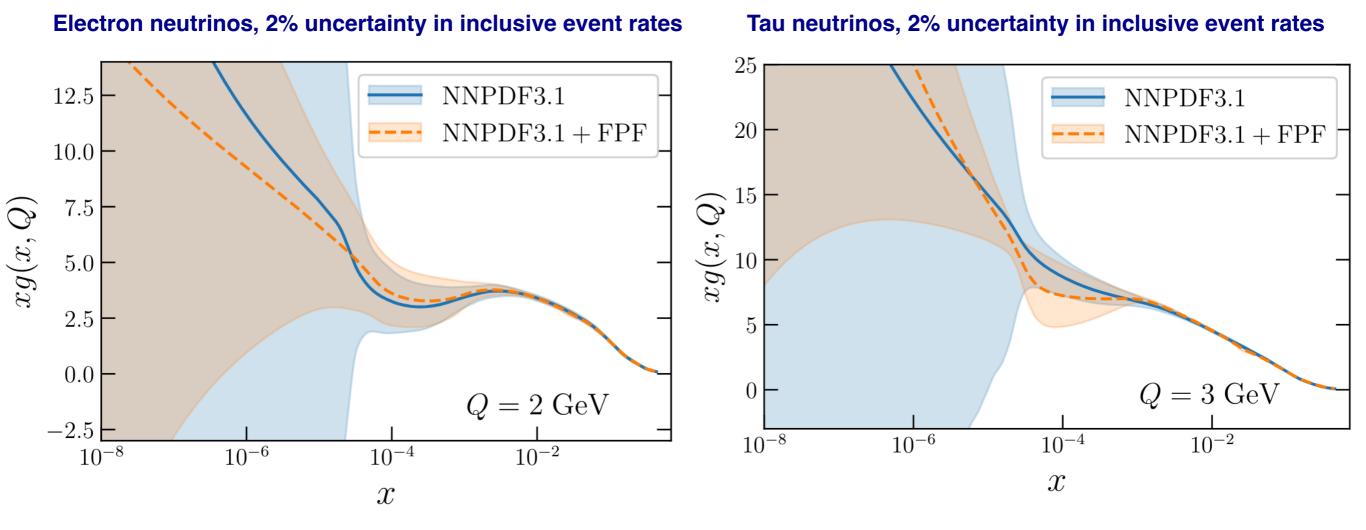


- When taking ratios of event rates (e.g. charm electron neutrinos vs charm muon neutrinos), QCD uncertainties reduced to O(few %)
- Strategy: assume a measurement of inclusive event rates as a function of neutrino energy with a given precision, quantify impact on PDFs via Bayesian reweighting

$$R_{\tau/e}(E_{\nu}) \equiv \frac{N(\nu_{\tau} + \bar{\nu}_{\tau}; E_{\nu})}{N(\nu_{e} + \bar{\nu}_{e}; E_{\nu})}, \qquad R_{\exp}^{\nu_{e}}(E_{\nu}) = \frac{N_{\text{FASER}\nu}(\nu_{e} + \bar{\nu}_{e}E_{\nu})}{N_{\text{SND@LHC}}(\nu_{e} + \bar{\nu}_{e}; E_{\nu})}$$

#### Retain PDF sensitivity while reducing the large QCD uncertainties in the theory prediction

#### Results



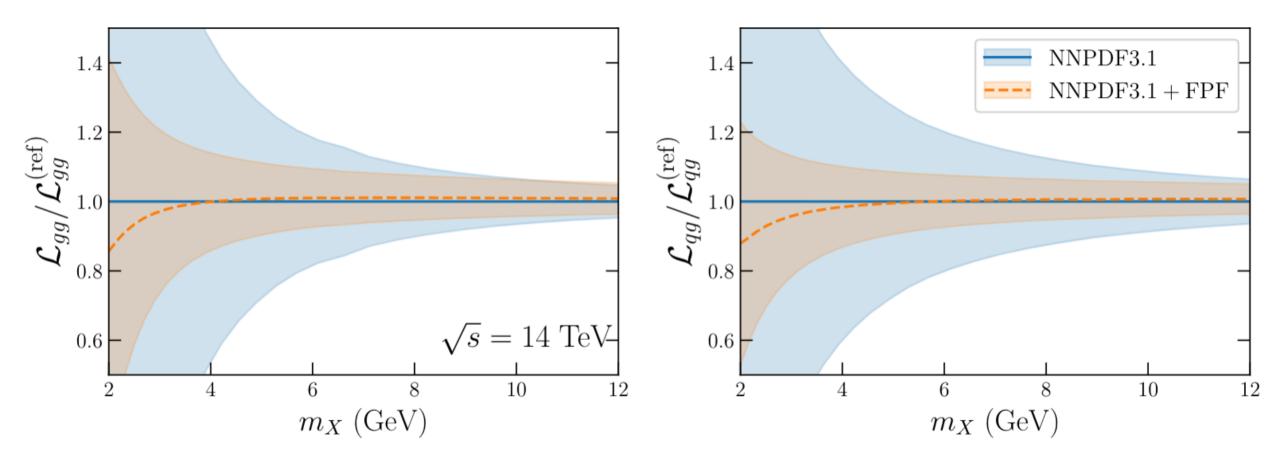
Results based on pseudo-data for a **measurement of the rapidity ratio** (proxy for experiment ratio)

$$R_{y}^{(e)} \equiv \frac{N_{\nu_{e}}(E_{\nu}, 7.5 < y_{u} < 8.0)}{N_{\nu_{e}}(E_{\nu}, 8.5 < y_{u} < 9.0)} \qquad \qquad R_{y}^{(\tau)} \equiv \frac{N_{\nu_{\tau}}(E_{\nu}, 7.5 < y_{u} < 8.0)}{N_{\nu_{\tau}}(E_{\nu}, 8.5 < y_{u} < 9.0)}$$

