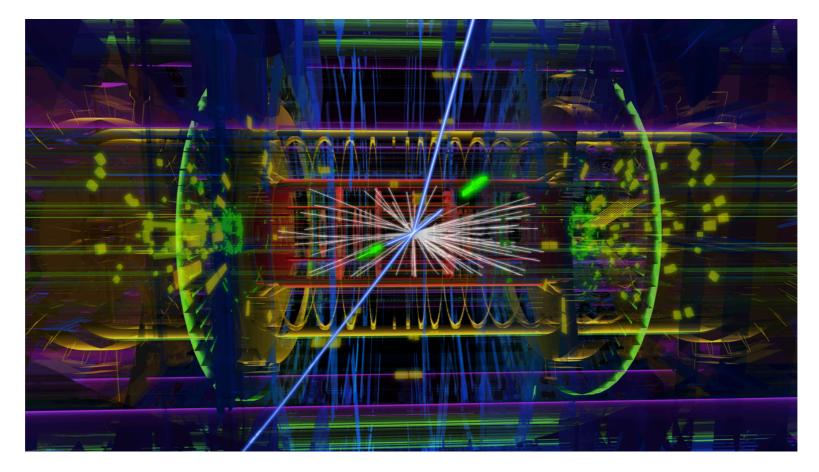
# **Precision Studies of the Higgs**

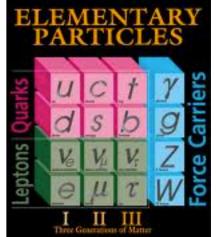


Giulia Zanderighi CERN & University of Oxford

Saturday Morning Lectures Oxford, 7<sup>th</sup> February 2015

# Status of particle physics

 Standard Model (SM): successful theory of strong (QCD), weak and electromagnetic (EW) elementary interactions



- Yet, no fundamental theory: theoretical issues + unexplained phenomena (e.g. gravity, matter anti-matter asymmetry, dark matter, dark energy, ...)
- The LHC is designed to

Investigation of the symmetry breaking (test origin of mass through the Higgs mechanism)

find physics beyond the SM (still to be done)

## BUT

Do we know what this really means?

What is the problem with particles having a mass?

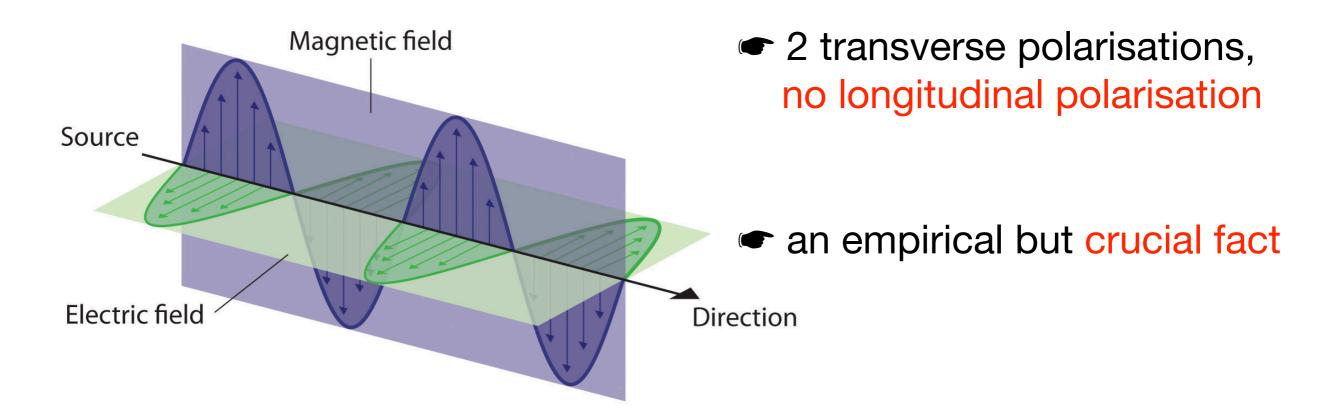
What is the Higgs mechanism & how does it solve the problem?

Why should there be New Physics at the TeV scale?

## Step back

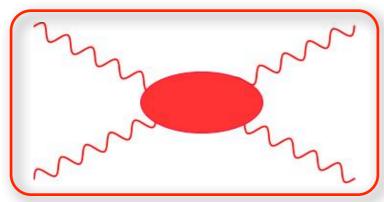
Duality in quantum field theory: wave ⇔ particle

### electromagnetic wave $\Leftrightarrow$ photon



# Gauge symmetry

If a 3<sup>rd</sup> longitudinal polarization existed

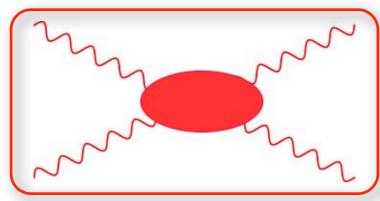


 $\Rightarrow$  scattering probability grows with energy

Violation of unitarity (probability > 1)  $\Rightarrow$  field theory breaks down

# Gauge symmetry

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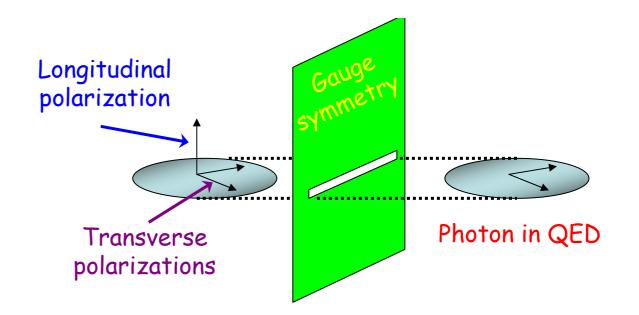


 $\Rightarrow$  scattering probability grows with energy

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### In QED:

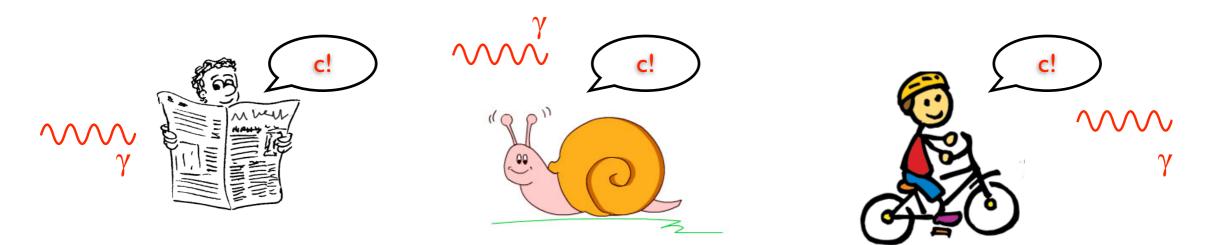
3rd polarization does not exist  $\leftrightarrow$  gauge symmetry



Gauge symmetry crucial to keep theory sensible at high energy

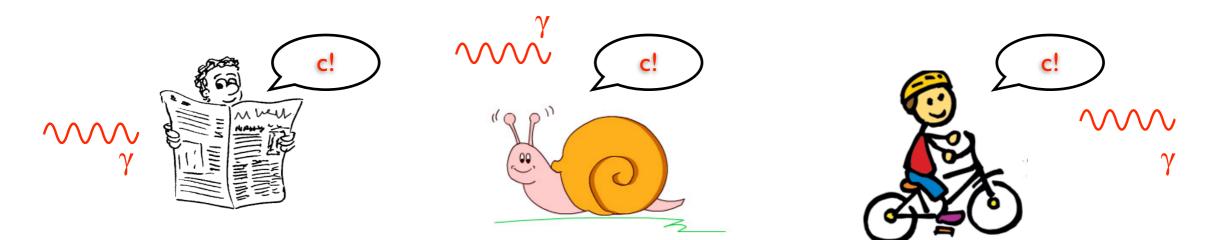
## Gauge filter and masses

From relativity: the speed of light is the same in all frames

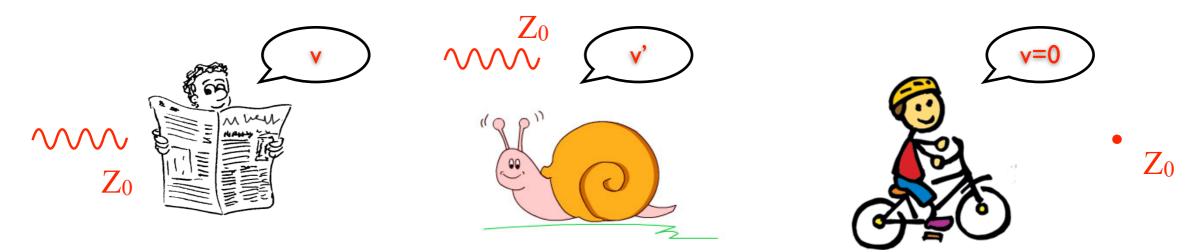


## Gauge filter and masses

From relativity: the speed of light is the same in all frames



For massive particle can choose a frame where the particle is at rest



In that frame, the distinction between transverse and longitudinal polarizations breaks rotational invariance

### Gauge trick does not work with massive particles

# EW symmetry breaking

SppS (1983-1985) pp collider at CERN, Geneva, running at  $E_{beam} = 450 \text{ GeV}$  LEP-II (1990-2001)  $e^+e^-$  collider at CERN, Geneva, running at E = 91.2  $\implies$  206 GeV

Z/W interactions are described by a EW gauge theory But Z/W masses break EW symmetry  $\Rightarrow$  theory breaks down at high E

gauge symmetry  $\leftrightarrow$  massless states  $\leftrightarrow$  sensible field theory

At what energy does this happen?



# EW symmetry breaking

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At what energy does this happen?



