





The structure of the proton in the LHC high-precision era

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Theoretical High-Energy Physics seminar

Lund University, 19/10/2018

The inner life of protons

See also ``The structure of the proton in the LHC precision era" Gao, Harland-Lang, JR (Physics Reports 17)

The Higgs boson

Huge gap between Higgs and Plank scales?

Elementary or composite? More Higgs bosons?

Coupling to Dark Matter? Role in cosmological phase transitions?

Is the vacuum state of the Universe stable?







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Dark matter

- Weakly interacting massive particles?
 Neutrinos? Ultralight particles (axions)?
- Interactions with SM particles? Selfinteractions?
- Structure of the Dark Sector?



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The Higgs boson

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Is the vacuum state of the Universe stable?

Quarks and leptons

Why **3 families?** Origin of **masses, mixings**?

Origin of Matter-Antimatter asymmetry?

Are **neutrinos Majorana or Dirac**? CP violation in the lepton sector?

Dark matter

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Crucial information on these fundamental questions will be provided by the LHC: the **exploration of the high-energy frontier** has just started!

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Parton distributions @ LHC

QCD Factorisation theorem:

Event rates = **parton distributions** + hard-scattering partonic cross-sections



Parton distributions @ LHC





DGLAP evolution (upwards in Q)

Momentum sum rule (energy conservation)

$$\frac{\partial}{\partial \ln Q^2} f_i(x, Q^2) = \int_x^1 \frac{dz}{z} P_{ij}\left(\frac{x}{z}, \alpha_s(Q^2)\right) f_j(z, Q^2)$$
$$\int_x^1 \int_x^{n_f} \sum_{i=1}^n f_i(x, Q^2) f_i(z, Q^2) \int_x^{n_f} f_i(z, Q^2) f_i(z, Q^2) \int_x^{n$$

$$\int_{0} dx \, x \left(\sum_{i=1}^{3} \left(q_{i}((x, Q^{2}) + \bar{q}_{i}(x, Q^{2})) + g(x, Q^{2}) \right) = 1$$

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QCD factorisation

The QCD factorization theorems guarantees PDF universality

$$\sigma_{lp \to \mu X} = \widetilde{\sigma}_{u\gamma \to u} \otimes u(x) \implies \sigma_{pp \to W} = \widetilde{\sigma}_{u\bar{d} \to W} \otimes u(x) \otimes \bar{d}(x)$$



Determine PDFs in **lepton-proton collisions** (deep-inelastic scattering) ...

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... and use them to compute predictions for **proton-proton collisions**

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From the proton mass to the LHC

- Extract PDFs at hadronic scales (few GeV), where non-perturbative QCD sets in
- Use perturbative evolution to compute PDFs at high scales as input to LHC predictions



Why better PDFs?



Parton Distributions and the SMEFT

Heavy bSM states beyond direct reach can still modify LHC cross-sections

$$\mathscr{L}_{\text{SMEFT}} = \mathscr{L}_{\text{SM}} + \sum_{i}^{N_{d6}} \frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \sum_{j}^{N_{d8}} \frac{b_j}{\Lambda^4} \mathcal{O}_j^{(8)} + \dots$$

Some operators induce energy-growing effects: exploit unique LHC kinematic reach

$$\sigma(\boldsymbol{E}) = \sigma_{\rm SM}(\boldsymbol{E}) \left(1 + \sum_{i}^{N_{d6}} \omega_i \frac{\boldsymbol{c}_i m_{\rm SM}^2}{\Lambda^2} + \sum_{i}^{N_{d6}} \widetilde{\omega}_i \frac{\boldsymbol{c}_i \boldsymbol{E}^2}{\Lambda^2} + \mathcal{O}\left(\Lambda^{-4}\right) \right)$$

LHC 13 TeV, NNLO, α_s=0.118

LHC 13 TeV, NNLO, α_s=0.118



The NNPDF approach to PDF fits



- Neural Networks as universal unbiased interpolants to parametrise PDFs: eliminate model assumptions
- Monte Carlo replicas to propagate uncertainties wo Gaussian assumptions
- Genetic algorithms and Machine
 Learning to explore parameter space

	Proton PDFs	Nuclear PDFs
Traditional	$g(\mathbf{x}) \simeq \mathbf{x}^{-b}(1-\mathbf{x})^c$	$R_g(x, A) \simeq (1 + bx + cx^2) \times A^d$
leural Nets	$g(x) \simeq NN(x)$	$R_g(x, A) \simeq NN(x, A)$

Λ

The NNPDF approach to PDF fits



Combine precision measurements and state-of-the-art theory within robust statistical framework

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The NNPDF approach to PDF fits



Highly non-trivial validation of the QCD factorisation framework: Including O(5000) data points, from O(40) experiments, some of them with $\approx 1\%$ errors, yet the global PDF fit achieves $\chi^2/N_{dat} \approx 1$!