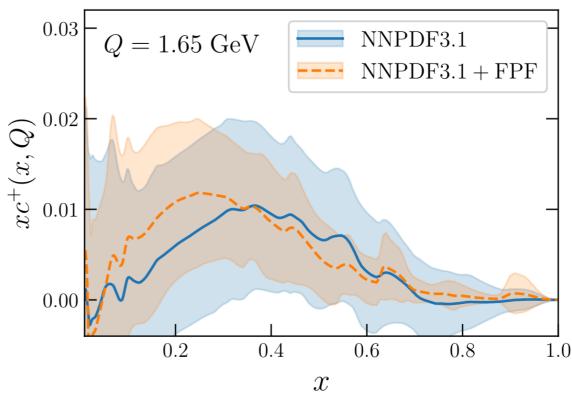
Sensitivity to small-x gluon outside coverage of any other (laboratory) experiment
 Study impact of different observables, QCD errors, and the precision of measurement

#### **Results**

Electron neutrinos, 2% uncertainty in inclusive event rates

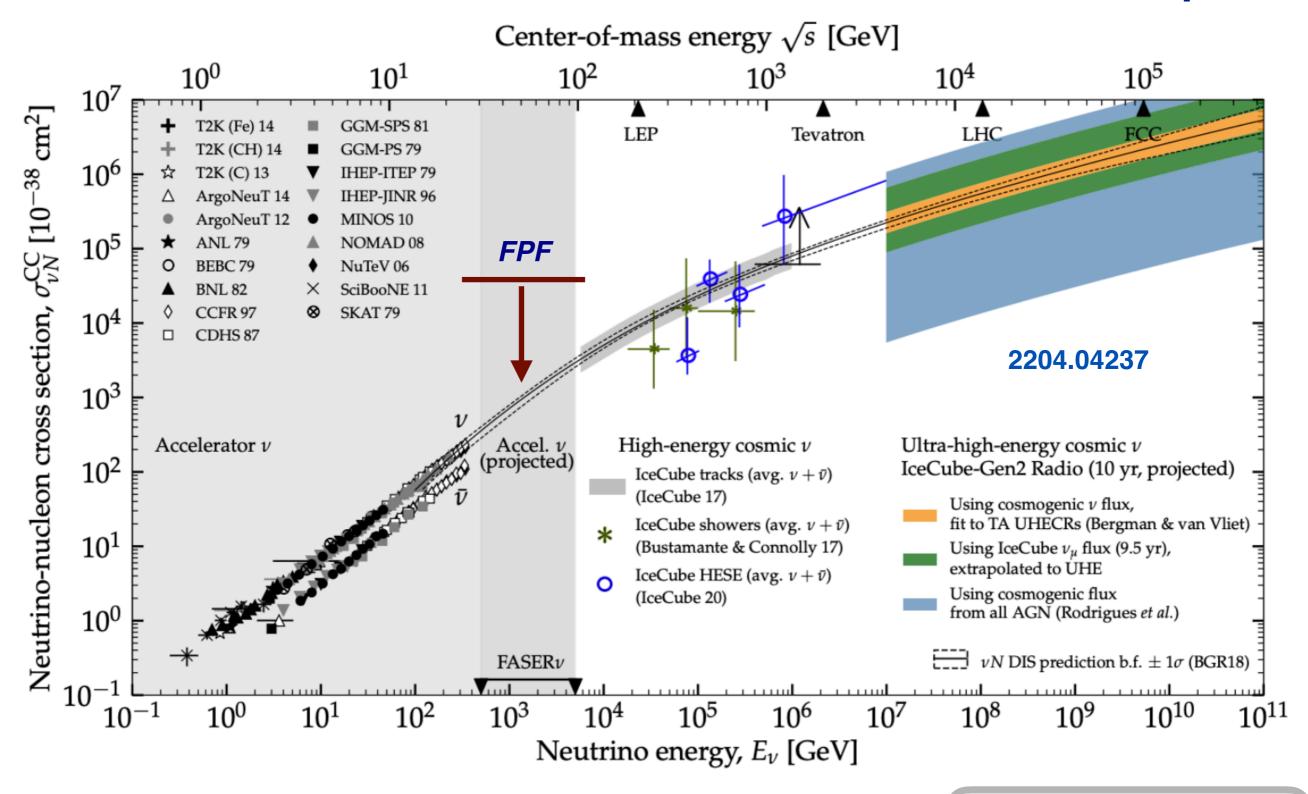


- General improvements of low-mass gluon-initiated processes at the LHC
- Constraints also on the charm PDF via the gluon-charm initial state
- Run III data on LHC neutrinos may provide first measurement of the gluon content of the proton at x=10<sup>-7</sup>