That's why the LHC was designed to investigate

- the mechanisms of mass generation
- how to keep the theory sensible at higher energy

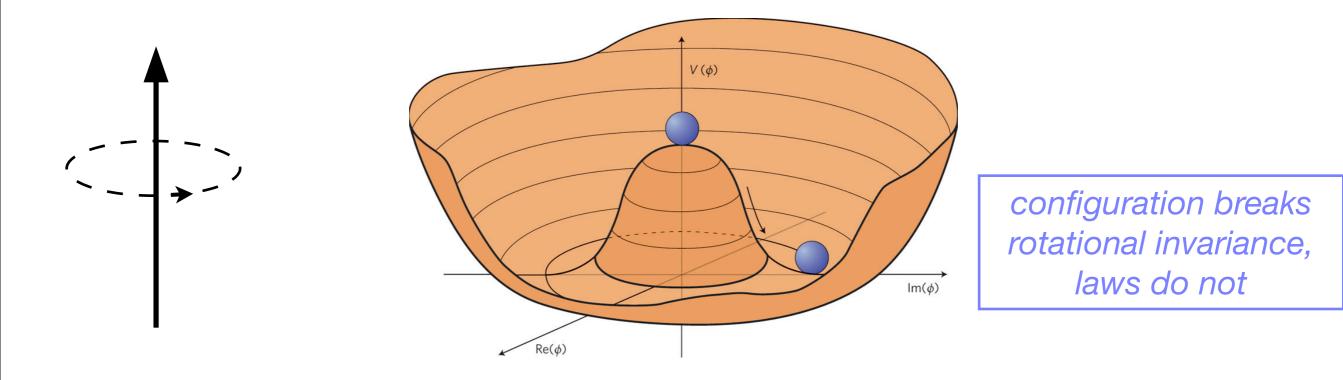
# Spontaneous symmetry breaking

Most popular solution:

Higgs mechanism, i.e. EW symmetry spontaneously broken

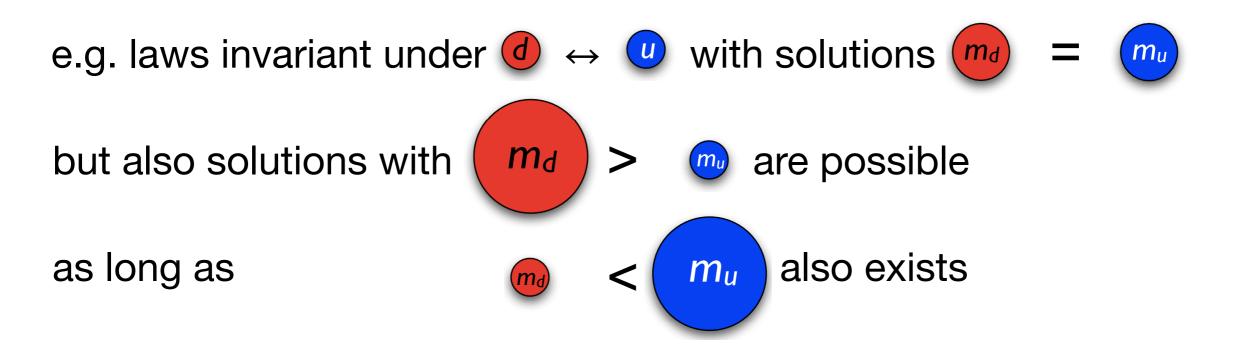
Spontaneous symmetry breaking (SSB): symmetry of equations but not of solutions

What does this mean ?



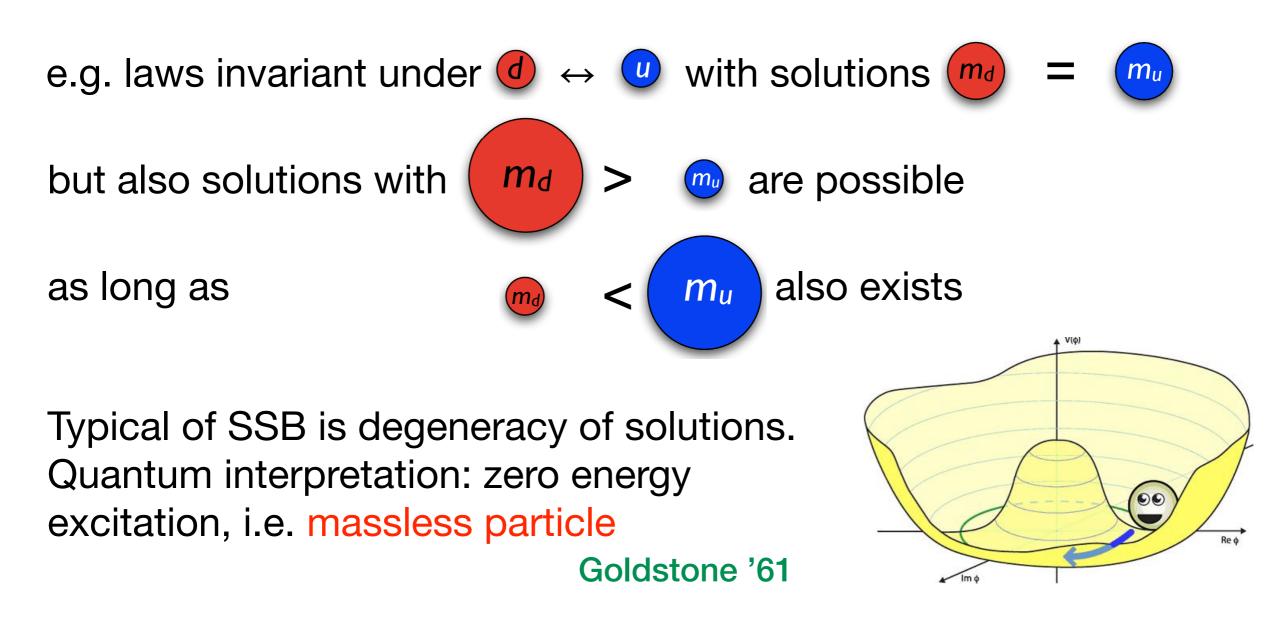
# Spontaneous symmetry breaking

With SSB relations implied by the exact symmetry can be modified



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With SSB relations implied by the exact symmetry can be modified



Problem: in Nature there is no massless Goldstone boson

# EW symmetry breaking

Higgs mechanism

with gauge interactions, zero-energy excitation absorbed by the gauge field  $\Rightarrow$  massive gauge particles and no Goldstone boson Brout, Englert, Higgs '64; Weinberg and Salam '67

# EW symmetry breaking

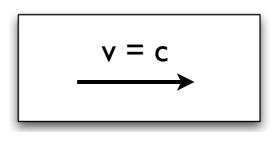
### <u>Higgs mechanism</u>

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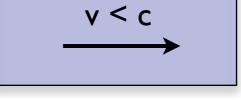
### <u>Higgs field</u>

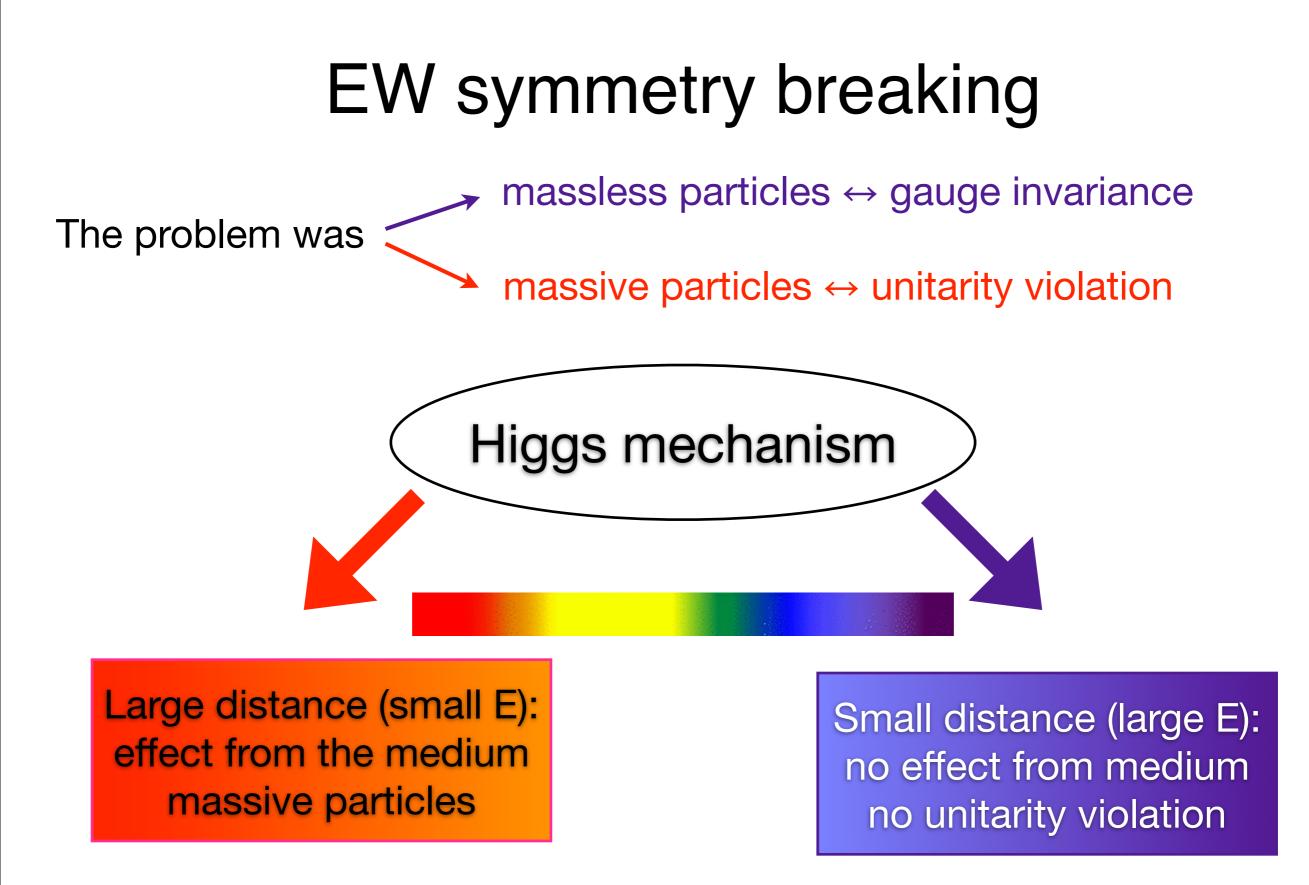
continuum medium pervading the whole universe. Particles interacting undergo a slow-down just as particles propagating in any medium do

