HOW TO USE RANDOM MATRIX THEORY IN THE DETECTION OF GRAVITATIONAL WAVES



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SOME HISTORY ON THE STATISTICAL DISTRIBUTION OF THE WIDTHS AND SPACINGS OF NUCLEAR RESONANCE LEVELS

By EUGENE P. WIGNER

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One of the first applications of random matrices to physics is due to Jenö Pál Wigner



Problem: how to determine the energy levels of heavy nuclei (such as Uranium)

Difficulties: it is a very complicated system, we do not know how to write the Hamiltonian

$$H\Psi_n = E_n\Psi_n$$

SOME HISTORY

Wigner observed that the space, s, between neighbouring energy levels were governed by the following probability density



I Idea: interpret the distance between energy levels as the distance between eigenvalues of the Hamiltonian of the system

II Idea: define an ensemble of Hamiltonians with statistical properties as those that the real Hamiltonian might have if it could be written explicitly

RANDOM MATRIX THEORY

give up intrinsic details of the system in order to collect information on average properties

CSe 73 h L fJI_ar Whe260 080 Len DF 7 x 937984 qZ14C7 187 b80 28 u N 888598 819f13 x18 Z42 JT 5 1.87.9 Lu7 az8 47u - p 82 77C 13772 gW3 17m 87 J Ev iTE 19x3= aC9 4tk 7E 7 uj 147 76x88 91C JTB Sn v xh Env 6Sual tqu Rb 10 k Vk ala 1bt3V u80	THEOR	YINA	NUTSF	6b so (7 90x50590 6b so (7 90x50590 48 4s \$ 1 730+9352 CF L9 B c 9507Ri70 16 s(F) e3597043 18 T\ 78 [(@1231s) 18 T\ 78 [(@1231s) 15 7\\ 000; D5+6 4x s2'0) (40669183 15 sTC7r 87(List)
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Gaussian Ensembles: independent, identically distributed Gaussian entries

<u>G</u>aussian <u>O</u>rthogonal <u>E</u>nsemble: real symmetric matrices <u>G</u>aussian <u>U</u>nitary <u>E</u>nsemble: complex Hermitian matrices <u>G</u>aussian <u>Symplectic E</u>nsemble: quaternionic Hermitian matrices



Dyson index β : counts the number of real components per matrix element

The eigenvalue spacing distribution of these ensembles is approximated by the Wigner surmise

THEORY IN A NUTSHELL

These matrices can be diagonalized as

$$M = U\Lambda U^{-1}$$
angular radial
decomposition
$$\Lambda = (\lambda_1, \dots, \lambda_n) \in \mathbb{R}^n$$

The matrix U is

- a real orthogonal matrix if $M \in \text{GOE}$
- a complex unitary matrix if $M \in GUE$
- a complex symplectic matrix if $M \in GSE$

Wigner semicircle law

Let X be a n-dimensional complex Hermitian matrix with

- $(X_{ij})_{i < j}$ i.i.d. with zero mean and finite variance σ ,
- (X_{kk}) i.i.d. with bounded mean and variance.

Then the eigenvalues of $n^{-1/2}X$ tend in distribution to the semicircle law



Gaussian Ensembles follow this eigenvalue distribution

APPLICATION TO GW DETECTION



Gravitational waves: perturbations of the gravitational field that travel at the speed of light

14 September 2015: FIRST direct detection



- two low mass black hole merger



first detection of this type of event



effective proof that these phenomena may take place within the current age of the universe

GRAVITATIONAL WAVE ASTRONOMY ERA

APPLICATION TO GW DETECTION

Gravitational wave <u>stochastic</u> background: superposition of signals from unresolved sources Gaussian

Problem: the stochastic background data may hide a gravitational wave signal

We need a method to look for long-memory effect in the background

Idea: RANDOM MATRICES

Random matrix approach in search of weak signals immersed in background noise D. Grech & J. Miśkiewicz, Europhysics Letters, Volume 97, Issue 3, February 2012

I STEP: BUILD THE MATRICES AND THE ENSEMBLE

 $\{x_i\}$ • time series • i=1,...,N+1 • N>>1 • increment time series • we have $\Delta x_i = x_{i+1} - x_i$

$$s_k = \left\{ \Delta x_{(k-1)L+1}, \dots \Delta x_{kL} \right\}$$

Renormalize according to

$$s_k \to \hat{s}_k = \frac{s_k}{\sqrt{L}\sigma_k}$$

•

I STEP: BUILD THE MATRICES AND THE ENSEMBLE

- the first L subseries build the first LxL matrix
- the second L subseries build the second LxL matrix
- thus, we get N/L^2 matrices $M^{(n)}$ (n=1,...,N/L^2)

Strategy: Further steps rely on examination of eigenvalue spectra properties for the ensamble of symmetrized matrices with (i,j) entries $(M_{ij}^{(n)} + M_{ji}^{(n)})/2$ and on a comparison with a spectrum known **a priori**. Any distortion from this **a priori** spectrum is interpreted as a weak signal hidden in the background

II STEP:SIMPLE EXAMPLE

- assume that the noise is pure white noise
- add a sinusoidal signal for different s/n
- if only white noise the eigenvalue spectrum should disappear for $|\lambda|>2$

$$\rho(\lambda) = \frac{1}{2\pi}\sqrt{-\lambda^2 + 4}$$

II STEP:SIMPLE EXAMPLE

ensemble of 1000 matrices of size 200x200



III STEP: CHARACHTERIZE THE NOISE OF A REAL EXPERIMENT

- NAUTILUS, ultra-cryogenic resonant gravitational wave detector
- built to detect gravitational bursts
- no longer running, yet it collected some data



III STEP: CHARACHTERIZE THE NOISE OF A REAL EXPERIMENT

- ensemble of 1000 matrices of dimension 200x200
- triangular shape in the spectrum due to the specifics of the instrument
- if we shuffle the data we get Wigner



III STEP: CHARACHTERIZE THE NOISE OF A REAL EXPERIMENT

if we compare real noise and simulated noise we see a difference in the tail



this can be important for the detection of periodic weak signals

IV STEP: A SIMPLE EXAMPLE



V STEP: A PHYSICAL EXAMPLE

- gravitational wave signal from a freely precessing antisymmetric star
- the level of background noise was reduced by a factor 10^k (k=0,-1,-2,-3)



V STEP: A PHYSICAL EXAMPLE



good performance of the method

- weak signal may be hidden in the background data
- random matrix approach can help in revealing this coded signal
- compare the density distribution of the eigenvalues with a known one
- pay attention to the tails