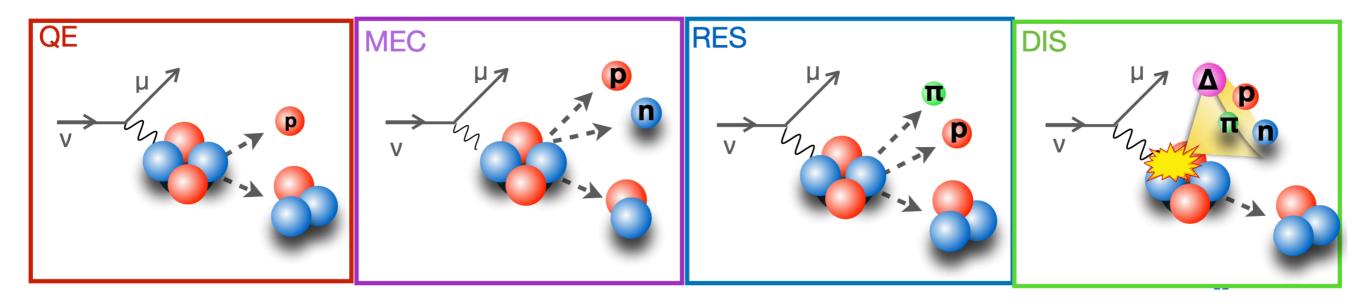


# Neutrino Structure Functions from GeV to EeV Energies

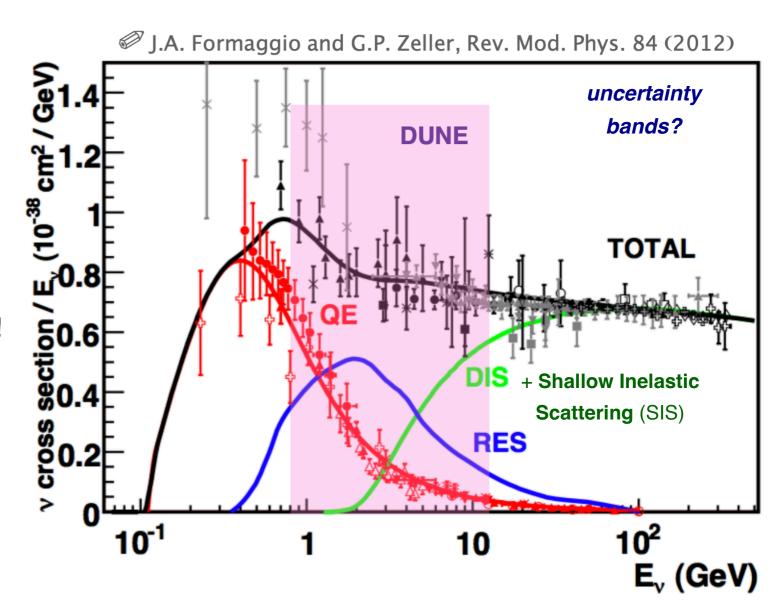
A. Candido, A. Garcia, G. Magni, T. Rabemananjara, J. Rojo, R. Stegeman, JHEP 23

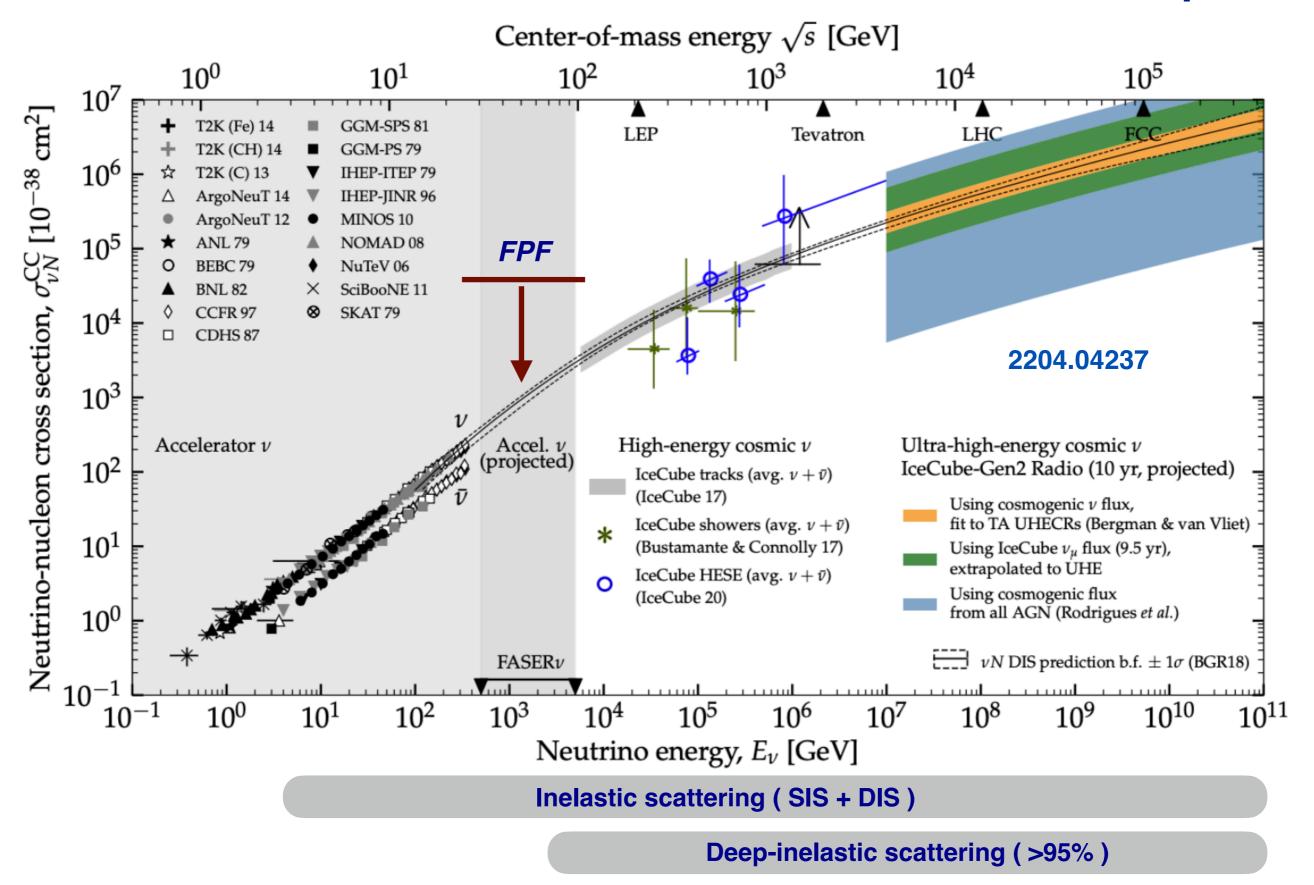


Depending on the neutrino energy, **different interaction mechanisms** dominate the neutrino-nucleus cross-section ``collider neutrinos" (FASER, SND@LHC, FPF) cover uncharted range