### slow down $\Rightarrow$ inertia $\Rightarrow$ mass



vacuum

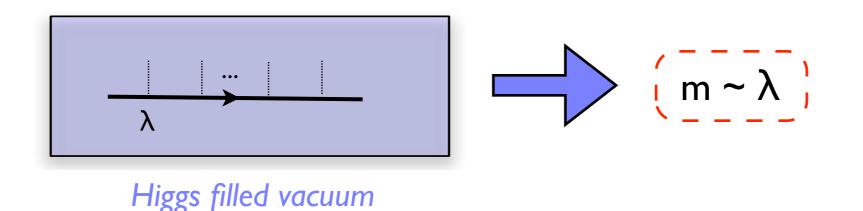




NB: EW charge distribution carries no electric charge  $\Rightarrow$  photon remains massless even after EWSB

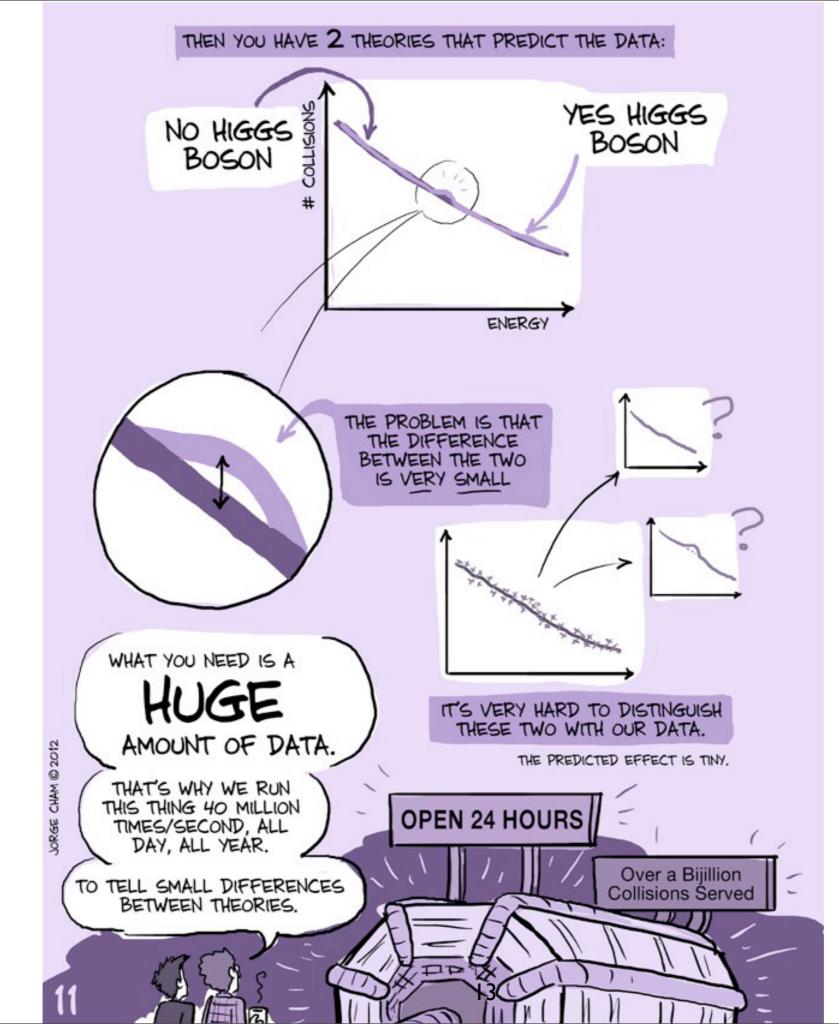
# Higgs mechanism in EW

- If the Higgs field exists, then quanta of the field must exist too
   ⇒ Higgs boson
- Coupling of a particle to Higgs is proportional to the particle's mass



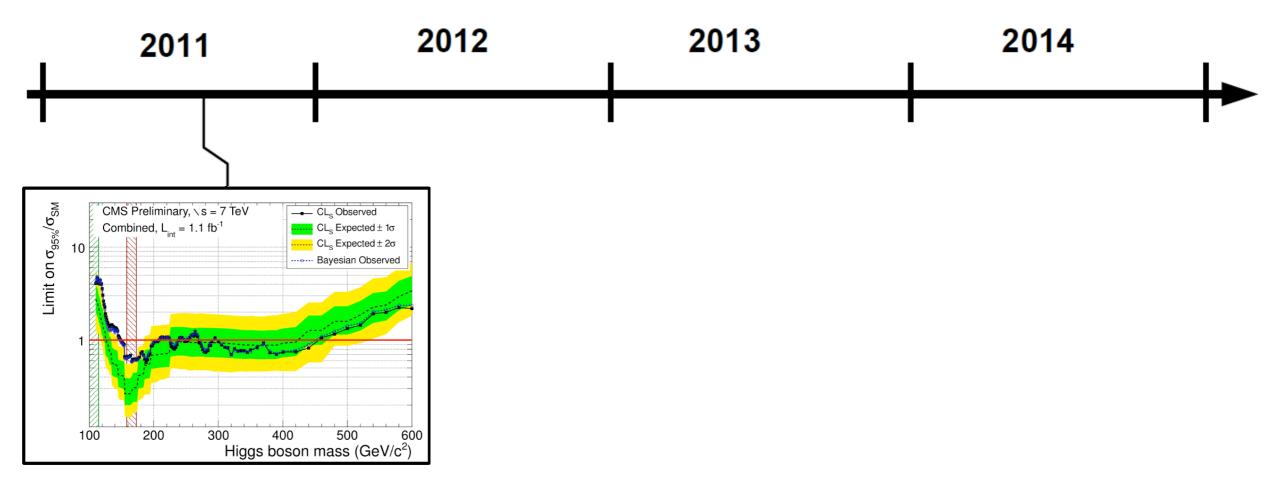
• The Higgs boson will have a mass too ... because the Higgs slows itself down as it propagates in the (Higgs) vacuum

 In the SM the Higgs mass is a free parameter, but once its value is determined everything else (couplings/masses) is fixed