PDF constraints from precision LHC data

PDF information from p+p collisions



Gluon PDF from top quarks

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Second Complementary probe of the large-x gluon

Included differential top distributions in NNPDF3.0 NNLO: constraints on large-x gluon comparable to inclusive jet production

Improved theory uncertainties in regions crucial for BSM searches, *i.e.*, m_{tt} > 1 TeV (while fitting only y_t and y_{tt})



Gluon PDF from direct photons

- Revisited the impact of LHC direct photon data into the global PDF fit Campbell, JR, Slade, Williams 18
- Theory based on NNLO QCD and LL electroweak calculations
- Moderate impact on medium-x gluon
- Good **consistency** with other gluonsensitive experiments in NNPDF3.1

	NNPDF3.1	NNPDF3.1+ATLAS γ
Fixed-target lepton DIS	1.207	1.203
Fixed-target neutrino DIS	1.081	1.087
HERA	1.166	1.169
Fixed-target Drell-Yan	1.241	1.242
Collider Drell-Yan	1.356	1.346
Top-quark pair production	1.065	1.049
Inclusive jets	0.939	0.915
$Z p_T$	0.997	0.980
Total dataset	1.148	1.146

Flavour separation from forward W,Z

Forward coverage of LHCb: unique sensitivity to small-x and large-x regions beyond that of ATLAS/CMS

Specially important to disentangle quark flavour at large-x

NNPDF3.1 NNLO, Q = 100 GeV

10⁻²

 10^{-3}

10⁻⁴

10⁻¹

NNPDF3.1 NNLO, Q = 100 GeV

Impact on the charm PDF

NNPDF3.1: independent parametrisation for charm PDF, constrained from data

NNPDF3.1 NNLO, Q = 1.7 GeV

Momentum Fraction of Charm Quarks

LHC electroweak data provide important info on charm content of protons

- Intrinsic' charm (IC) bounded to < 0.7% of proton momentum from LHC data
- Indications of a small but non-zero IC

$$C(Q^2) \equiv \int_0^1 dx \ x \ \left(c(c,Q^2) + \bar{c}(x,Q^2)\right)$$

Large intrinsic charm allowed in CT14IC **disfavoured** by LHC data

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Parton Distributions and theoretical uncertainties

NNPDF, in preparation

PDF uncertainties

PDF uncertainties receive contributions from different sources:

Theory uncertainties on PDFs from **Missing Higher Orders** (MHOs) never quantified!

Theory uncertainties from MHOs

At any finite order, perturbative QCD calculations depend on the unphysical **renormalisation** and **factorisation scales**

$$\sigma(\mu_R, \mu_F) = \sum_{k=0}^n \sum_{i,j}^{n_f} \alpha_s^{p+k}(\mu_R) \,\widetilde{\sigma}^{(k)}(\mu_R, \mu_F) \otimes q_i(\mu_F) \otimes q_j(\mu_F) + \mathcal{O}\left(\alpha_s^{p+n+1}\right)$$

In PDF fits, both scales are set to a fixed value, the typical **momentum transfer of the process** *Q*, and the MHOUs are ignored

$$\sigma(\mu_R = Q, \mu_F = Q) = \sum_{k=0}^n \sum_{i,j}^{n_f} \alpha_s^{p+k}(Q) \,\widetilde{\sigma}^{(k)}(Q) \otimes q_i(Q) \otimes q_j(Q)$$

At order N^kLO, the dependence on the two scales is determined by the N^{k-1}LO splitting functions and partonic cross-sections by imposing:

$$\sigma(\boldsymbol{\mu}_{R}, \boldsymbol{\mu}_{F}) = \sigma(\boldsymbol{Q}, \boldsymbol{Q}) + \mathcal{O}\left(\boldsymbol{\alpha}_{s}^{p+n+1}\right)$$

Scale variations provide an estimate of the perturbative MHOs

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Theory uncertainties from MHOs

How severe is **ignoring MHOUs** in modern global PDFs fits?