- For energies > 5 GeV, inelastic scattering dominates the inclusive cross-section
- Common misconception: inelastic scattering does not coincide with deep-inelastic scattering (DIS) where pQCD can be applied!
- How robust is our theoretical understanding of neutrino inelastic scattering interactions?





sizable kinematic region where Shallow Inelastic Scattering (SIS) cannot be neglected

#### The role of the low-Q region

inclusive neutrino cross-sections receives **sizeable contributions from** *Q* < 2 GeV **region**, where structure functions cannot be evaluated in the pQCD framework

$$\sigma(\mathbf{E}_{\nu}) = \int_{Q_{\min}^2}^{2m_N \mathbf{E}_{\nu}} dQ^2 \left[ \int_{Q^2/(2m_N y \mathbf{E}_{\nu})}^1 dx \, \frac{d^2 \sigma}{dx dQ^2}(x, Q^2, \mathbf{E}_{\nu}) \right]$$

$$\frac{d^2 \sigma^{\nu A}(x,Q^2,y)}{dxdy} = \frac{G_F^2 s/4\pi}{\left(1 + Q^2/m_W^2\right)^2} \left[Y_+ F_2^{\nu A}(x,Q^2) - y^2 F_L^{\nu A}(x,Q^2) + Y_- x F_3^{\nu A}(x,Q^2)\right]$$

-1

**Deep-Inelastic Scattering** 

$$F_i^{\nu A}(x,Q^2) = \sum_{j=q,\bar{q},g} \int_x^1 \frac{dz}{z} C_{i,j}^{\nu N}(z,\alpha_s(Q^2)) f_j^{(A)}\left(\frac{x}{z},Q^2\right)$$

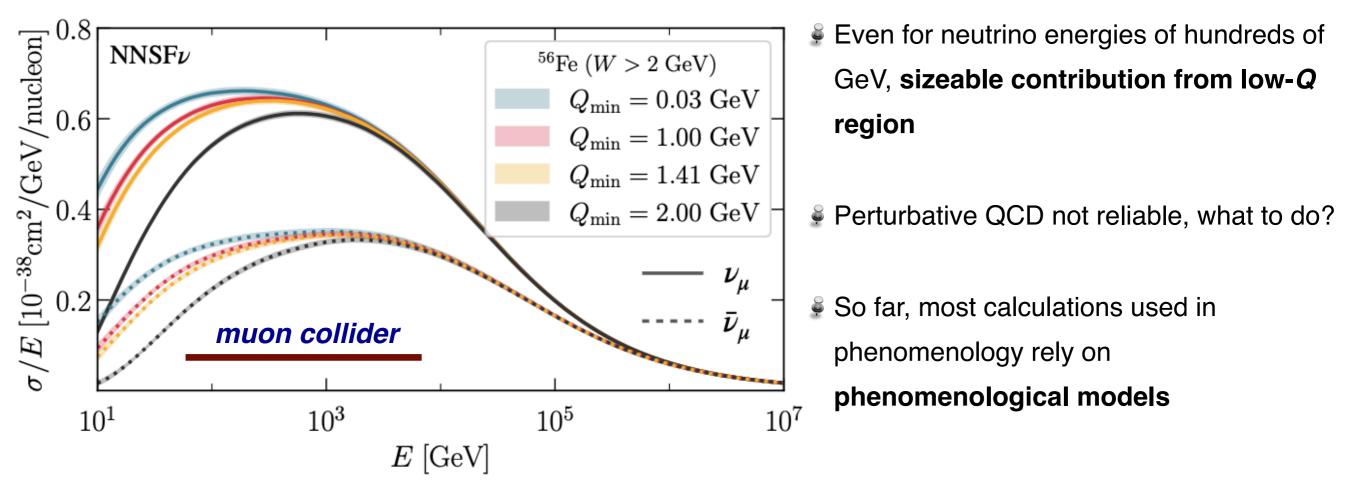
**Shallow-Inelastic Scattering** 

$$F_i^{\nu A}(x,Q^2) = ?$$

#### The role of the low-Q region

inclusive neutrino cross-sections receives **sizeable contributions from** *Q* < 2 GeV **region**, where structure functions cannot be evaluated in the pQCD framework

$$\sigma(\boldsymbol{E}_{\boldsymbol{\nu}}) = \int_{Q_{\min}^2}^{2m_N \boldsymbol{E}_{\boldsymbol{\nu}}} dQ^2 \left[ \int_{Q^2/(2m_N y \boldsymbol{E}_{\boldsymbol{\nu}})}^1 dx \, \frac{d^2 \sigma}{dx dQ^2}(x, Q^2, \boldsymbol{E}_{\boldsymbol{\nu}}) \right]$$



# The Bodek-Yang model

The Bodek-Yang model is popular to describe inelastic neutrino DIS structure functions

based on **effective leading-order PDFs** (GRV98LO) supplemented to phenomenological scaling variables and *K*-factors to improve agreement with data

$$f_i^{\text{LO}}(x,Q^2) \to f_i^{\text{LO,BY}}(\xi,Q^2) \qquad \xi = \frac{2x(Q^2 + m_f^2 + B)}{Q^2 \left[1 + \sqrt{1 + (2m_N x)^2/Q^2}\right] + 2Ax}$$