# Before the discovery

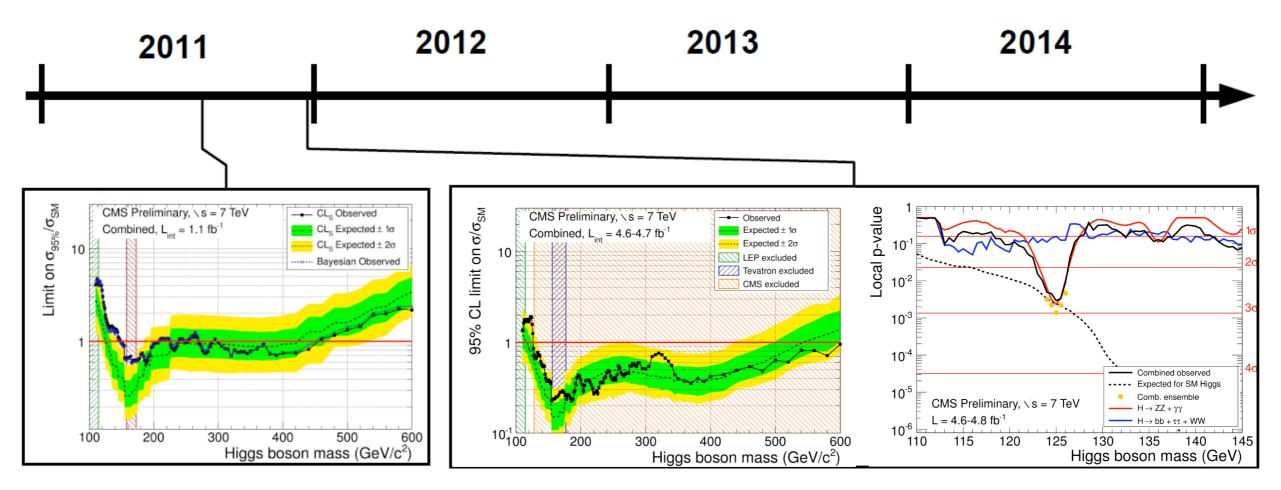
### First combined exclusion limits



First 1fb<sup>-1</sup> (7TeV): no Higgs boson between 160 and 500GeV

EPS-HEP '11 Lepton-Photon '11

### First hints



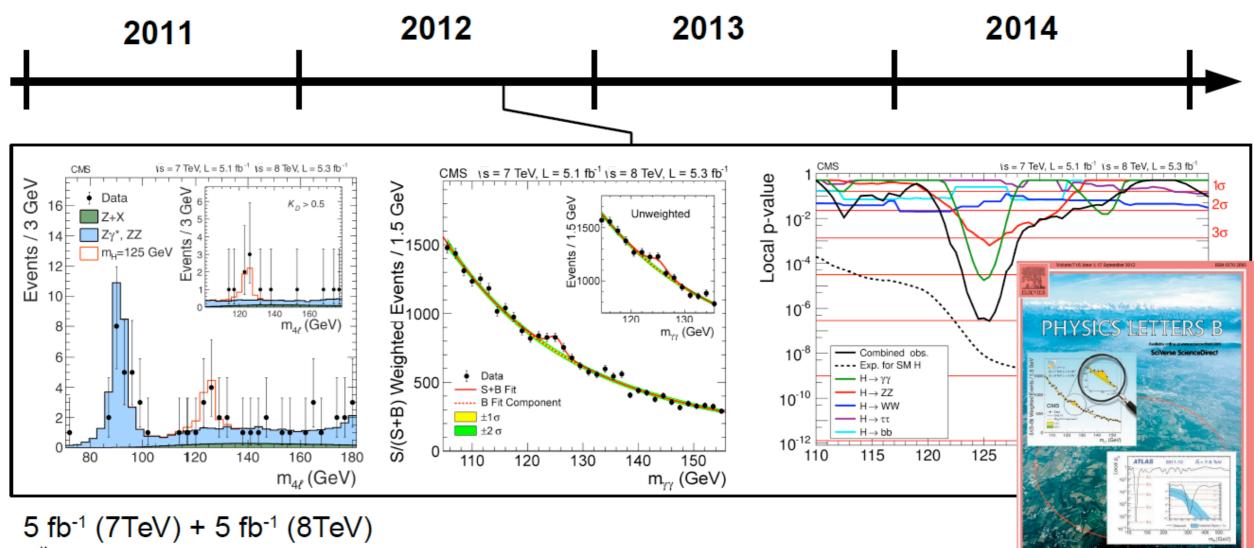
First 5fb<sup>-1</sup> (7TeV):

SM Higgs boson excluded for  $127 < m_{H} < 600 \text{GeV}$ 

Excess (local significance 2.8 $\sigma$ ) for m<sub>H</sub>~125GeV

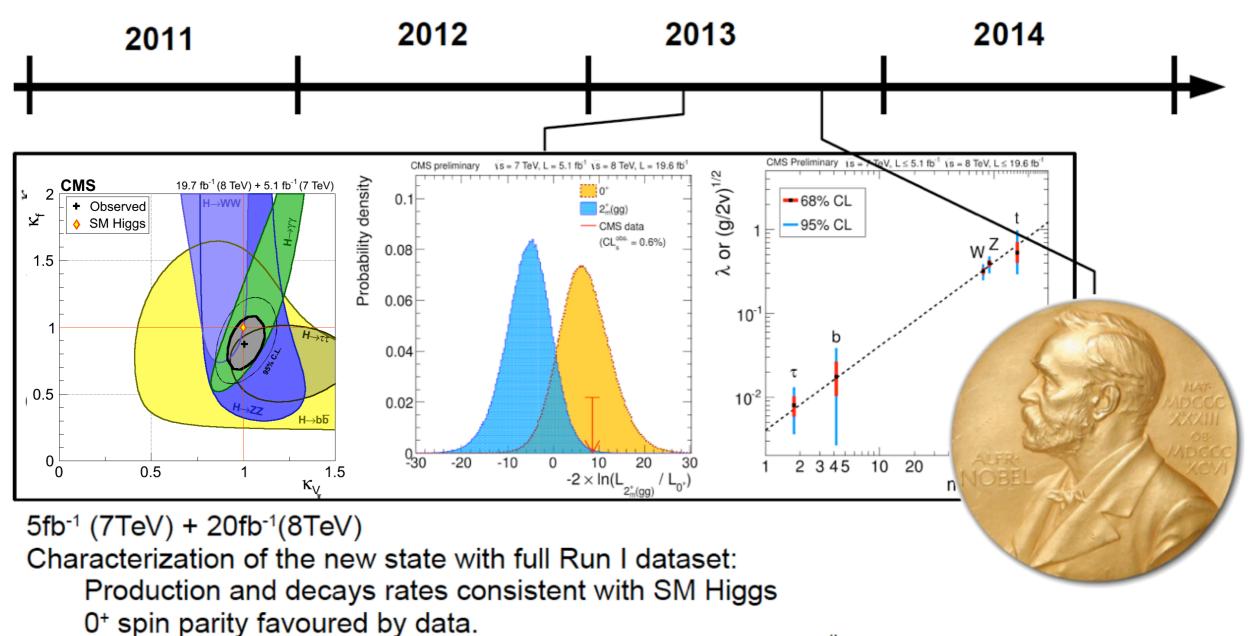
CMS/ATLAS Higgs Jamboree Moriond 2012

### Evidence of a new boson



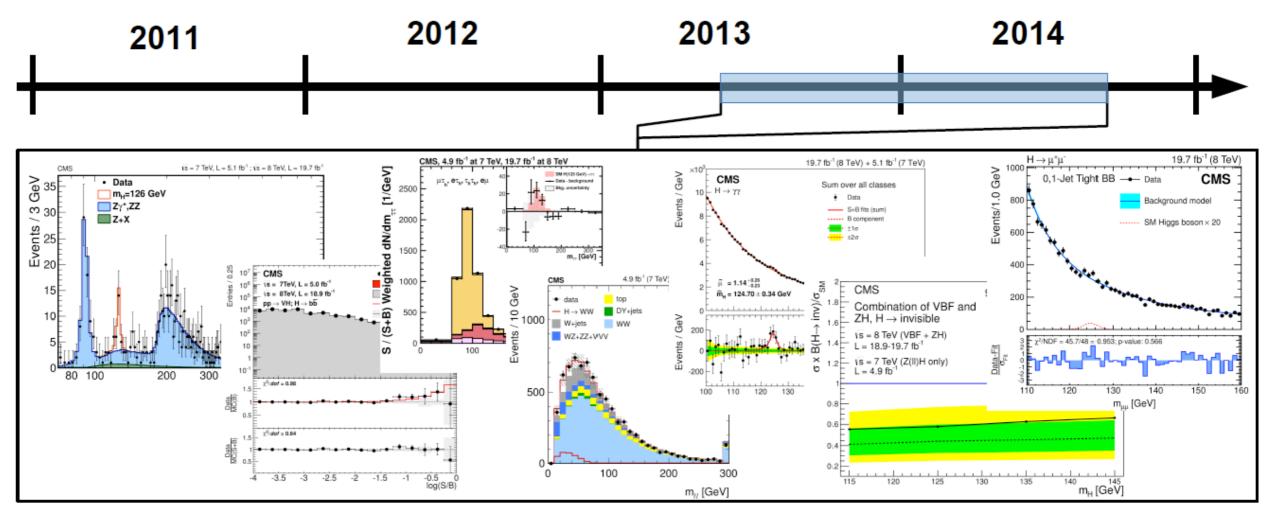
4<sup>th</sup> July 2012: CMS and ATLAS announce Evidence for a new boson.

### Identification of the Higgs boson



8<sup>th</sup> October 2013: Nobel Prize for Physics awarded to prof. Higgs and Englert.

### The Run I legacy



5fb<sup>-1</sup> (7TeV) + 20fb<sup>-1</sup>(8TeV)

Final results on Run I full dataset published 1-2 years after the discovery

of the new boson.

Ultimate precision for this dataset attained.

Preliminary combined analysis of all channels presented in July 2014.

# 2012-2014 remarkably intense and exciting years for particle physics





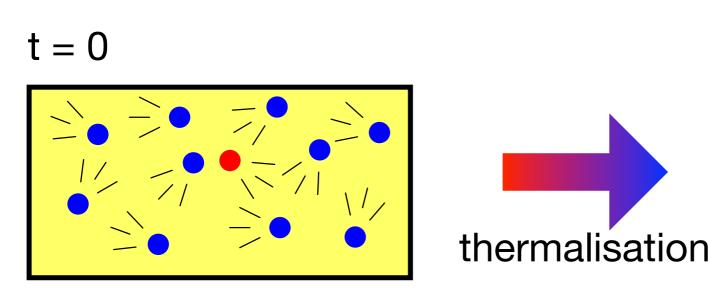
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- the Higgs mass receives corrections from vacuum fluctuations
- the size of the correction should be proportional to the maximum allowed energy M<sub>Planck</sub>, M<sub>GUT</sub>, . . .
- $M_H \ll M_{Planck}$  requires fine-tuning up to 17 digits or New Physics!

### Analogy with thermal fluctuations



At large t expect to have 
$$E_{\bullet} \sim E_{\bullet}$$
  
While the observation is  $E_{\bullet} \sim 10^{-17} E_{\bullet}$ 

While there is no inconsistency, it just seems hard to believe!