Shift between NLO and NNLO PDFs comparable or larger than PDF errors

Given the high precision of modern PDF determinations, **accounting for MHOUs** is most urgent

Option #1: PDF fits with scale variations

Perform PDF fits where **theory calculations** are constructed by exploring a range of values for the **renormalisation** and **factorisation scales**

Define the theory error due to MHOUs from envelope of the scale-varied PDFs

Option #1: PDF fits with scale variations

Perform PDF fits where **theory calculations** are constructed by exploring a range of values for the **renormalisation** and **factorisation scales**

The 7pt envelope seems to work fine in most cases (perhaps too conservative?)
Non-trivial theory-induced correlations between e.g. DIS and collider processes
CPU-intensive, and cumbersome or LHC applications - can we do better?

Construct a theory covariance matrix from scale-varied cross-sections and combine it with the experimental covariance matrix

$$\chi^{2} = \frac{1}{N_{\text{dat}}} \sum_{i,j=1}^{N_{\text{dat}}} \left(D_{i} - T_{i} \right) \left(\text{cov}^{(\text{exp})} + \text{cov}^{(\text{th})} \right)_{ij}^{-1} \left(D_{j} - T_{j} \right)$$

Different scale variation prescriptions possible, e.g, the 5-point prescription

$$\operatorname{cov}_{ii}^{(\text{th})} = \frac{1}{2} \left(\Delta_i(+,0)^2 + \Delta_i(-,0)^2 + \Delta_i(0,+)^2 + \Delta_i(0,+)^2 \right)$$
$$\operatorname{cov}_{jj}^{(\text{th})} = \frac{1}{2} \left(\Delta_i(+,0)\Delta_j(+,0) + \Delta_i(-,0)\Delta_j(-,0) + \frac{1}{4}(\Delta_i(0,+) + \Delta_i(0,-))(\Delta_j(0,+) + \Delta_j(0,-)) \right)$$
$$\Delta_i(-,0) \equiv \sigma(\mu_R = Q/2, \mu_F = Q) - \sigma(\mu_R = Q, \mu_F = Q)$$
$$\Delta_i(+,+) \equiv \sigma(\mu_R = 2Q, \mu_F = 2Q) - \sigma(\mu_R = Q, \mu_F = Q)$$

Construct a theory covariance matrix from scale-varied cross-sections and combine it with the experimental covariance matrix

Systematic validation of the theory covariance matrix on the `exact' result, the NNLO-NLO shift, with the O(5000) data points of the global fit

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Experiment correlation matrix NMC 0.75 SLAC BCDMS 0.50 CHORUS 0.25 NTVDMN 0.00 HERACOMB -0.25 HERAFZOHAB -0.50ATLAS CMS -0.75 CDF HERAF280149040 ATLAS LHCb CHORUS HERACOMB NTVDMN CMS NMELAC BCDMS OFFCD

1.00

The small-x gluon and neutrino astronomy

Gauld, JR, Rottoli, Talbert 15 Gauld, JR, Rottoli, Sarkar, Talbert 15 Gauld, JR 17 Bertone, Gauld, Rojo 18

The small-x gluon from HERA data

- Small-x gluon unconstrained: information from HERA ends for x<10⁻⁴
- Very large uncertainties in global fits
- Need processes covering x<10⁻⁴ region

- Include LHCb D meson production at 5, 7, 13 TeV
- Fit normalised distributions & ratios between CoM energies to reduce MHOUs

$$N_X^{ij} = \frac{d^2\sigma(\text{X TeV})}{dy_i^D d(p_T^D)_j} \left/ \frac{d^2\sigma(\text{X TeV})}{dy_{\text{ref}}^D d(p_T^D)_j} \right|_{T_{\text{ref}}}$$
$$R_{13/X}^{ij} = \frac{d^2\sigma(13 \text{ TeV})}{dy_i^D d(p_T^D)_j} \left/ \frac{d^2\sigma(\text{X TeV})}{dy_i^D d(p_T^D)_j} \right|_{T_{\text{ref}}}$$

gluon PDF uncertainties reduced by factor 10 at $x \approx 10^{-6}$

Excellent description of all LHCb datasets

and ratios (after errata corrected)