**Limitations** of the BY model of neutrino structure functions:

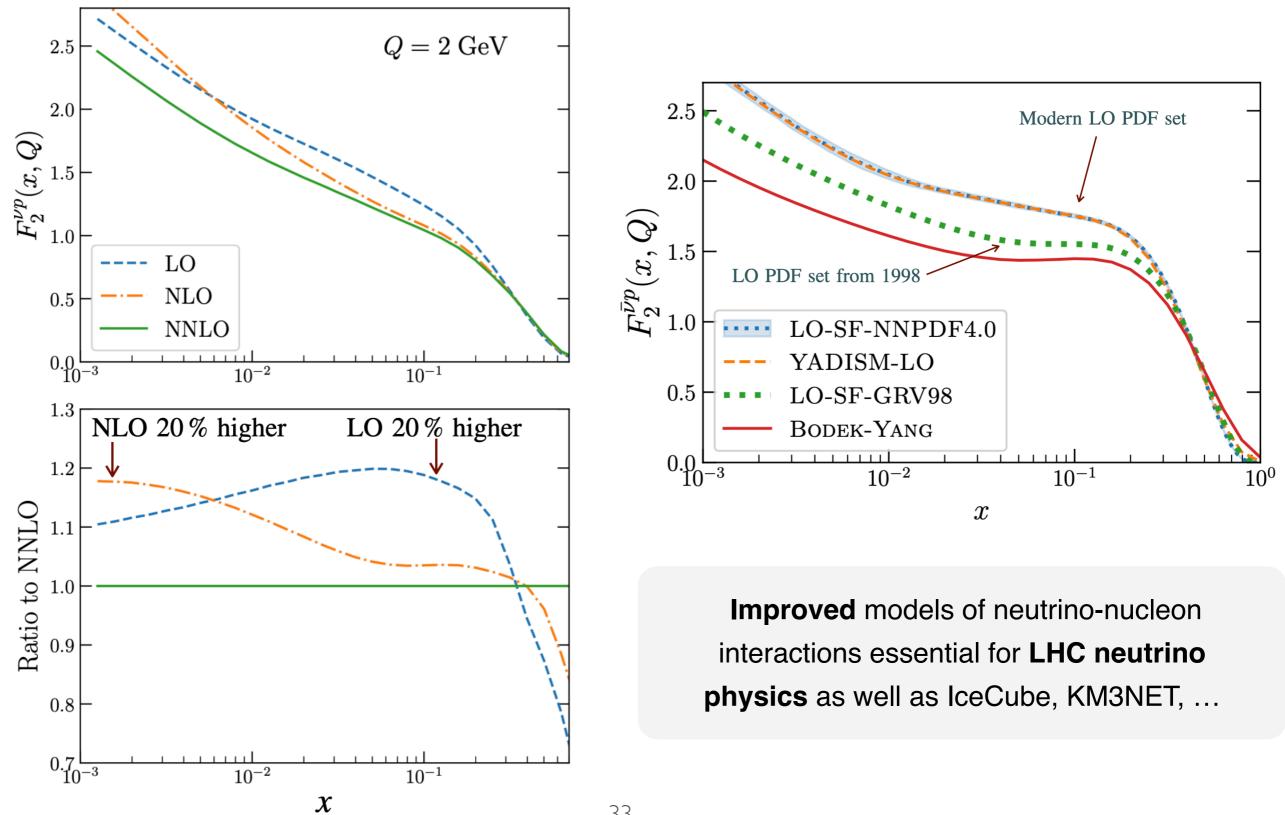
Obsolete PDF parametrisation that ignores constraints from the last 25 years

Neglects higher-order QCD corrections (can be up to 100%)

- Seannot be used above 100 TeV: not an option for UHE neutrinos
- Does not provide **uncertainty estimate**, difficult to assess its accuracy and precision
- Cannot be systematically improvable e.g. by new data

# The Bodek-Yang model

The Bodek-Yang model is popular to describe inelastic neutrino DIS structure functions