# Explanations for gauge hierarchy

 In the analogy: natural explanation could be that red does not really interact with blue because the interaction is screened

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- In the Higgs case: similarly, the interaction could be screened by new forces/particles

# Explanations for gauge hierarchy

- In the analogy: natural explanation could be that red does not really interact with blue because the interaction is screened
- In the Higgs case: similarly, the interaction could be screened by new forces/particles
  - A variety of possible explanations exist to protect the Higgs mass from having a sensitivity to high-energy scales
  - (supersymmetry, technicolour, Randall-Sundrum warped space, pseudo-Goldstone Higgs, Little Higgs, ... )

Currently these are all speculations. Only experimental data can discriminate between the predictions of various models

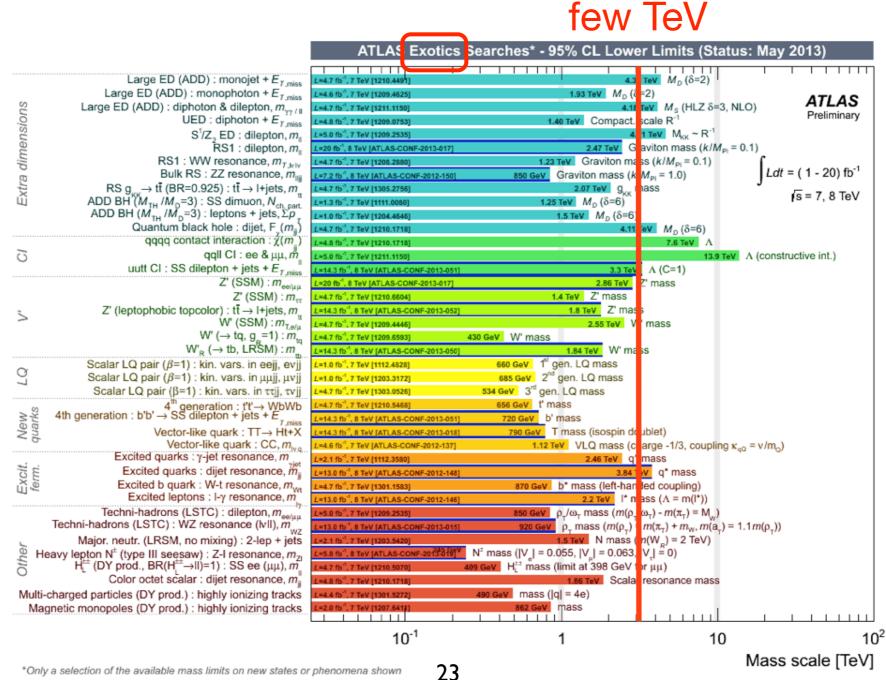
### Status of New Physics searches

### Unfortunately, direct searches are so far not successful

ATLAS SUSY Searches* - 95% CL Lower Limits $1 \text{ TeV}$ ATLAS Preliminary $\sqrt{s} = 7, 8 \text{ TeV}$							
	Model	$e,\mu,\tau,\gamma$	Jets	$E_{ m T}^{ m miss}$	∫£ dt[ft	<sup>1</sup> ] Mass limit	Reference
Inclusive Searches	$\begin{array}{c} \text{MSUGRA/CMSSM} \\ \text{MSUGRA/CMSSM} \\ \text{MSUGRA/CMSSM} \\ \overline{q}\bar{q}, \overline{q} \rightarrow q \overline{V}_1^0 \\ \overline{g}\bar{x}, \overline{g} \rightarrow q \overline{q} \overline{V}_1^0 \\ \overline{g}\bar{x}, \overline{g} \rightarrow q \overline{V}_1^0 \\ \overline{g}\bar$	$\begin{matrix} 0 \\ 1 \ e, \mu \\ 0 \\ 0 \\ 1 \ e, \mu \\ 2 \ e, \mu \\ 2 \ e, \mu \\ 1 - 2 \ \tau + 0 - 1 \ \ell \\ 2 \ \gamma \\ 1 \ e, \mu + \gamma \\ \gamma \\ 2 \ e, \mu (Z) \\ 0 \end{matrix}$	2-6 jets 3-6 jets 7-10 jets 2-6 jets 2-6 jets 3-6 jets 0-3 jets 0-3 jets mono-jet	Yes Yes - Yes Yes Yes Yes Yes Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.3	$\begin{array}{c c c c c c c c c c c c c c c c c c c $	1405.7875 ATLAS-CONF-2013-062 1308.1841 1405.7875 1405.7875 ATLAS-CONF-2013-062 ATLAS-CONF-2013-089 1208.4688 1407.0603 ATLAS-CONF-2014-001 ATLAS-CONF-2012-144 1211.1167 ATLAS-CONF-2012-152 ATLAS-CONF-2012-152
3 <sup>rd</sup> gen. <u>§</u> med.	$\tilde{g} \rightarrow b \bar{b} \tilde{\chi}_{1}^{0}$ $\tilde{g} \rightarrow t \bar{t} \tilde{\chi}_{1}^{0}$ $\tilde{g} \rightarrow t \bar{t} \tilde{\chi}_{1}^{0}$ $\tilde{g} \rightarrow b \bar{t} \tilde{\chi}_{1}^{0}$	0 0 0-1 e,μ 0-1 e,μ	3 b 7-10 jets 3 b 3 b	Yes Yes Yes Yes	20.1 20.3 20.1 20.1	k         1.2:         TeV         m(k <sup>2</sup> <sub>1</sub> )<400 GeV           k         1.1 TeV         m(k <sup>2</sup> <sub>1</sub> )<350 GeV	1407.0600 1308.1841 1407.0600 1407.0600
3 <sup>rd</sup> gen. squarks direct production	$ \begin{array}{l} \tilde{b}_{1}\tilde{b}_{1}, \tilde{b}_{1} \rightarrow b\tilde{\chi}_{1}^{0} \\ \tilde{b}_{1}\tilde{b}_{1}, \tilde{b}_{1} \rightarrow t\tilde{\chi}_{1}^{-} \\ \tilde{r}_{1}\tilde{r}_{1}(\text{light}), \tilde{r}_{1} \rightarrow b\tilde{\chi}_{1}^{-} \\ \tilde{r}_{1}\tilde{r}_{1}(\text{light}), \tilde{r}_{1} \rightarrow b\tilde{\chi}_{1}^{0} \\ \tilde{r}_{1}\tilde{r}_{1}(\text{light}), \tilde{r}_{1} \rightarrow b\tilde{\chi}_{1}^{0} \\ \tilde{r}_{1}\tilde{r}_{1}(\text{medium}), \tilde{r}_{1} \rightarrow t\tilde{\chi}_{1}^{0} \\ \tilde{r}_{1}\tilde{r}_{1}(\text{meday}), \tilde{r}_{1} \rightarrow t\tilde{\chi}_{1}^{0} \\ \tilde{r}_{1}\tilde{r}_{1}\tilde{r}_{1} \rightarrow t\tilde{\chi}_{1}^{0} \\ \tilde{r}_{1}\tilde{r}_{1}\tilde{r}_{1} \rightarrow t\tilde{\chi}_{1}^{0} \\ \tilde{r}_{1}\tilde{r}_{1}\tilde{r}_{1}\tilde{\tau}_{1} \rightarrow t\tilde{\chi}_{1}^{0} \\ \tilde{r}_{1}\tilde{r}_{1}\tilde{r}_{1}\tilde{\tau}_{1} \rightarrow t\tilde{\chi}_{1}^{0} \\ \tilde{r}_{1}\tilde{r}_{1}\tilde{r}_{1}\tilde{\tau}_{2}\tilde{\tau}_{1}\tilde{\tau}_{1}\tilde{\tau}_{2} \\ \tilde{r}_{1}\tilde{r}_{1}\tilde{r}_{2}\tilde{\tau}_{2}\tilde{\tau}_{1}\tilde{\tau}_{2} \rightarrow t\tilde{t} + Z \end{array} $	$\begin{array}{c} 0\\ 2\ e,\mu\ ({\rm SS})\\ 1-2\ e,\mu\\ 2\ e,\mu\\ 2\ e,\mu\\ 0\\ 1\ e,\mu\\ 0\\ 1\ e,\mu\\ 0\\ 3\ e,\mu\ (Z) \end{array}$	2 b 0-3 b 1-2 b 0-2 jets 2 jets 2 b 1 b 2 b tono-jet/c-t 1 b 1 b	Yes Yes Yes Yes Yes Yes Yes ag Yes Yes Yes	20.1 20.3 4.7 20.3 20.3 20.1 20 20.1 20.3 20.3 20.3	$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	1308.2631 1404.2500 1208.4305, 1209.2102 1403.4853 1403.4853 1308.2631 1407.0583 1406.1122 1407.0608 1403.5222 1403.5222
EW direct	$ \begin{array}{c} \tilde{\ell}_{LR} \tilde{\ell}_{LR}, \tilde{\ell} \rightarrow \ell \tilde{\chi}_1^0 \\ \tilde{\chi}_1^+ \tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow \ell \nu (\ell \tilde{\nu}) \\ \tilde{\chi}_1^+ \tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow \tau \nu (\tau \tilde{\nu}) \\ \tilde{\chi}_1^+ \tilde{\chi}_2^0 \rightarrow \tilde{\ell}_L \nu \tilde{\ell}_L ((\tilde{\nu}), \ell \tilde{\nu} \tilde{\ell}_L \ell (\tilde{\nu} \nu) \\ \tilde{\chi}_1^+ \tilde{\chi}_2^0 \rightarrow W \tilde{\chi}_1^0 D \tilde{\chi}_1^0 \\ \tilde{\chi}_1^+ \tilde{\chi}_2^0 \rightarrow W \tilde{\chi}_1^0 h \tilde{\chi}_1^0 \\ \tilde{\chi}_2^0 \tilde{\chi}_2^0 , \tilde{\chi}_2^0 \rightarrow \tilde{\ell}_R \ell \end{array} $	2 e,µ 2 e,µ 2 τ 3 e,µ 2-3 e,µ 1 e,µ 4 e,µ	0 0 - 0 2 b 0	Yes Yes Yes Yes Yes Yes Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.3	$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	1403.5294 1403.5294 1407.0350 1402.7029 ATLAS-5294, 1402.7029 ATLAS-CONF-2013-093 1405.5086
Long-lived particles	$\begin{array}{l} \text{Direct} \tilde{X}_{1}^{+} \tilde{X}_{1}^{-} \text{ prod., long-lived } \tilde{\chi}_{1}^{\pm} \\ \text{Stable, stopped } \tilde{g} \text{ R-hadron} \\ \text{GMSB, stable } \tilde{\tau}, \tilde{X}_{1}^{0} \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \\ \text{GMSB, } \tilde{X}_{1}^{0} \rightarrow \gamma \tilde{G}, \text{ long-lived } \tilde{\chi}_{1}^{0} \\ \tilde{q}\tilde{q}, \tilde{X}_{1}^{0} \rightarrow qq\mu \text{ (RPV)} \end{array}$	Disapp. trk 0 μ) 1-2 μ 2 γ 1 μ, displ. vtx	1 jet 1-5 jets - -	Yes Yes - Yes -	20.3 27.9 15.9 4.7 20.3	$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	ATLAS-CONF-2013-069 1310.6584 ATLAS-CONF-2013-058 1304.6310 ATLAS-CONF-2013-092
RPV	$ \begin{array}{l} LFV pp \rightarrow \bar{\nu}_\tau + X, \bar{\nu}_\tau \rightarrow e + \mu \\ LFV pp \rightarrow \bar{\nu}_\tau + X, \bar{\nu}_\tau \rightarrow c(\mu) + \tau \\ Bilinear \ RPV \ CMSSM \\ \bar{\chi}_1^+ \bar{\chi}_1^-, \bar{\chi}_1^+ \rightarrow W \bar{\chi}_1^0, \bar{\chi}_1^0 \rightarrow ee\bar{\nu}_\mu, e\mu\bar{\nu}_e \\ \bar{\chi}_1^+ \bar{\chi}_1^-, \bar{\chi}_1^+ \rightarrow W \bar{\chi}_1^0, \bar{\chi}_1^0 \rightarrow er\bar{\nu}_\tau, er\bar{\nu}_\tau \\ \bar{g} \rightarrow qqq \\ \bar{g} \rightarrow \bar{t}_1 t, \bar{t}_1 \rightarrow bs \end{array} $	$\begin{array}{c} 2 \ e, \mu \\ 1 \ e, \mu + \tau \\ 2 \ e, \mu \ (\text{SS}) \\ 4 \ e, \mu \\ 3 \ e, \mu + \tau \\ 0 \\ 2 \ e, \mu \ (\text{SS}) \end{array}$	- 0-3 b - - 6-7 jets 0-3 b	- Yes Yes - Yes	4.6 4.6 20.3 20.3 20.3 20.3 20.3	\$\vec{v}_r\$         1.61 TeV $\lambda_{111}^*=0.10, \lambda_{132}=0.05$ \$\vec{v}_r\$         1.1 TeV $\lambda_{311}^*=0.10, \lambda_{122}=0.05$ \$\vec{v}_r\$         1.1 TeV $\lambda_{311}^*=0.10, \lambda_{122}=0.05$ \$\vec{v}_r\$         1.5 TeV         m(\$\vec{v}_{2,11}=0.10, \lambda_{122}=0.05           \$\vec{v}_r\$         1.5 TeV         m(\$\vec{v}_{2,11}=0.10, \lambda_{122}=0.05           \$\vec{v}_r\$         1.5 TeV         m(\$\vec{v}_{2,11}=0.10, \lambda_{122}=0.05           \$\vec{v}_r\$         750 GeV         m(\$\vec{v}_{1,11}=0.10, \lambda_{122}=0.05           \$\vec{v}_r\$         450 GeV         m(\$\vec{v}_{1,11}=0.10, \lambda_{122}=0.05           \$\vec{v}_r\$         450 GeV         m(\$\vec{v}_{1,11}=0.2, xm(\$\vec{v}_{1,11}, \lambda_{121}=0           \$\vec{v}_r\$         916 GeV         BR(t)=BR(b)=BR(c)=0%           \$\vec{v}_r\$         850 GeV         BR(t)=BR(b)=BR(c)=0%	1212.1272 1212.1272 1404.2500 1405.5086 1405.5086 ATLAS-CONF-2013-091 1404.250
Other	Scalar gluon pair, sgluon $\rightarrow q\bar{q}$ Scalar gluon pair, sgluon $\rightarrow t\bar{t}$ WIMP interaction (D5, Dirac $\chi$ )	$2 e, \mu (SS)$ $\sqrt{s} = 8 \text{ TeV}$	4 jets 2 b mono-jet	- Yes Yes 8 TeV	4.6 14.3 10.5	sgluon         100-287 GeV         incl. limit from 1110.2693           sgluon         350-800 GeV         1           M* scale         704 GeV         m(χ)<80 GeV, limit of <687 GeV for D8	1210.4826 ATLAS-CONF-2013-051 ATLAS-CONF-2012-147
		$y_s = 8$ TeV artial data		data		10 <sup>-1</sup> 1 Mass scale [TeV]	