$N_5(84)$	$N_{7}(79)$	$N_{13}(126)$	$R_{13/5}(107)$	$R_{13/7}(102)$
1.97	1.21	2.36	1.36	0.80
0.86	0.72	1.14	1.35	0.81
1.31	0.91	1.58	1.36	0.82
0.74	0.66	1.01	1.38	0.80
1.08	0.81	1.27	1.29	0.80
1.53	0.99	1.73	1.30	0.81
1.07	0.81	1.34	1.35	0.81
0.82	0.70	1.07	1.35	0.81
0.84	0.71	1.10	1.36	0.81

Cross-section ratios: improved perturbative stability

Cross-section ratios: improved perturbative stability

'New Physics' within QCD

Science Life and Physics

The Guardian

Jon Butterworth

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✓ @jonmbutterworth Thu 28 Dec 2017 17.30 GMT

After 40 years of studying the strong nuclear force, a revelation

This was the year that analysis of data finally backed up a prediction, made in the mid 1970s, of a surprising emergent behaviour in the strong nuclear force

In the mid 1970s, four Soviet physicists, Batlisky, Fadin, Kuraev and Lipatov, made some predictions involving the strong nuclear force which would lead to their initials entering the lore. "BFKL" became a shorthand for a difficult-to-

BFKL dynamics at small-x

- QCD calculations in the DGLAP factorisation framework successful in describing data from proton-proton and electron-proton collisions
- Need to go beyond DGLAP: at small-x, logarithmically enhanced terms in 1/x become dominant and need to be resummed to all orders
- BFKL (high-energy, small-x) resummation can be matched to DGLAP collinear framework and included into PDF fits

$$\begin{array}{ll} \begin{array}{l} \textbf{DGLAP} \\ \textbf{Evolution in } Q^2 \end{array} & \frac{\partial}{\partial \ln Q^2} f_i(x,Q^2) = \int_x^1 \frac{dz}{z} P_{ij}\left(\frac{x}{z},\alpha_s(Q^2)\right) f_j(z,Q^2) \\ \\ \textbf{BFKL} \\ \textbf{Evolution in } x \end{array} & \frac{\partial}{\partial \ln 1/x} f_+(x,Q^2) = \int_0^\infty \frac{d\nu^2}{\nu^2} K\left(\frac{Q^2}{\nu^2},\alpha_s(Q^2)\right) f_+(x,\nu^2) \\ \\ \textbf{ABF, CCSS, TW} \\ \textbf{+ others, 94-08} \end{array} & P_{ij}^{N^k LO + N^h LLx}(x) = P_{ij}^{N^k LO}(x) + \Delta_k P_{ij}^{N^h LLx}(x) \end{array}$$

BFKL dynamics at small-x

Ball, Bertone, Bonvini, Marzani, JR, Rottoli 17

more data from the **small-x region**

Best description of **small-***x***HERA data** only possible with **BFKL effects!**

Forward charm production revisited

LHCb D meson production included in NNPDF3.1sx (N)NLO+NLLx fits

Similar reduction of gluon PDF errors at **small-***x* + **increase in central value**

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Forward charm production revisited

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Q = 1.7 GeV

Neutrino telescopes

Ultra-high energy (UHE) neutrinos: novel window to the extreme Universe!

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Unveiling cosmic neutrino origin

Neutrino telescopes as QCD microscopes

signal: cosmic neutrino - nucleus scattering

background: prompt charm production

Neutrino telescopes as QCD microscopes

signal: cosmic neutrino - nucleus scattering

background: prompt charm production

Sensitive to **small-***x* **quarks** (and thus gluons via evolution) down to $\mathbf{x} \approx \mathbf{10^{-8}}$ and $\mathbf{Q} \approx \mathbf{M_W}$

Sensitive to small-x gluons down to $x \approx 10^{-6}$ and $Q \approx M_{charm}$ in the centre-of-mass frame

Neutrino telescopes as QCD microscopes

signal: cosmic neutrino - nucleus scattering

background: prompt charm production

UHE neutrino-nucleus cross-section

Bertone, Gauld, JR 17

State-of-the-art predictions for ultra-high energy neutrino interactions

- BFKL small-x effects in PDFs and deep-inelastic structure functions
- Constraints on small-x PDFs from LHCb charm production
- Accounting for **nuclear corrections** and heavy-quark-initiated contributions

UHE neutrino-nucleus cross-section

- Differences both at intermediate (better PDFs, improved treatment of heavy quarks) and high energies (LHCb constraints, BFKL effects)
- Nuclear effects important: constrain them with LHCb charm production in p+Pb
- IceCube and other neutrino telescopes are the ultimate QCD microscopes!