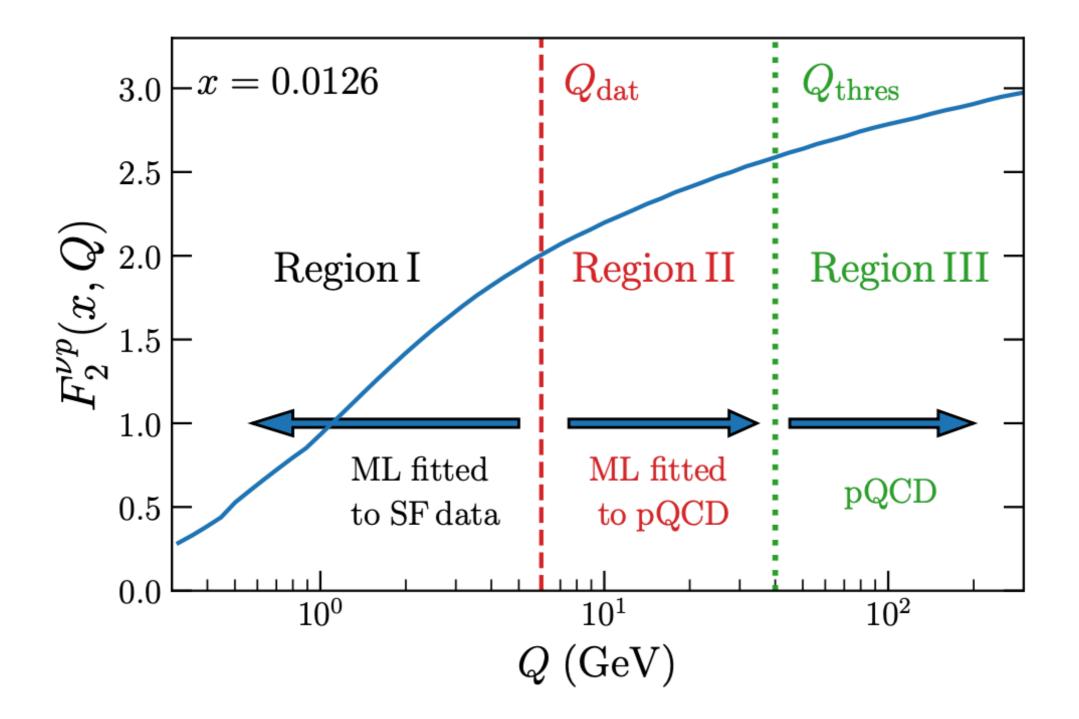


# The NNSFv approach

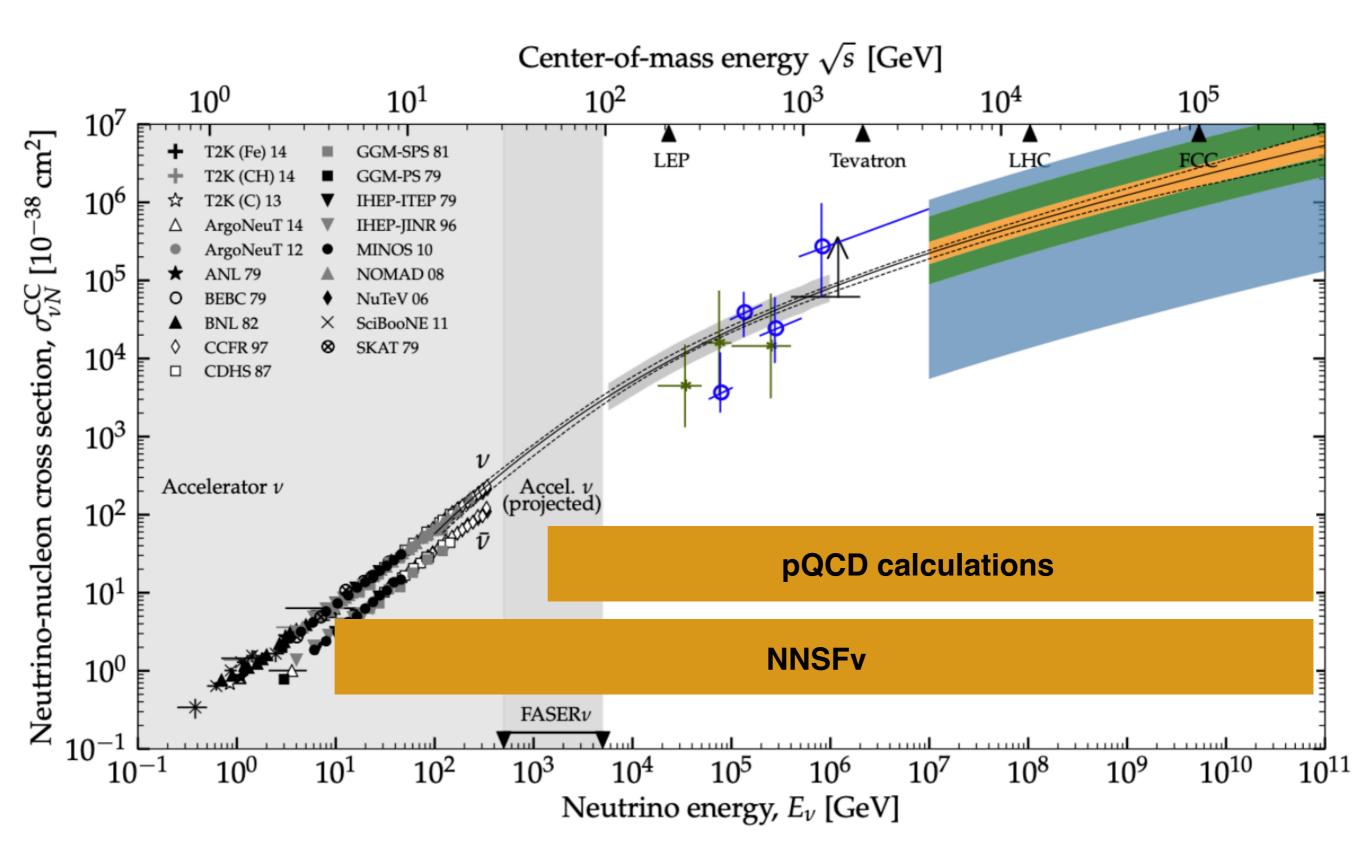
Motivation: realise the first determination of neutrino structure functions valid from

photoproduction Q = 0 all the way to Q = 100 TeV, enabling calculation of inclusive inelastic

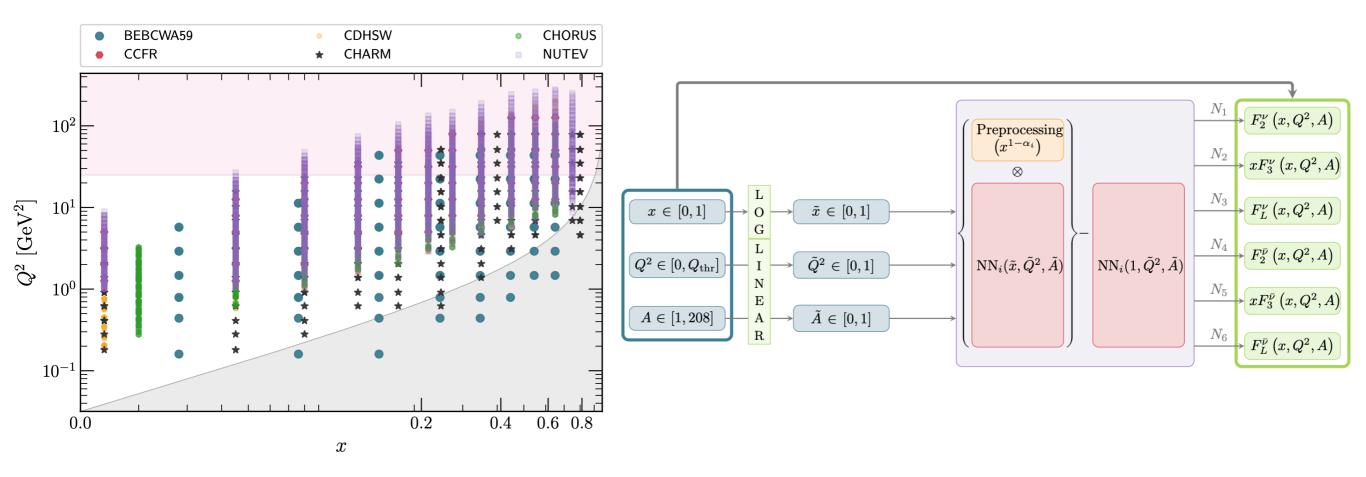
cross-sections for neutrinos from 5 GeV to 1012 GeV energies