## Status of New Physics searches

### Unfortunately, direct searches are so far not successful

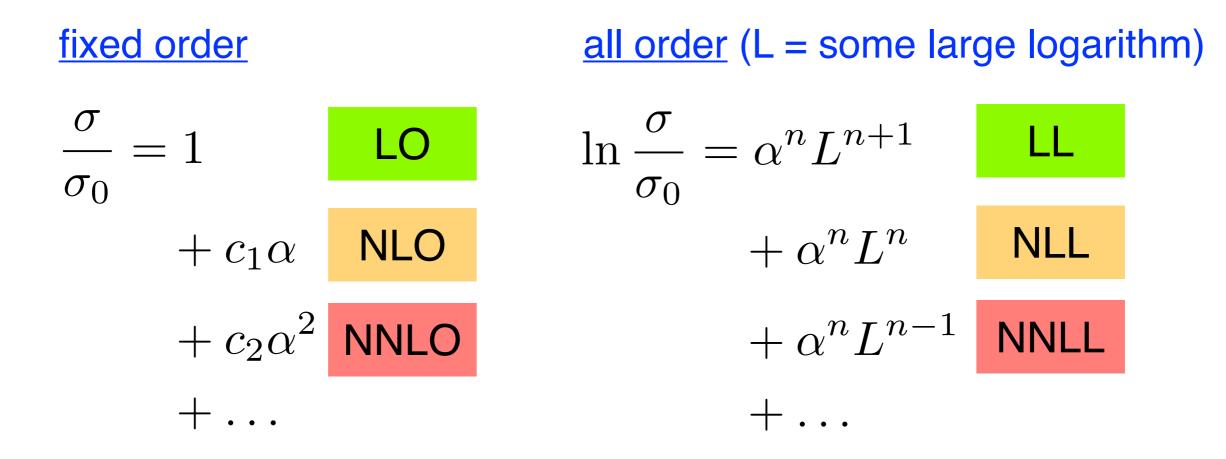


## Future direction

- Run II at almost twice the energy will allow us to push the reach of these direct searches considerably
- Yet, the possibility must be taken into account, that no new state is produced directly (simply because the energy is not enough)
- Indirect searches and precision tests more prominent in Run II
- the Higgs sector in particular will undergo scrupulous precision tests (remember: given the Higgs mass everything is predicted in the SM, so everything can and must be tested)
- Precision tests require both accurate measurements and precise theoretical predictions

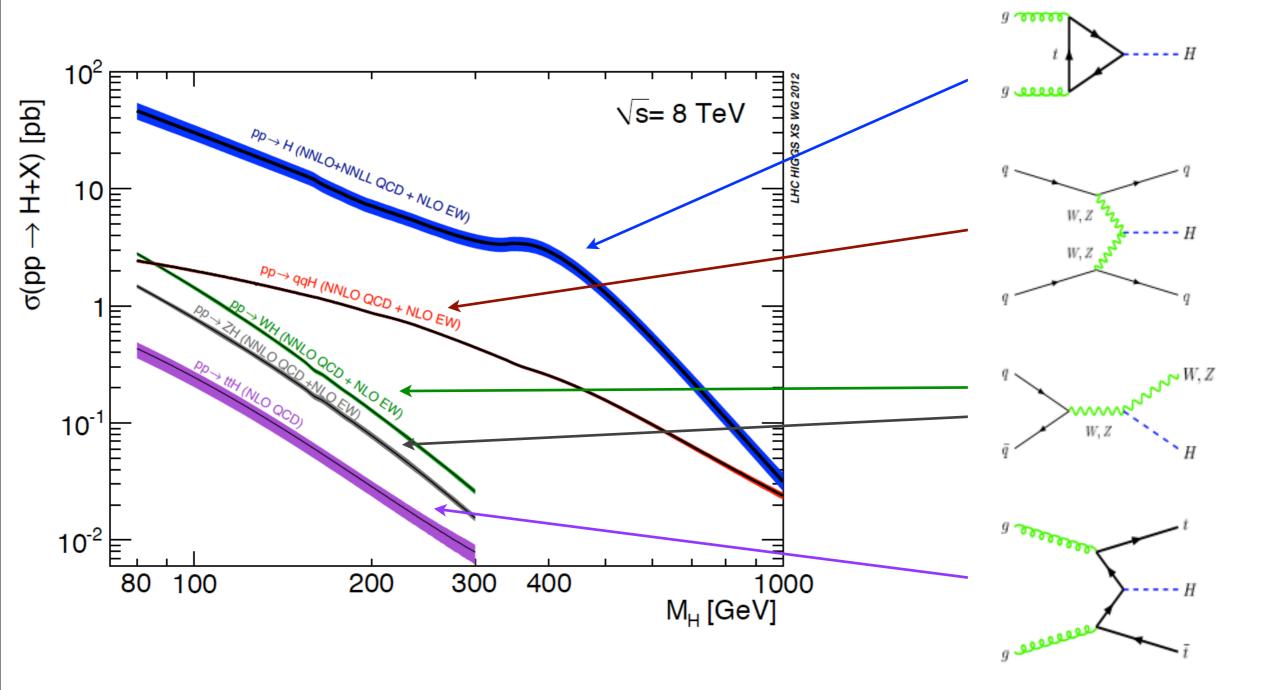
# Precision through Perturbation

At the LHC, QCD and electroweak (EW) interactions are weak. We can compute perturbative expansions in the (small) coupling. Higher-order terms will improve predictions. Different expansions:



<u>QCD:</u>  $\alpha \sim 0.1$  expect NLO to be O(10%) correction, NNLO O(1%) ... <u>EW:</u>  $\alpha \sim 0.01$  expect NLO to be O(1%) correction, ...