Towards Ultimate PDFs at the High Lumi LHC

Abdul-Khalek, Bailey, Gao, Harland-Lang, JR 18 + HL/HE-LHC Yellow Report, to appear

A luminous future

In the framework of the update of the European Strategy for Particle Physics, a CERN Yellow Report will evaluate the physics potential of the HL-LHC (to appear in Dec 2018)

We have studied the impact of HL-LHC data on PDFs, including projections with (future) LHCb measurements.

What is the **ultimate precision** that can be expected for PDFs from **hadron collider data?**

Process	Kinematics	$N_{\rm dat}$
$Z p_T$	$\begin{array}{ c c c c } 20\mathrm{GeV} \leq p_T^{ll} \leq 3.5\mathrm{TeV} \\ 12\mathrm{GeV} \leq m_{ll} \leq 150\mathrm{GeV} \\ y_{ll} \leq 2.4 \end{array}$	162
high-mass Drell-Yan	$\begin{vmatrix} p_T^{l1(2)} \ge 40(30) \text{GeV} \\ \eta^l \le 2.5, m_{ll} \ge 116 \text{GeV} \end{vmatrix}$	21
top quark pair	$m_{t\bar{t}} \simeq 5 \text{ TeV}, y_t \le 2.5$	26
W+charm (central)	$ \begin{array}{ c c } p_T^{\mu} \geq 26 \mathrm{GeV}, p_T^c \geq 5 \mathrm{GeV} \\ \eta^{\mu} \leq 2.4 \end{array} $	6
W+charm (forward)	$p_T^{\mu} \ge 20 \text{GeV}, p_T^c \ge 20 \text{GeV}$ $p_T^{\mu+c} \ge 20 \text{GeV}$ $2 \le \eta^{\mu} \le 5, 2.2 \le \eta^c \le 4.2$	12
Direct photon	$E_T^{\gamma} \lesssim 3 \text{ TeV}, \eta_{\gamma} \leq 2.5$	53
Forward W, Z	$\begin{vmatrix} p_T^l \ge 20 \text{ GeV}, \ 2.0 \le \eta^l \le 4.5 \\ 2.0 \le y_{ll} \le 4.5 \\ 60 \le m_{ll} \le 120 \text{ GeV} \end{vmatrix}$	90
Inclusive jets	$ y \le 3, R = 0.4$	54
Total		424

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A luminous future

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HL-LHC constraints from LHCb

Projected forward W+charm data

Projected invariant tt mass data

Forward W+charm

Top quark pair production

HL-LHC measurements will be specially useful to constrain the **gluon** and **quark flavour separation** in the large-*x* region, including strangeness

HL-LHC constraints on PDFs

PDF uncertainties HLLHC / Current	10 GeV < M _X < 40 GeV	40 GeV < M _X < 1 TeV	1 TeV < M _X < 6 TeV
g-g luminosity	0.58 (0.49)	0.41 (0.29)	0.38 (0.24)
q-g luminosity	0.71 (0.65)	0.49 (0.42)	0.39 (0.29)
quark-quark luminosity	0.78 (0.73)	0.46 (0.37)	0.60 (0.45)
quark-antiquark luminosity	0.73 (0.70)	0.40 (0.30)	0.61 (0.50)

gg => h+jet @ HL-LHC √s=14 TeV

HL-LHC data will lead to stringent constraints on PDFs, reducing uncertainties by up to a factor 4, and making possible precise predictions of central processes such as Higgs p_T

Summary and outlook

Quark/gluon structure of the proton: essential ingredient for LHC phenomenology

LHC data provides stringent constraints on PDFs, from the small-x gluon to the large-x antiquarks and to intrinsic charm

PDFs with theoretical uncertainties: the next milestone for global QCD analyses

Seutrino astrophysics requires direct input from small-x QCD

Precision determination of the small-x gluon and the UHE neutrino-nucleon cross-section thanks to LHCb charm data + BFKL resummation

At the **HL-LHC**, theory calculations with 1% PDF uncertainties are within reach

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