# The NNSFv approach



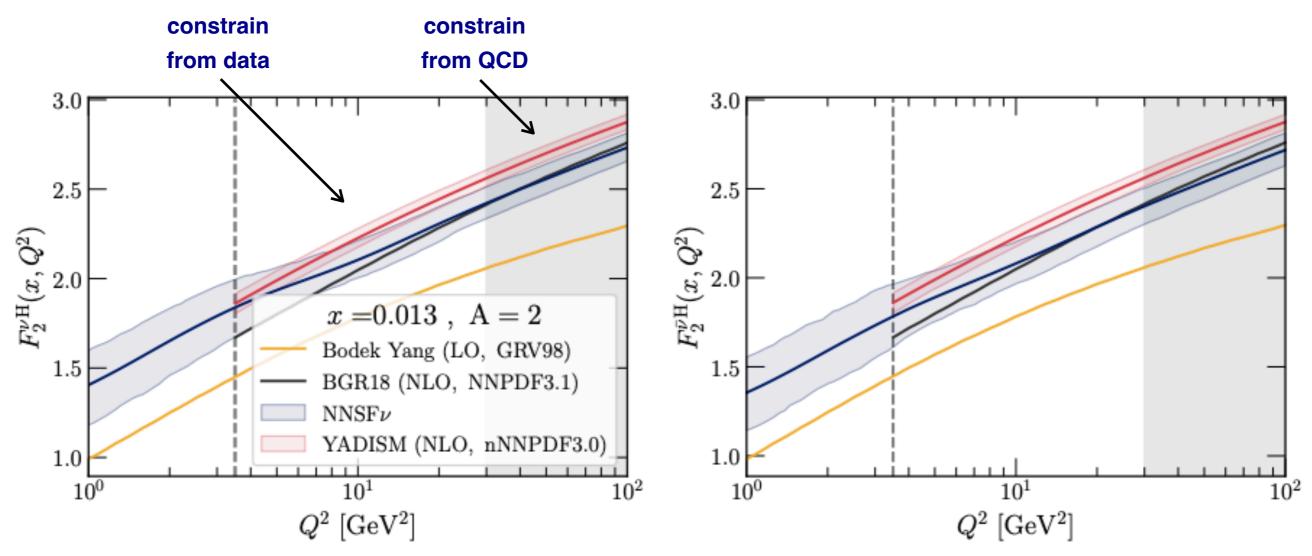
Use available data on neutrino-nucleus scattering to parametrise and determine inelastic structure functions by means of the NNPDF fitting methodology



- Finily  $\Im$  This data-driven parametrisation is made to **converge to the pQCD calculation** for large enough  $Q^2$  values as implemented with Lagrange multipliers
- $\frac{1}{2}$  In the neutrino energy region sensitive only to Q > few GeV, replace by pQCD calculation

consistent determination of neutrino structure functions valid for 12 orders of magnitude from  $E_{nu} = few \text{ GeV}$  up to  $E_{nu}=10^{12} \text{ GeV}$ 