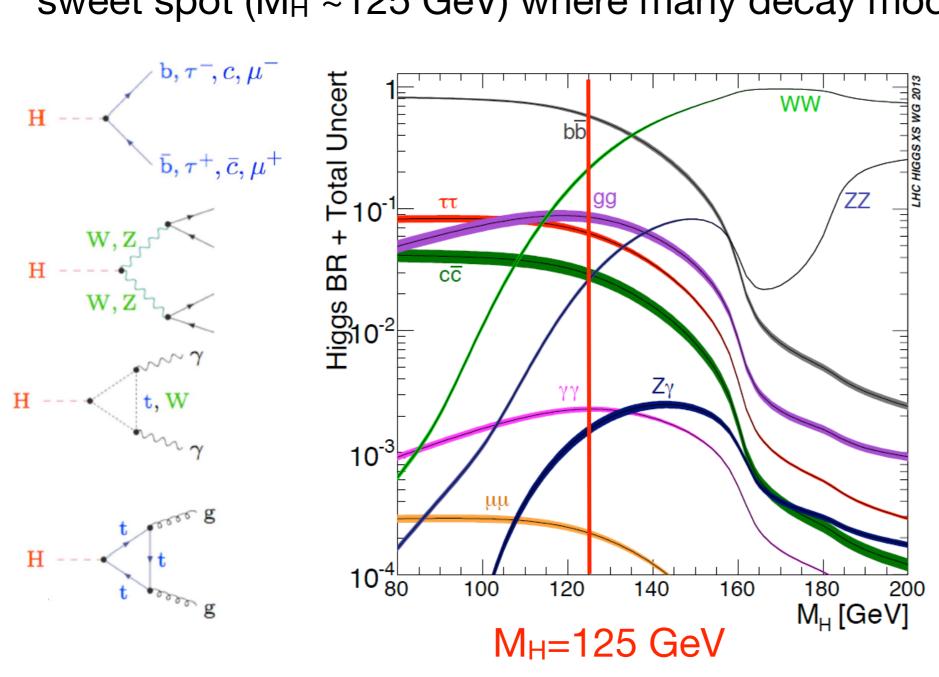
### Higgs production at the LHC



## Higgs decay modes

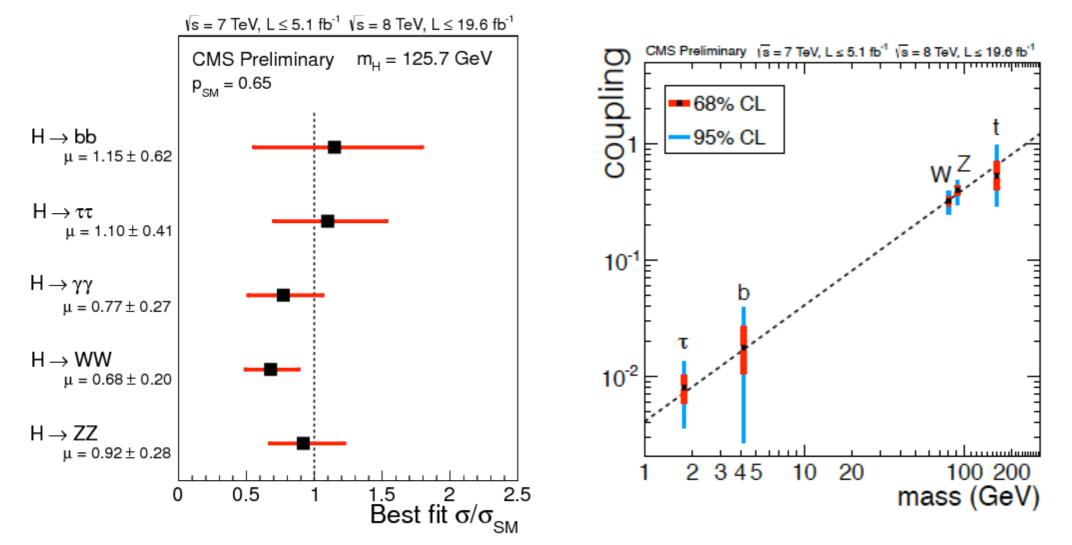
i.e. what is actually seen in detectors

Higgs decays very very quickly.... fortunately, its mass lies in a sweet spot (M<sub>H</sub>  $\sim$ 125 GeV) where many decay modes are available



# Status of Higgs measurements

Precision Higgs phenomenology (based on full 7 and 8 TeV data) shows so far no departure from a plain SM Higgs boson pattern

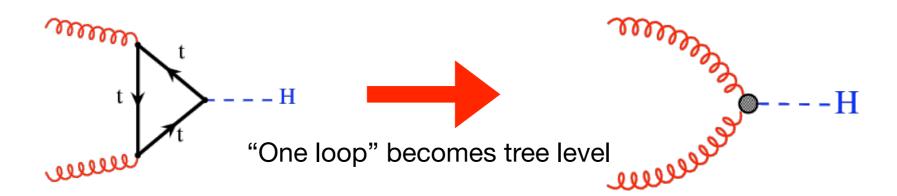


Run II at the LHC about to start: focus will be on accurate Higgs measurements using high-precision theory

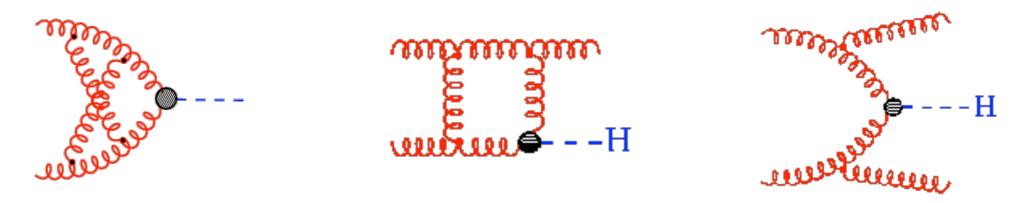
# Inclusive Higgs production

<u>One example:</u> the simplest (and dominant) Higgs production mechanism via gluon-gluon fusion (no decays). How well do we know this cross-section?

Most calculations based on the large mt-limit effective theory:

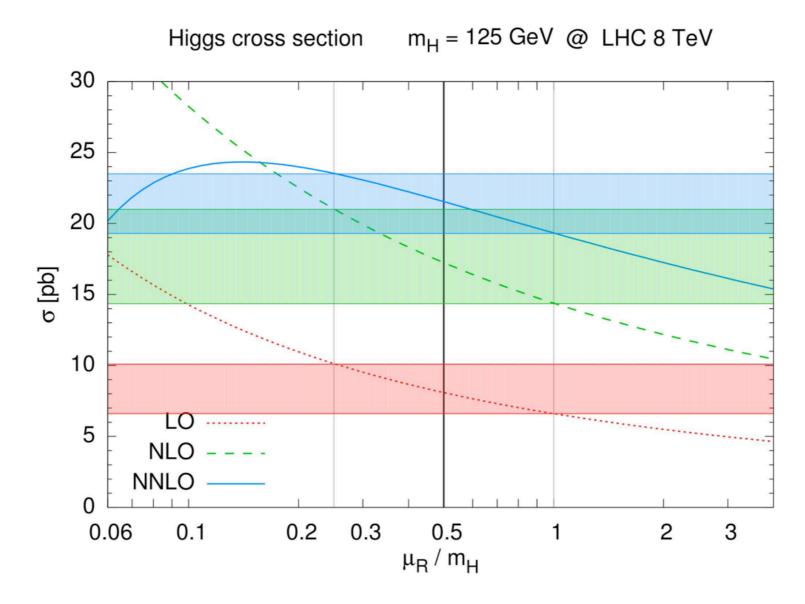


In this limit, NNLO corrections known for many years:



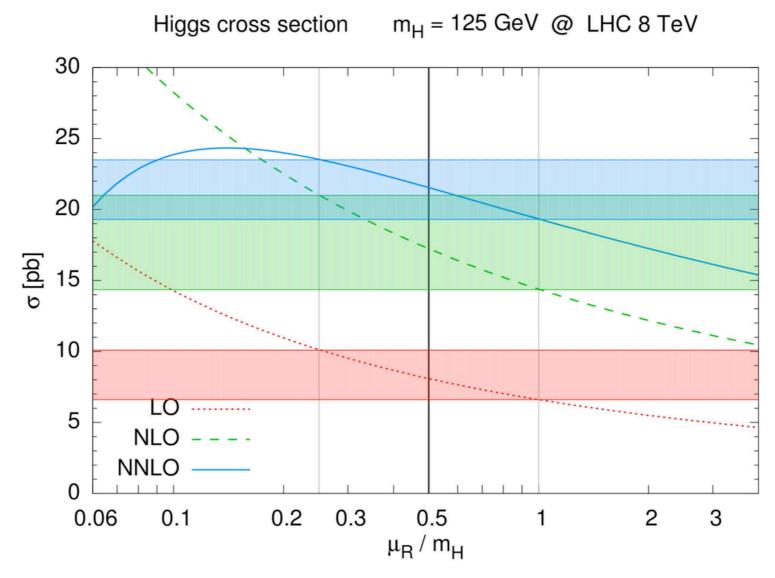
29

# Inclusive Higgs production



- perturbative series for  $gg \rightarrow H$  converges very slowly
- renormalization scale variation (commonly used to estimate theory uncertainty) underestimates the shift to the next order

# Inclusive Higgs production

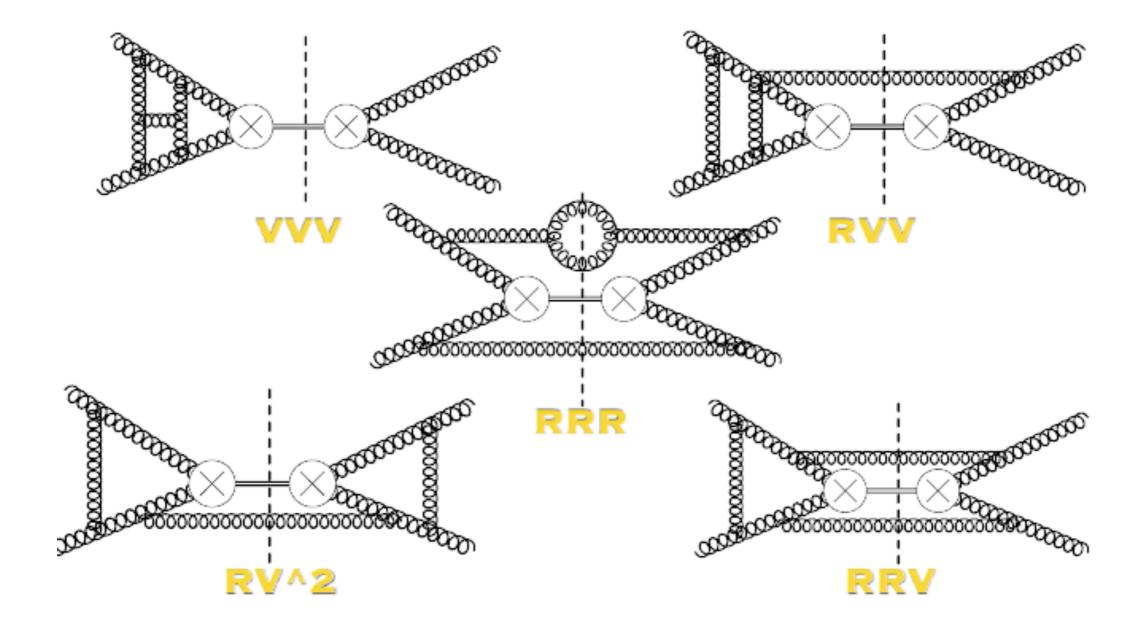


### Two ways to go:

- try to compute higher orders approximately (resummations)
- try to compute exact N<sup>3</sup>LO i.e.  $O(\alpha_s^3)$  correction. Is it that difficult?

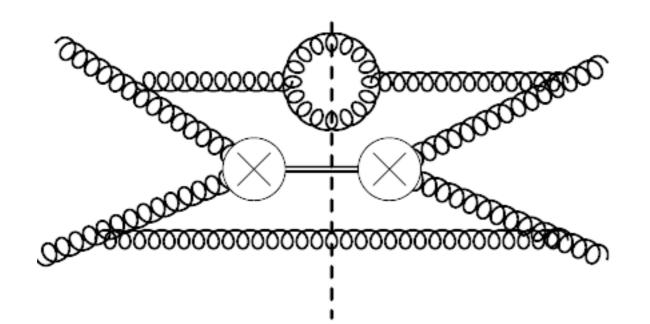
## Facts about N<sup>3</sup>LO

• O(100000) interference diagrams (1000 at NNLO)



## Facts about N<sup>3</sup>LO

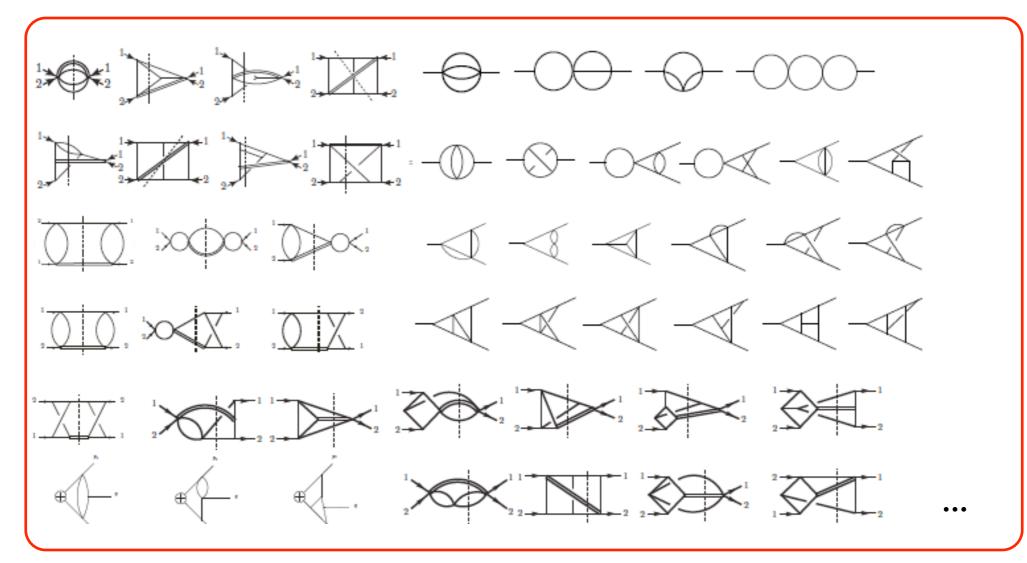
- O(100000) interference diagrams (1000 at NNLO)
- 68273802 loop and phase space integrals (47000 at NNLO)



+ 68273801 integrals

## Facts about N<sup>3</sup>LO

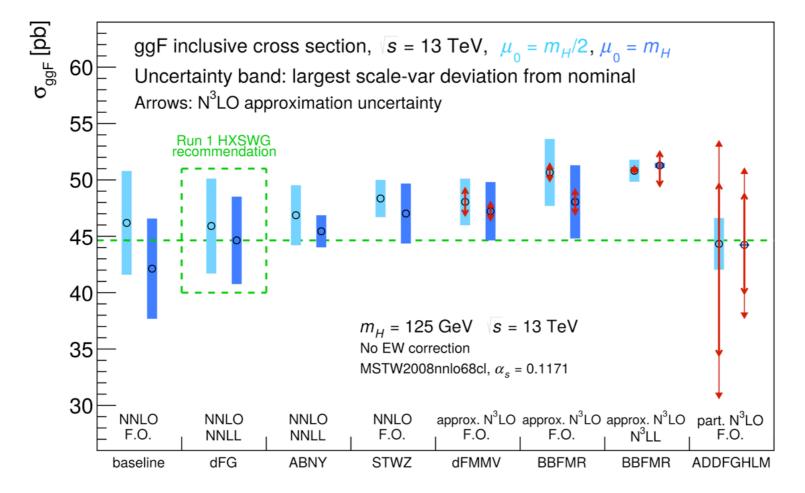
- O(100000) interference diagrams (1000 at NNLO)
- 68273802 loop and phase space integrals (47000 at NNLO)
- about 1000 master integrals (26 at NNLO)



## Approximate N<sup>3</sup>LO

Approximate N<sup>3</sup>LO from different groups (possibly including higher order logarithmic terms) together with their uncertainty

approx N<sup>3</sup>LO from Anastasiou et al '14



What is the most reasonable approximation based on our current knowledge? Central value and size of uncertainty hotly debated!

### THE SOCRATIC PROBLEM

HOW DO WE ESTIMATE THE AMOUNT OF OUR IGNORANCE?

[21δ] ἐντεῦθεν οὖν τούτω τε ἀπηχθόμην καὶ πολλοῖς τῶν παρόντων: πρὸς ἐμαυτὸν δ' οὖν ἀπιών ἐλογιζόμην ὅτι τοὐτου μὲν τοῦ ἀνθρώπου ἐγώ σοφώτερός είμι: κινδυνεύει μὲν γὰρ ἡμῶν οὐδέτερος οὐδὲν καλὸν κἀγαθὸν εἰδέναι, ἀλλ' οὐτος μὲν οἴεταί τι εἰδέναι οὐκ εἰδώς, ἐγώ δέ, ώσπερ οὖν οὐκ οἶδα, οὐδὲ οἴομαι: ἔοικα γοῦν τούτου γε σμικρῷ τινι αὐτῷ τούτω σοφώτερος εἶναι, ὅτι ἂ μὴ οἶδα οὐδὲ οἴομαι εἰδέναι. ἐντεῦθεν ἐπ' ἄλλον ἦα τῶν ἐκείνου δοκούντων σοφωτέρων εἶναι καί

### Plato. Platonis Opera, ed. John Burnet. Oxford University Press. 1903.

I am wiser than this man; for neither of us really knows anything fine and good, but this man thinks he knows something when he does not, whereas I, as I do not know anything, do not think I do either. I seem, then, in just this little thing to be wiser than this man at any rate, that what I do not know I do not think I know either.

talk given by S. Forte at the 8th Workshop of the Higgs Cross Section Working Group, 22<sup>nd</sup> Jan. '15

## Conclusions

- Fantastic data available and expected from LHC (restarts operation for three years this summer)
- Higgs discovery was a true milestone for particle physics, but also leaves many questions open (hierarchy problem, naturalness, ...)
- Run II will focus on precision studies: what does the future hold?