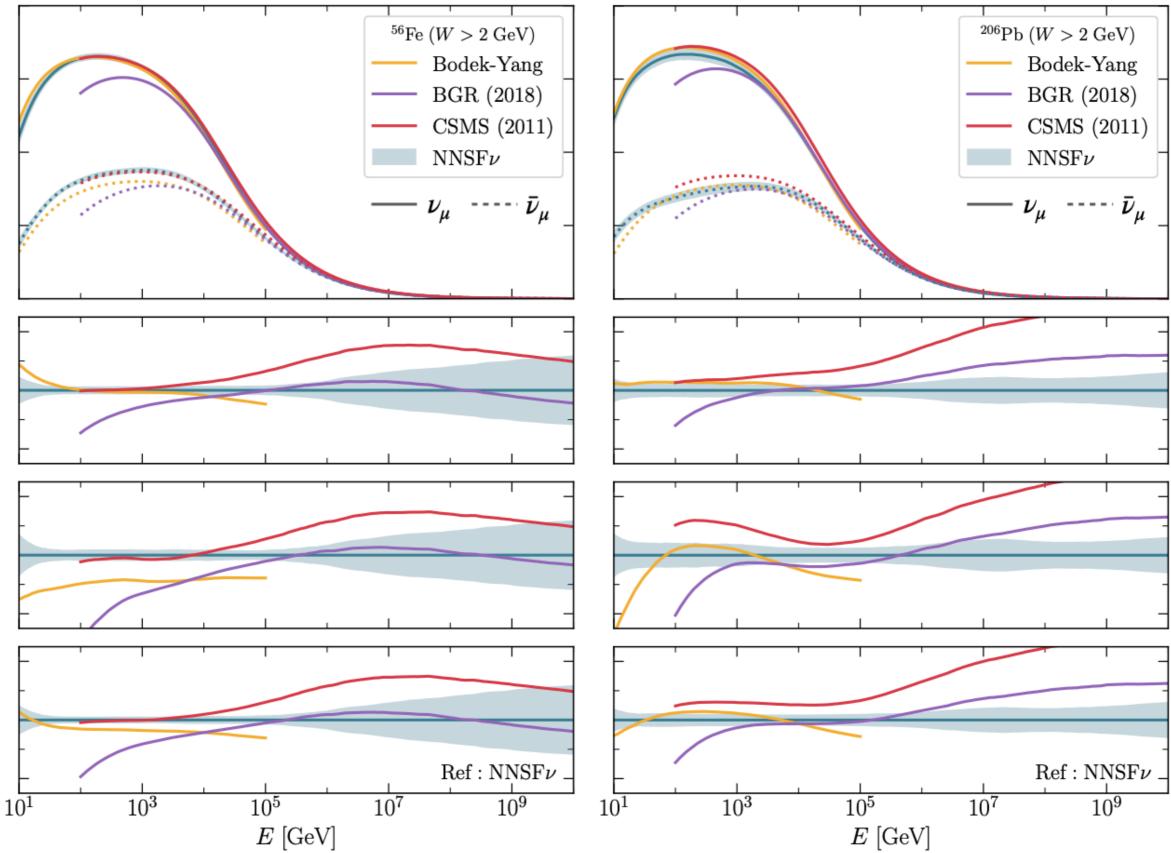
# The NNSFv results



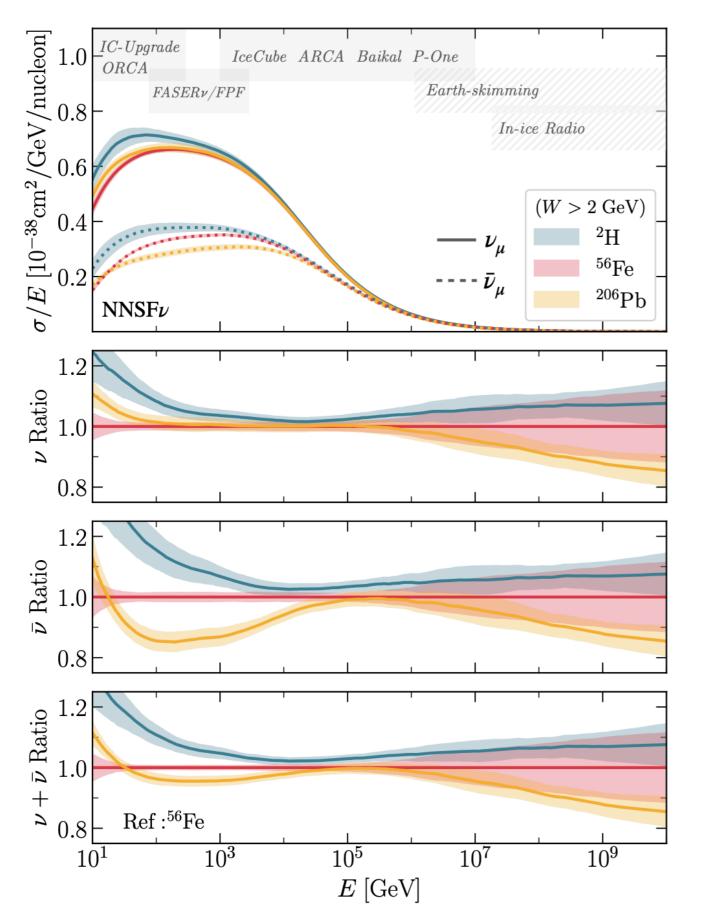
Smooth matching between data-driven and pQCD regions, uncertainty estimate in whole energy range

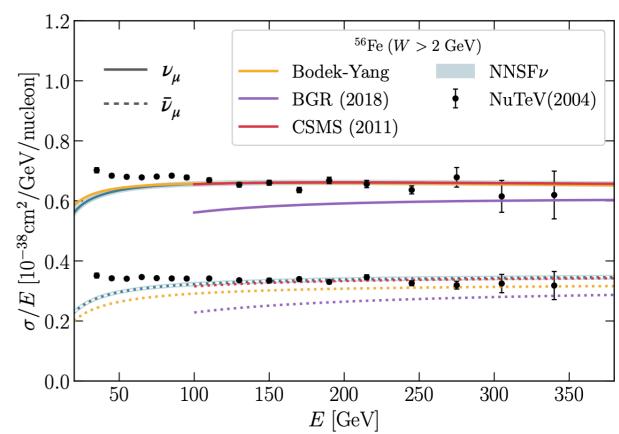
- Structure functions and integrated cross-sections available via user-friendly LHAPDF grids
- For the first time, a **unique theory prediction** for neutrino inelastic scattering suitable for neutrinos with energies from a few GeV up to the multi-EeV region

#### The NNSFv results



#### The NNSFv results





- Agreement with available neutrino structure function and cross-section data
- Estimate experimental & theory uncertainties
- Model-independent determination of nuclear corrections

Implemented in GENIE: ready to be used in your neutrino simulations!

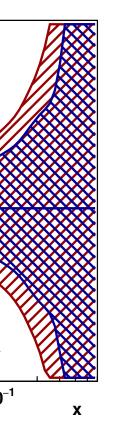
### Summary and outlook

LHC neutrinos realise an exciting program in a broad range of topics from **BSM and long-lived cles** to **neutrinos, QCD, and hadron structure**, with connections to astroparticle physics

surements of **muon neutrino DIS structure functions** at the LHC open a new probe to n and nuclear structure: a charged-current counterpart of the Electron Ion Collider

surements of **electron and tau neutrino event rates** at the LHC can constrain the small-x n in unexplored regions by using **dedicated observables** where QCD errors cancel out

nproved modelling of the **low-Q region in inelastic scattering** required for many neutrino riments and will be precisely tested with LHC neutrinos



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